

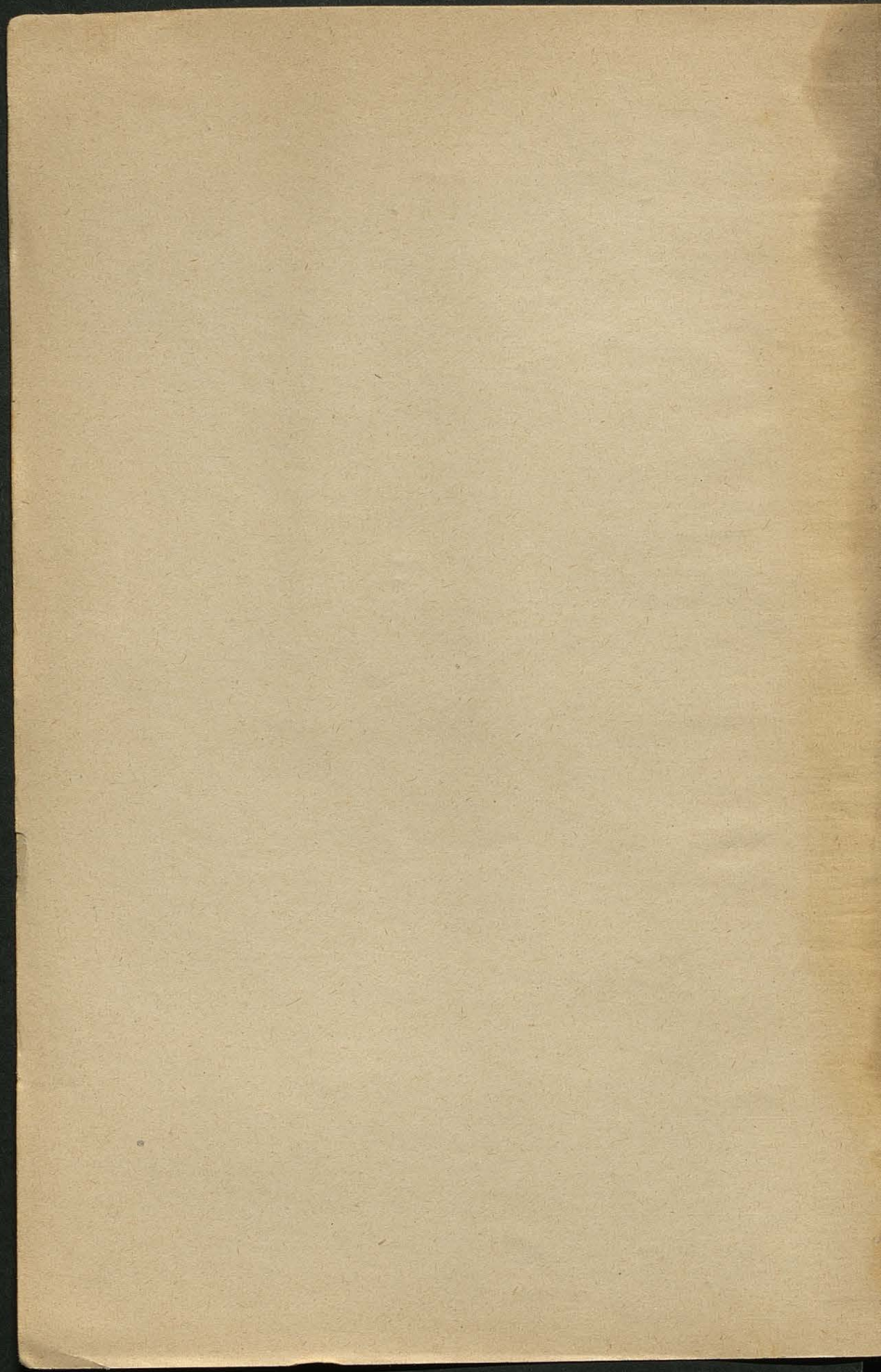
9451



10

6973

J. LUXANBY
WIEN
IV. Wiedener Hauptstr. 29



1.11

Das Voltier ges $\frac{1}{2}$ $\frac{1}{2}$

für $100 / - 100 \text{ } \frac{1}{2} \text{ } \frac{1}{2} \text{ } \frac{1}{2} \text{ } \frac{1}{2}$

Schumann, Palmieri, Secchi

$V \leftarrow \text{sch. } \frac{1}{2}$



$$V + \frac{q_1}{R} = \text{Volte } \frac{1}{2} \text{ } \frac{1}{2} = 0.$$

$0 = \frac{1}{2}$

$$V_2 + \frac{q_2}{R} = 0$$

die $\frac{1}{2}$ ist $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

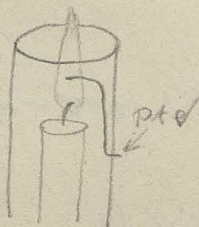
$\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

$q_2 - q_1 = \frac{1}{2}$; $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

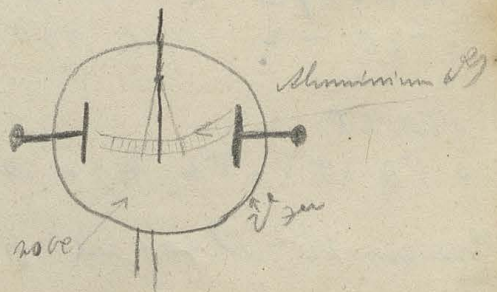
$$q_2 - q_1 = R \cdot (V_1 - V_2)$$

die $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$; $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$
 $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

Genes'side Anordnung:

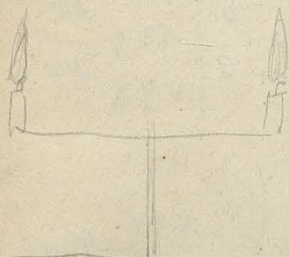


Genes'side Elektrometer



Anordnung $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

1. 10. 2007; ...
 ...
 ... 4

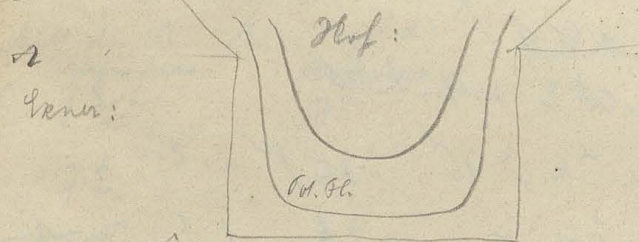


Im Luftballon ...
 10. ...

...
 ...
 ... x 20m (in ...)

e ... + ...
 ...
 ...

...
 ...



2m	20m	48	68 K
	15m	17	32
	10m	7	11
	5m	2	5
	0	0	0

✓ Felswand:

40 m	5 m L 000	0
30 m		0
25 m		0

35 m	80 h
	80
	70

1/2 100 m:

230 v.

200

150

50 2/3 100 m Kissen v. g.

w. p. 2 5/8 v. m. v. g. des Kissenflecks 0 1/2.

w. v. e. - 100 20 m 0 2 v. - Reduktion 1/4.

1/4

1/4

w. v. m. = 0.1/2 p. [2 1/2 2 0. 2 0. 2 0. 2 0.]

2 1/2 p. 1 1/2 v. 1/2 p. 1/2

en v. 1/4 v. 1/2 p.

Unter 5 Böteln = 2 1/2 p. 1/2

2 1/2 p. 1/2 5-6 16. v. 1/2 p. 1/2 2 m

103 v. Red. p. 1/2 v. 1/2 p.; 1/2 p. 1/2 5 m. v. 1/2

v. 1/2 p. 1/2 v. 1/2 p. 5'0-

Jänner	991	voll postm
Febr	299	
März	294	
April	138	
Mai	110	
Juni	<u>102</u>	
Juli	123	
August	121	
September	121	
October	188	
November	260	
December	<u>470</u>	

absolute Nim 81 [im Juni]
Nax 638 [- Dec]

✓ 2/2 N 2 Nax 5 2 Nim.
~ Nim 2 1/2 e e ~ 1/2
2 - 1/2 N 1/2 N = Darm.
ab regente. 203
M. Perpignan

1	2	3	4	5	6	7	8	9	10	11	12
43	<u>40</u>	<u>39</u>	40	43	50	59	<u>63</u>	61	55	53	55
54	<u>53</u>	54	58	62	69	<u>72</u>	71	66	60	53	48

Umfahrt der A. A.	~ 1/2 h	Hand Nax	Früh Nim	Nachm Nim
W	1 3/4 h	7 h	2 1/2	2 h
W	7 1/2	7 1/2 h	3 h	2 h
W	7 1/2	8 h	3 h	3 h
W	8 h	6 3/4 h	3 h	12 h

2 W W 1/2 23ten Dec 2 1/2

1 e. Pol. für 2 Jhr 12

Wolfsbühel in 220 V. pro km Höhe best.

Wien 140

St. Sigen

Lido bei Venedig

Ceylon, Bombay 55 u

Sonnst. 1100 Lu. best. [gef. 10. Thunus]

Ägypt. Pyramide 1100 Exner

Waffelth. 10000 Pol. # e. t

Crejo 2/3 Bremen, 2 Chops, 2 Jhr u. Semmer

Bahnw. u. 2 Korr. u. b.

es 4. für 26 - 1 Thunus 100

Die 26

~ 7 + Pol. 100 u. 2 Jhr 2 Jhr 26

f. 2000 12/12 100 26 26

5 Jhr 100 2000

10. Kohn 200 2000 - 10000

26 - 50 V.

26 26 26 26 26 26 26 26 26 26

Polymerization

at 100°C

5.1 g. pro 1 cm³ 5.31

5.2 - - 92

7.8 - - 48

Theorie

Volle ——— Hermann's Polymer

↑ p₁ + p₂ + ... + p_n = 1
er eff. + s = 0 -

Donillet's experiment with various

of polymerization conditions, in which
p₁ + p₂ + ... + p_n = 1

Hermann's Polymerization of styrene = Influence
of the temperature

of the reaction - which is a polymerization

$$\frac{7}{2} V = \frac{Q}{D} \quad D = R + h$$

$$V = \frac{Q}{R+h} \quad \frac{\partial V}{\partial h} = -\frac{Q}{(R+h)^2} = -\frac{QR^2}{R^2(R+h)^2}$$

$$V_0 = \frac{Q}{R} = -\frac{V_0}{R} \left(\frac{R^2}{R+h} \right)$$

$$\frac{\partial V}{\partial h} = -\frac{V_0}{R} \quad \text{when } h \text{ is small}$$

is small

is small

$$R = 200 \text{ Volt per m} = -V_0$$

$$V_0 = 200 \cdot 64000000$$

$$= -130,000,000 \text{ Volt}$$

also it is the same as the potential of the wire

Redirection of the Thomson

for the case of the Thomson

of Thomson

$$V = -\frac{Q}{R+h} + \frac{Q}{R_2} \left(\frac{R^2}{R+h} \right)$$

$\frac{\partial V}{\partial h}$ will be zero when $h = R$

Then the potential of the wire

D. Edmund ~~was~~ $u = \frac{1}{2} \dots$

$u = \frac{1}{2} \dots$

$u = \frac{1}{2} \dots$ [Zurück]

$u = \frac{1}{2} \dots$

$u = \frac{1}{2} \dots$

The unipolar Induction

D. Edmund $u = \frac{1}{2} \dots$

$u = \frac{1}{2} \dots$ Volt

Gener, Induktion $u = \frac{1}{2} \dots$

Sohn

4. Simplex $u = \frac{1}{2} \dots$ F. Gener

$u = \frac{1}{2} \dots$ -- of 2

First $u = \frac{1}{2} \dots$

$u = \frac{1}{2} \dots$ -- of 2

$u = \frac{1}{2} \dots$ -- of 2.

Induktion $u = \frac{1}{2} \dots$ Hypothese

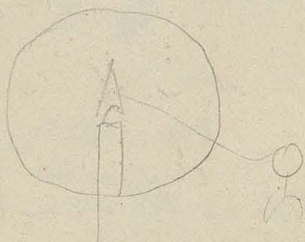
$u = \frac{1}{2} \dots$ Induktion; $u = \frac{1}{2} \dots$

o 1/20 of a lb - 100 g

Descant:

1/20 of a lb of a lb - 100 g

Descant = 1/20 of a lb of a lb - 100 g
inclusion by 1/20



Descant per 1/20 of a lb
100 g

Descant on 1/20 of a lb

1/20 of a lb of a lb - 100 g

1/20 of a lb of a lb - 100 g

Descant per 1/20 of a lb of a lb - 100 g

1/20 of a lb of a lb - 100 g

1/20 of a lb of a lb - 100 g

1/20 of a lb of a lb - 100 g

1/20 of a lb of a lb - 100 g

EVPS	100g	V per lb	814	106
2.3	325		85	97
3.8	297		10.4	89
4.4	197		11.9	74
5.5	166		12.5	68
6.8	116			

$\frac{\partial V}{\partial h} = -4\pi\sigma$... conductors

... $\frac{\partial V}{\partial h} = -4\pi\sigma_0$

... $b_0 - b'$... $\frac{\partial V}{\partial h}$

... $b_0 - b'$... $\frac{\partial V}{\partial h}$

... $b' = \epsilon \frac{\partial V}{\partial h}$... $\frac{\partial V}{\partial h}$

$\frac{\partial V}{\partial h} = -4\pi\sigma_0 + 4\pi k \epsilon \frac{\partial V}{\partial h}$

$\frac{\partial V}{\partial h} [1 - 4\pi k \epsilon] = -4\pi\sigma_0$

$\frac{\partial V}{\partial h} = -\frac{4\pi\sigma_0}{1 - 4\pi k \epsilon}$

... $\frac{\partial V}{\partial h}$...

... $\frac{\partial V}{\partial h}$...

... $\frac{\partial V}{\partial h} = \frac{\frac{\partial V}{\partial h}}{1 + 4\pi k \epsilon}$

... $\frac{\partial V}{\partial h}$...

... $b' = -$... $\frac{\partial V}{\partial h}$...

... $\frac{\partial V}{\partial h}$...

28/5

$$k = 1.31$$

$$\frac{\partial V}{\partial z} = \frac{A}{1 + 4\rho z}$$

$$A = 1330$$

$$\left(\frac{\partial V}{\partial z}\right)_0 = 1330$$

$$\frac{\partial V}{\partial z} = -426 = \frac{E}{2z} = -\frac{V_0}{z}$$

$$V_0 = -8 \cdot 10^9 \text{ Volt} \quad \text{mit } \epsilon \text{ in } E, \text{ Pol. } \times \times \times \times$$

4 Elektro's Gertel, Länge 100

mit je 1.6 mm Coeff 4 p 6 V pro m 420

10.4 "

pp.

89

60 da Länge + Anzahl Wd:

$$A = 1550 \text{ V.}$$

$$k = 1.31$$

A = C in Coeff 4 p 6 V pro m 420

200 p, A 2 p ~ November:

2 Nov - November

$$A = 1170$$

den 1930

Jänner 2640

Febr-Apr 1180

1/1, kann C in Coeff 4 p 6 V pro m 420

Exner et al - v. 2. 176th Edition

9

2668

$$A = 1420$$

$$k = 1.5$$

copy	Vollstrom	herkunft
17 m	432	441
3'4	323	276
6'3	191	168
9'3	112	120
12'4	82	93
15'4	76	76
17'5	66	69
19'5	60	62
21'4	56	57
23'5	51	52

empirical value for σ
on 26
 $V_0 \approx -9.109 V$

Theorie von Schenke

$f \cdot \rho = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... [Faraday]

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$f \cdot \rho = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

Schenke ... $\rho \cdot f = \rho \cdot f$

... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

Der ... $\rho \cdot f = \rho \cdot f$

Die Theorie ... $\rho \cdot f = \rho \cdot f$

Theorie von Arhenius

$f \cdot \rho = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

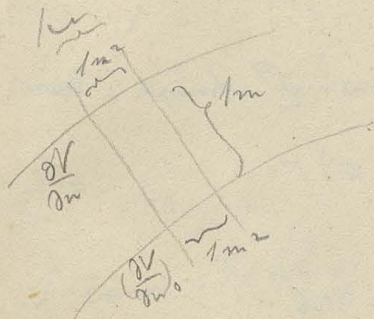
$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

$\rho \cdot f = \rho \cdot f$... $\rho \cdot f = \rho \cdot f$

What is the ... $\Delta x = -\text{neg. of } \Delta t$ 10
 ... $\Delta x = \Delta t$ if ...

for us ... then



$$\frac{\partial V}{\partial x} - \left(\frac{\partial V}{\partial x}\right)_0 = -400m$$

for ...
 ...

as ...

as ...

$$\frac{\partial V}{\partial x} = \left(\frac{\partial V}{\partial x}\right)_0 = -400m$$

with a ...
 $0m = \rho h$

$$\left(\frac{\partial V}{\partial x}\right)_0 = -400$$

↑ Neg of ρ

$$\frac{\partial V}{\partial x} = -400 - 400\rho h$$

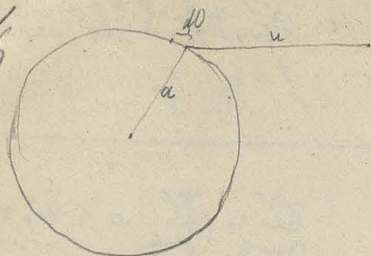
cost of ... ρh

... E

...

...

4/6



$$U = \int \delta \frac{d\theta}{u}$$

1) U_e 1) U_o
obtain U_i

$$V = U + W \quad W = \text{Pot. en. zw. } r \text{ u. } \sqrt{e} \text{ Stromph. - Strom.}$$

A) mit Hilfe d. Statik

$$\frac{U_e}{a} + 2 \left(\frac{\partial U_i}{\partial r} \right)_{r=a} = 4\pi\delta$$

1) \vec{c} um r konst. aus \vec{g} p

$$U_i + W_i = \text{const} = k$$

$$U_e + W_e = k$$

$$U_o + W_o = k$$

$$U_e = W_o + U_o - W_e$$

$$\frac{\partial U_i}{\partial r} r = - \frac{\partial W_i}{\partial r}$$

$$\frac{W_o - W_e + U_o}{a} + 2 \left(\frac{\partial W_i}{\partial r} \right)_{r=a} = 4\pi\delta$$

$$U_o = \frac{1}{2} \int \delta d\theta = \frac{E}{a}$$

$$4\pi\delta = \frac{E}{a} + \frac{W_o - W_e}{a} - 2 \frac{\partial W_e}{\partial r}$$

$$\frac{dV}{dt} = -\frac{E}{a} + \frac{W_0 - W_1}{a} + 2 \cdot \frac{dW_0}{dt} - 4n \rho d h \quad 11$$

$W = 2 \rho g C_1 \cdot 2 \rho g \cdot 5 \cdot 10^4$ Fragestellung

$e = 10^{-10} \text{ m}^2 = e \cdot f$

$$\frac{dV}{dt} = C_{\text{eff}} \cdot h - 4n \cdot \rho \cdot h \cdot a \cdot g / a$$

Δ Exner, 9e. Die Höhe von 60 m unter W_1 & W_2 ist

0 e 4 mph ρg \sim $n = \rho g$ e. Die W_1 & W_2 sind

f. Die W_1 & W_2 sind $n = \rho g$ e. Die W_1 & W_2 sind

Lecher:

57	93	1 m	Wasser
440 m	193		
550 m	193		
660 m	193		

Time

410 m	370 V	1300	554
500	406	1900	647
750	434		
820	480		
1000	490		
1120	508		

ρg \sim $6 \cdot 10^4$
 $\sim 4 \text{ mph} + \sim 1 \text{ m}^2$
 $-- 3 \text{ m}^2$

Les 277 1500 m

Sp 0.18 V. pro km. = 60m - für 11 Vill.

→ Paver 1000 m 400 = 0.7 für 1000 Paver [5%]

Passch 21 2 Feb - 1/2 m 1000 m 1000 m
je 20 < 1000 m 1000 m 1000 m

1000

Dionsten 1/2 m < 1000 m

Andale je 1000 m

→ 1000 m

Elster, Bettel

Peter Lehner am Brunnblick

6h 10h 12h 2h 4h 6h 8h 10h

→ 1/2 m 1000 m 1000 m 1000 m

1000 m 1000 m 1000 m 1000 m

Jahr Gang

	1. Okt.	11.	8.	7.	Fr	11.	11.	11.	11.
Walf	0.85	1.17	2.12	1.77	1.53	1.22	0.62	0.50	0.46
Sommt	0.87	0.84	1.05	0.98	1.13	1.00	1.16	1.07	1.01

2. $\rho = \rho_0$ in ρ constant
 $\rho = \rho_0 + \rho_1 \cos(kz - \omega t)$

$\frac{\partial v}{\partial z} = \left(\frac{\partial v}{\partial z}\right)_0 - \cos \rho h$
constant \uparrow

$2 \cos \left(\frac{\partial v}{\partial z}\right)_0$

$\rho = \rho_0 + \rho_1 \cos(kz - \omega t)$

$2 \cos \left(\frac{\partial v}{\partial z}\right)_0 \rho + \omega$

$h=0$

$\left(\frac{\partial v}{\partial z}\right)_{h=0} = -\frac{E}{2v} + \frac{W_0 - W_1}{e} + 2 \frac{\partial W_0}{\partial z}$

$\# \sim E - \text{just for } t_2$

$2 \rho \sim \text{interference pattern } \rho \sim \cos(kz)$

$< 2 \rho \sim \text{interference pattern } \rho \sim \cos(kz)$

$\rho \sim \text{interference pattern } \rho \sim \cos(kz)$

$\rho \sim \text{interference pattern } \rho \sim \cos(kz)$

$\rho \sim \text{interference pattern } \rho \sim \cos(kz)$

$$\frac{1}{6} \left(\frac{\partial V}{\partial h} \right)_{h_0} = - \frac{E}{A^2} + \frac{W_0 - W_0}{A} + 2 \left(\frac{\partial W}{\partial L} \right)_{L=A}$$

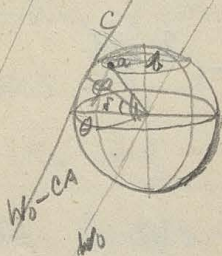
~~XXXXXXXXXXXX~~

$p_{avg} \sim \frac{1}{2} \rho g h_0 E$ or $I C$

Π of p_{avg} or \sim NSS of p

\sim ρ in 20 m cut back; \sim ρ of h_0 Δh

\sim $2 \rho t$ \sim ρ in 20 m cut back \sim ρ in m



$\sim \{ \varphi, \theta \}$

$W_0 - CA + c p$ \leftarrow Proj. of $c p$ \leftarrow $\rho g h_0$

$$p = bc - \text{Proj. of } bc$$

$$ab = \text{emp}$$

$$bc = \text{emp} \cos \theta$$

$$\text{Proj. of } bc = \text{emp} \cos^2 \theta$$

$$p = \text{emp} \sin^2 \theta$$

$$bc = a - \text{emp} \cos^2 \theta$$

$$p = a - \text{emp} \cos^2 \theta - \text{emp} \cos^2 \theta$$

$$W_0 - W_0 - c a [\sin \alpha - \delta + \sin \alpha \cos \alpha]$$

B

$a < \sin \alpha < 0$

$$\frac{\sin \alpha - \delta}{\sin \alpha} = \frac{\sin \alpha - \delta}{\sin \alpha}$$

$\sin \alpha > \delta$ so $\sin \alpha - \delta > 0$

$$W_0 = W_0 - c a [\dots]$$

$$\frac{W_0 - W_0}{a} = -c [\sin \alpha - \delta]$$

$$\frac{\partial W}{\partial a} da = -c [\sin \alpha - \delta]$$

$$\frac{W_0 - W_0}{a} + 2 \left(\frac{\partial W}{\partial a} \right)_a = -3c [\sin \alpha - \delta + \sin \alpha \cos \alpha]$$

$$\left(\frac{\partial W}{\partial a} \right)_a = -\frac{E}{A^2} - 3c [\dots]$$

$\sin \alpha > \delta$ so $\sin \alpha - \delta > 0$

$\sin \alpha > \delta$ so $\theta = 0$

$$\left(\frac{\partial W}{\partial a} \right)_a = -\frac{E}{A^2} - 3c \sin(\alpha - \delta)$$

$\sin \alpha > \delta$ so $\sin(\alpha - \delta) > 0$

$\sin \alpha > \delta$

$a < c > 0$ $\sin \alpha > \delta$ $\sin(\alpha - \delta) > 0$ $\left(\frac{\partial W}{\partial a} \right)_a < 0$

of the ...

... ..

for ... $-3c[\cos \theta]$
... $\phi = 0$... $\delta = 0$

$$\therefore -3c[\cos \theta]$$

of ... $-3c$

MV $+3c$

... ..

... ..

of

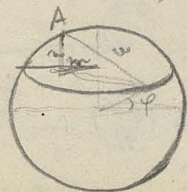
... ..
... ..

... ..

$$0 = \frac{W_0 - W_0}{r_0} + 2 \frac{dW}{r} = 0$$

... ..

Why



... ..

$W_0 = \text{rad. g. } \sim c A$

$$W = \int \frac{e m \varphi d\theta}{u}$$

$$u = \sqrt{m^2 + m^2 + a^2 \sin^2 \theta - 2am \cos \theta}$$

$$\int \frac{6a e m \varphi d\theta}{\sqrt{m^2 + m^2 + a^2 \sin^2 \theta} \sqrt{1 + k \cos \theta}}$$

$$\left\| \begin{aligned} p &= -\frac{2am \cos \theta}{m^2 + m^2 + a^2 \sin^2 \theta} \\ (p < 1) \end{aligned} \right.$$

$$W = \frac{6a e m \varphi}{\sqrt{m^2 + m^2 + a^2 \sin^2 \theta}} \int_0^{2\pi} \frac{d\theta}{\sqrt{1 + k \cos \theta}}$$

$$W = \frac{6a e m \varphi}{\sqrt{\dots}} \ln \left\{ 1 + k_1 p^2 + k_2 p^4 + \dots \right\}$$

$\sim \ln \left(\frac{1 + k_1 p^2 + k_2 p^4 + \dots}{1 - k_1 p^2 + k_2 p^4 + \dots} \right)$

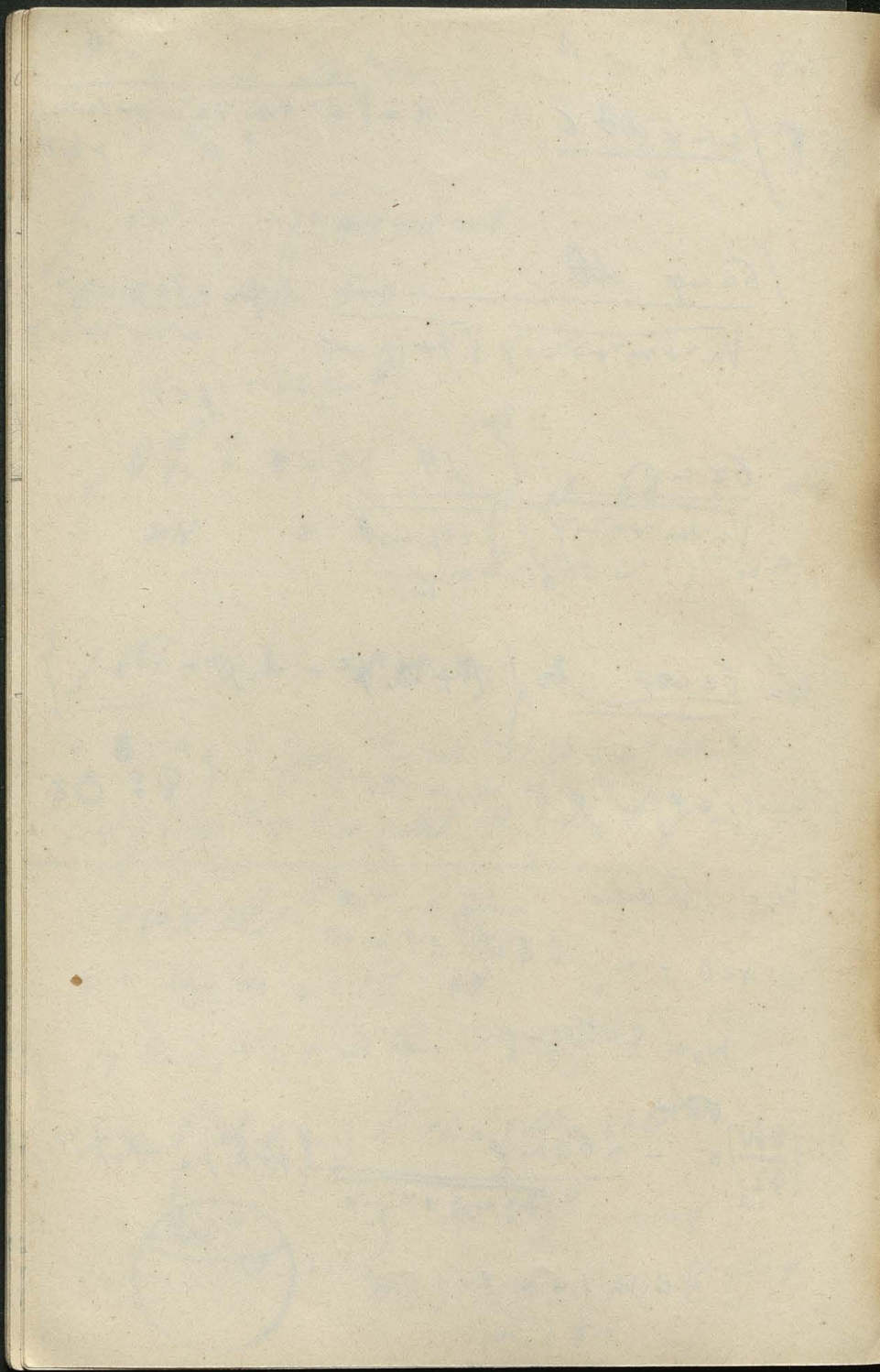
$W_0 = \text{Index}$

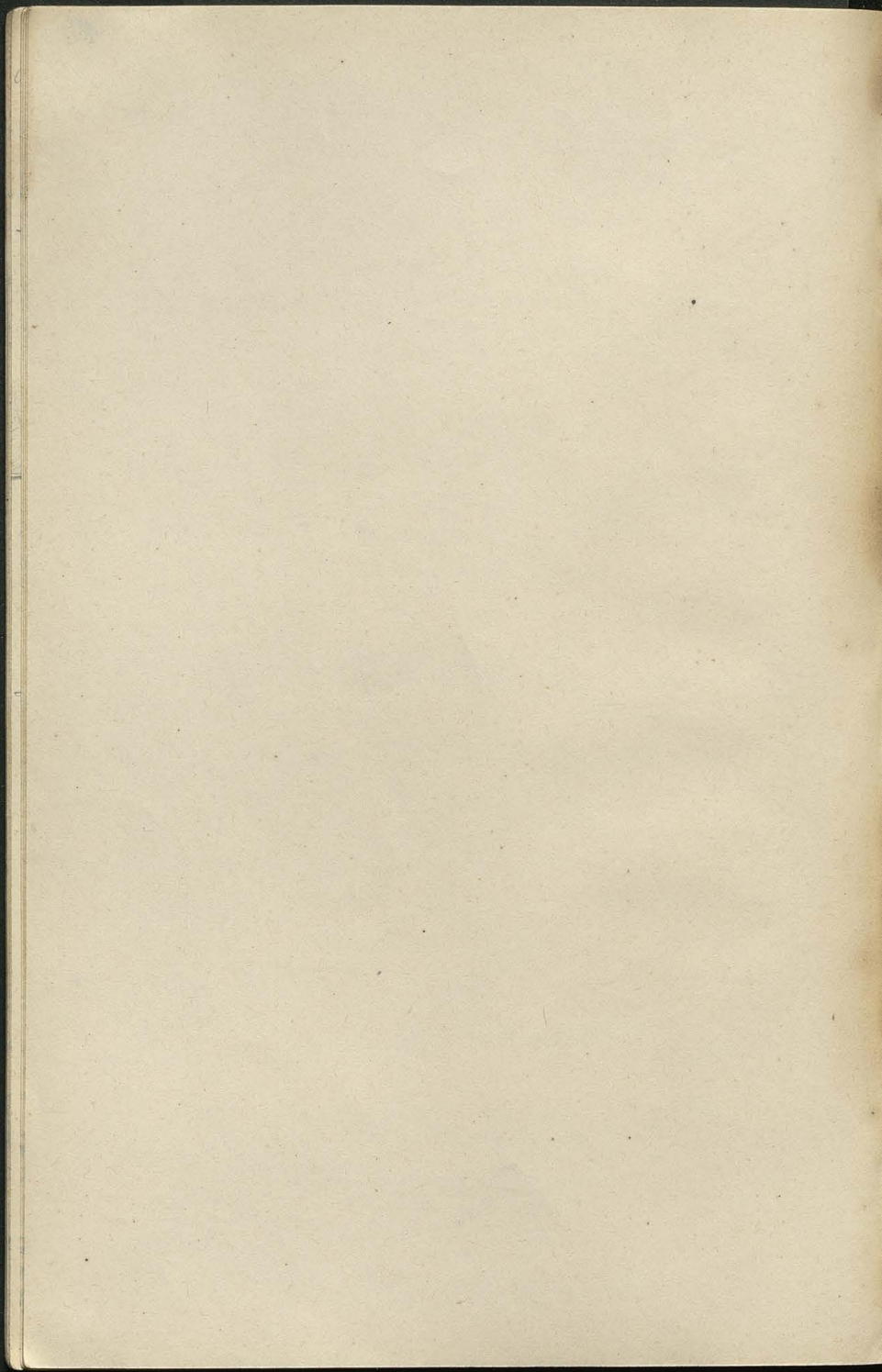
$$a r = 0 = a_c \left(\frac{1}{2} \delta \right) \cdot \delta + m = 0 \quad \text{mean}$$

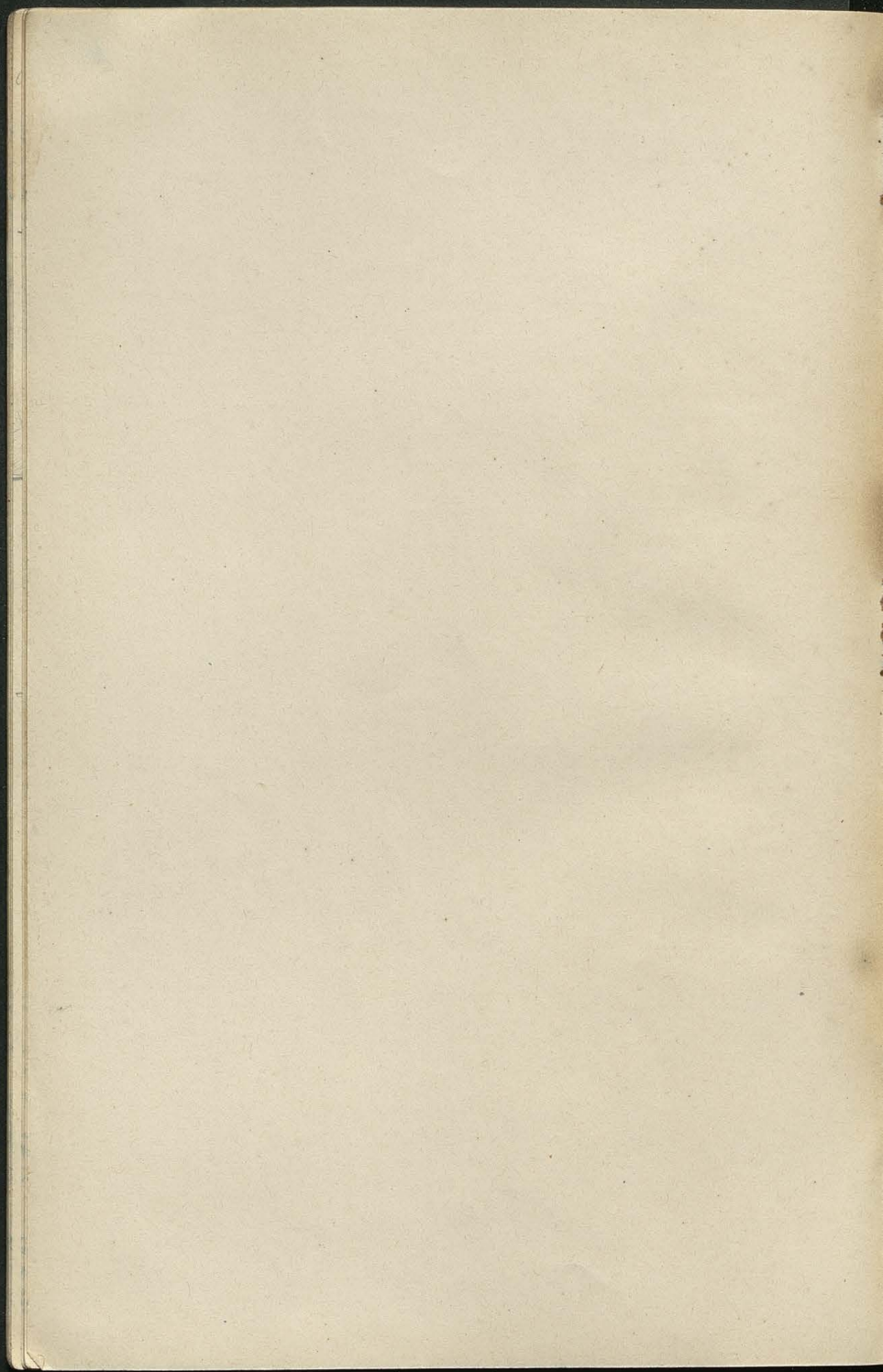
$$W_0 = \frac{2\pi 6a e m \varphi}{a}$$

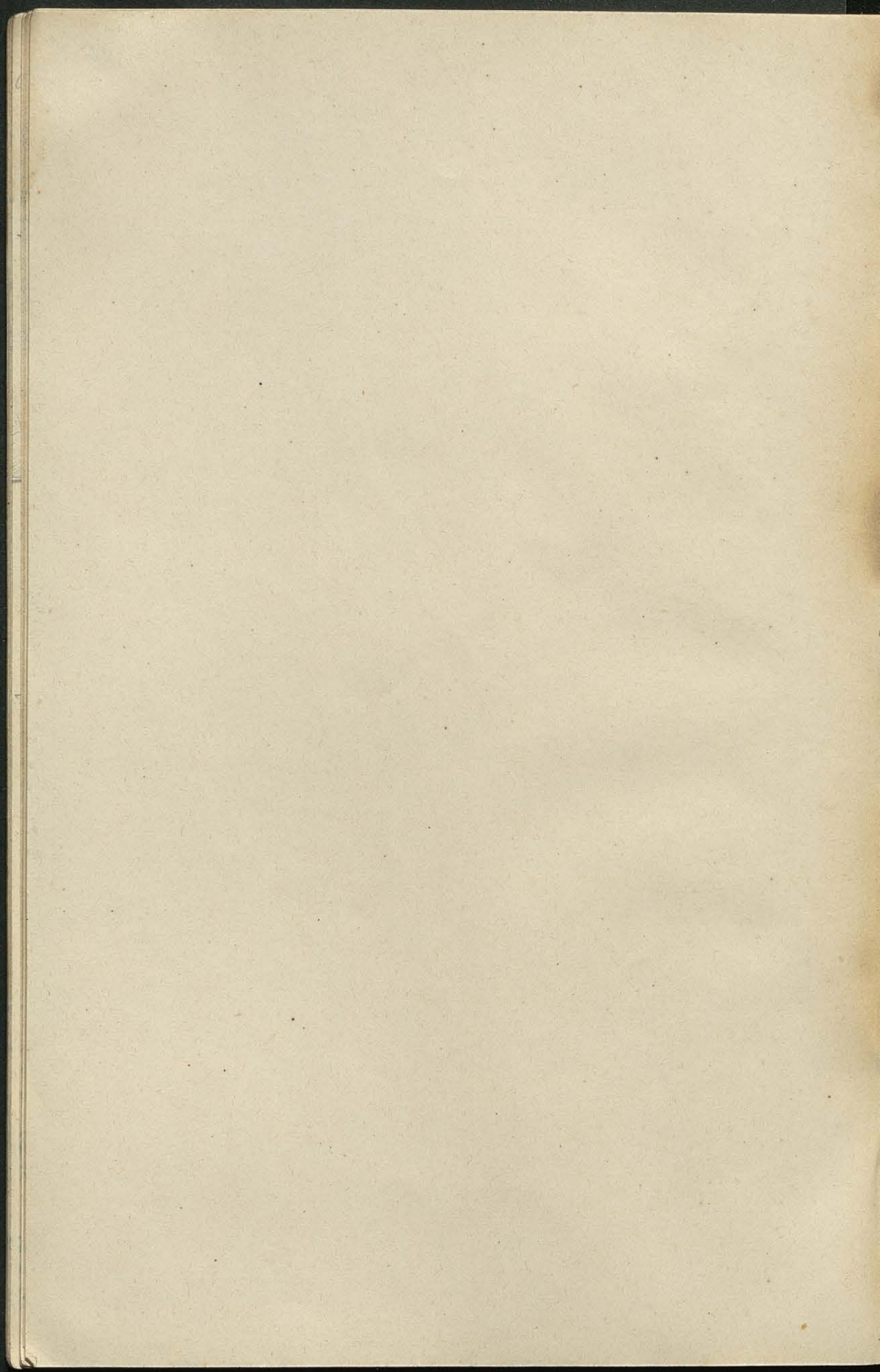
$$\left(\frac{\partial W}{\partial r} \right)_{r=0} = - \frac{2\pi 6a e m \varphi}{a \sqrt{m_0^2 + m_0^2 + a^2 \sin^2 \theta}} \left\{ 1 + k_1 p_0^2 + k_2 p_0^4 + \dots \right\}$$

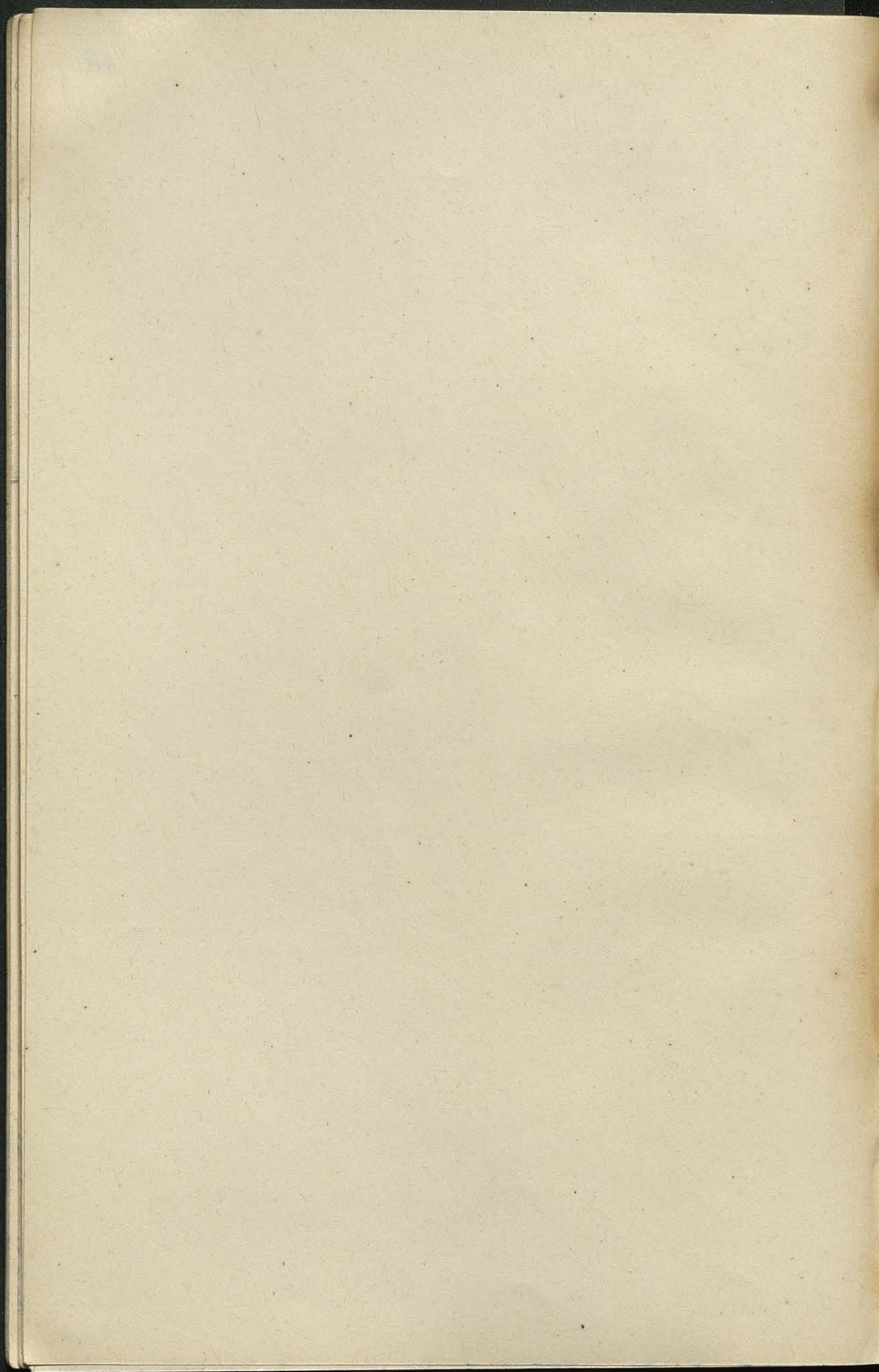
$\sim \ln \left(\frac{1 + k_1 p_0^2 + k_2 p_0^4 + \dots}{1 - k_1 p_0^2 + k_2 p_0^4 + \dots} \right)$

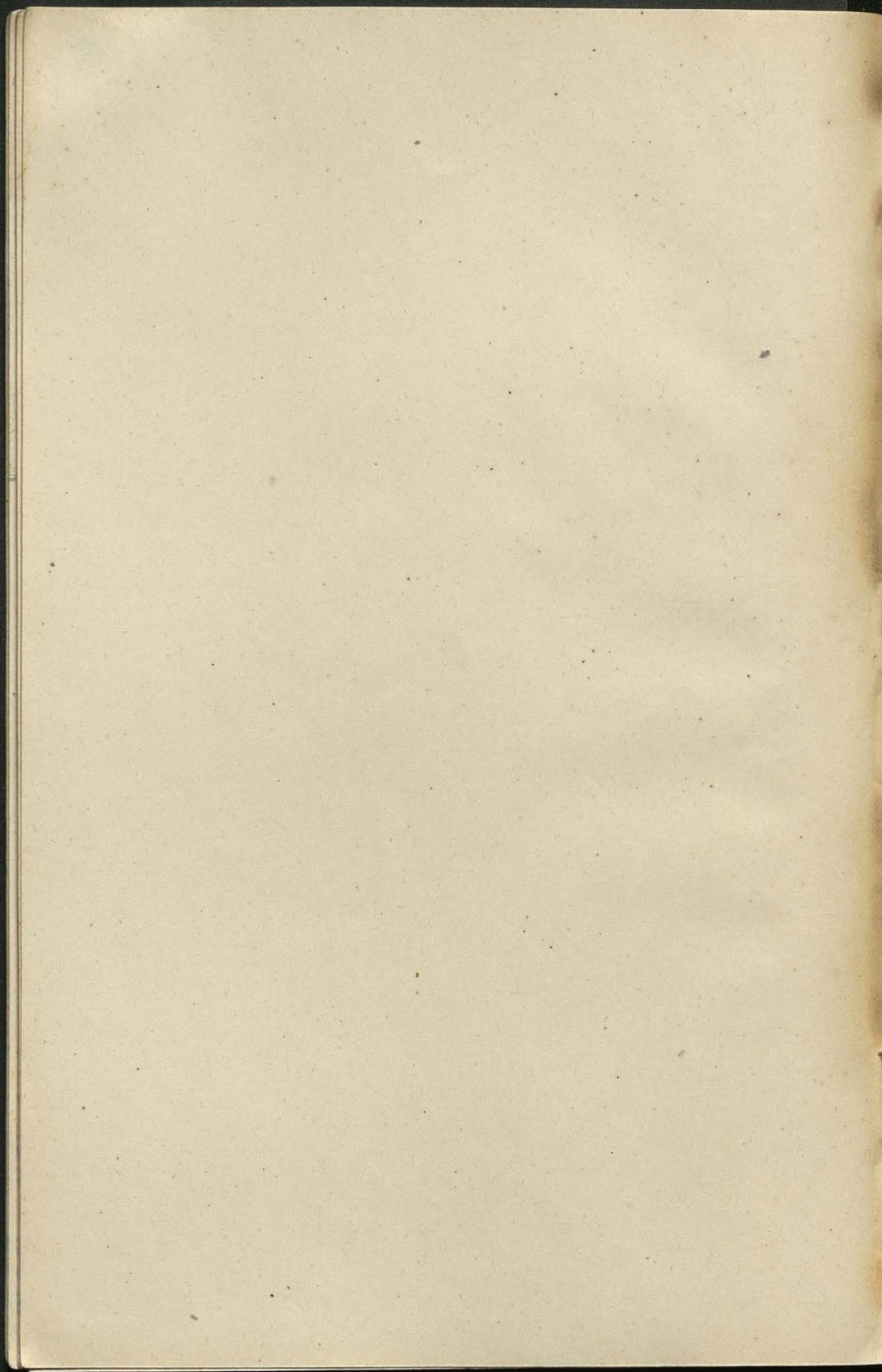


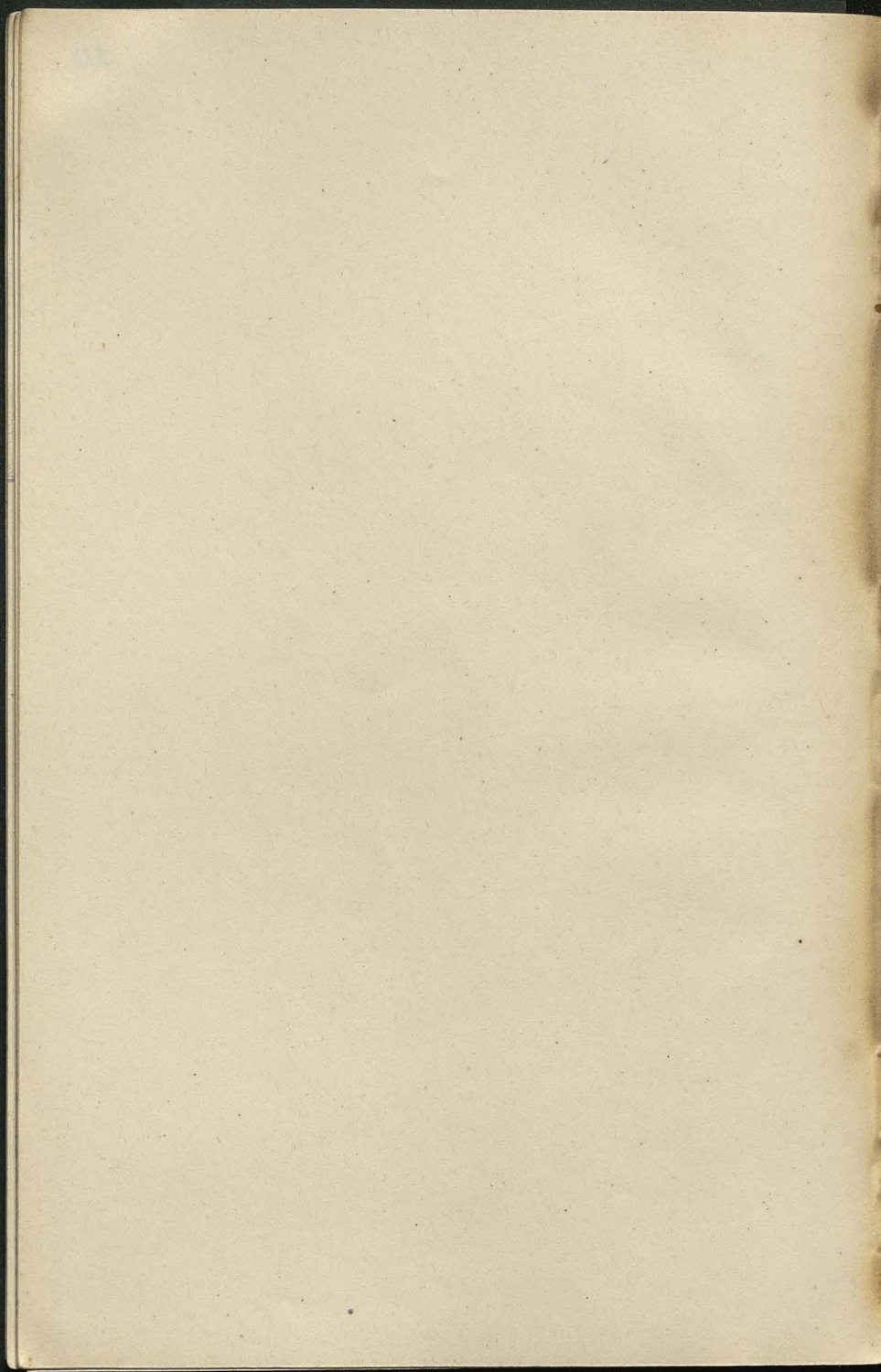


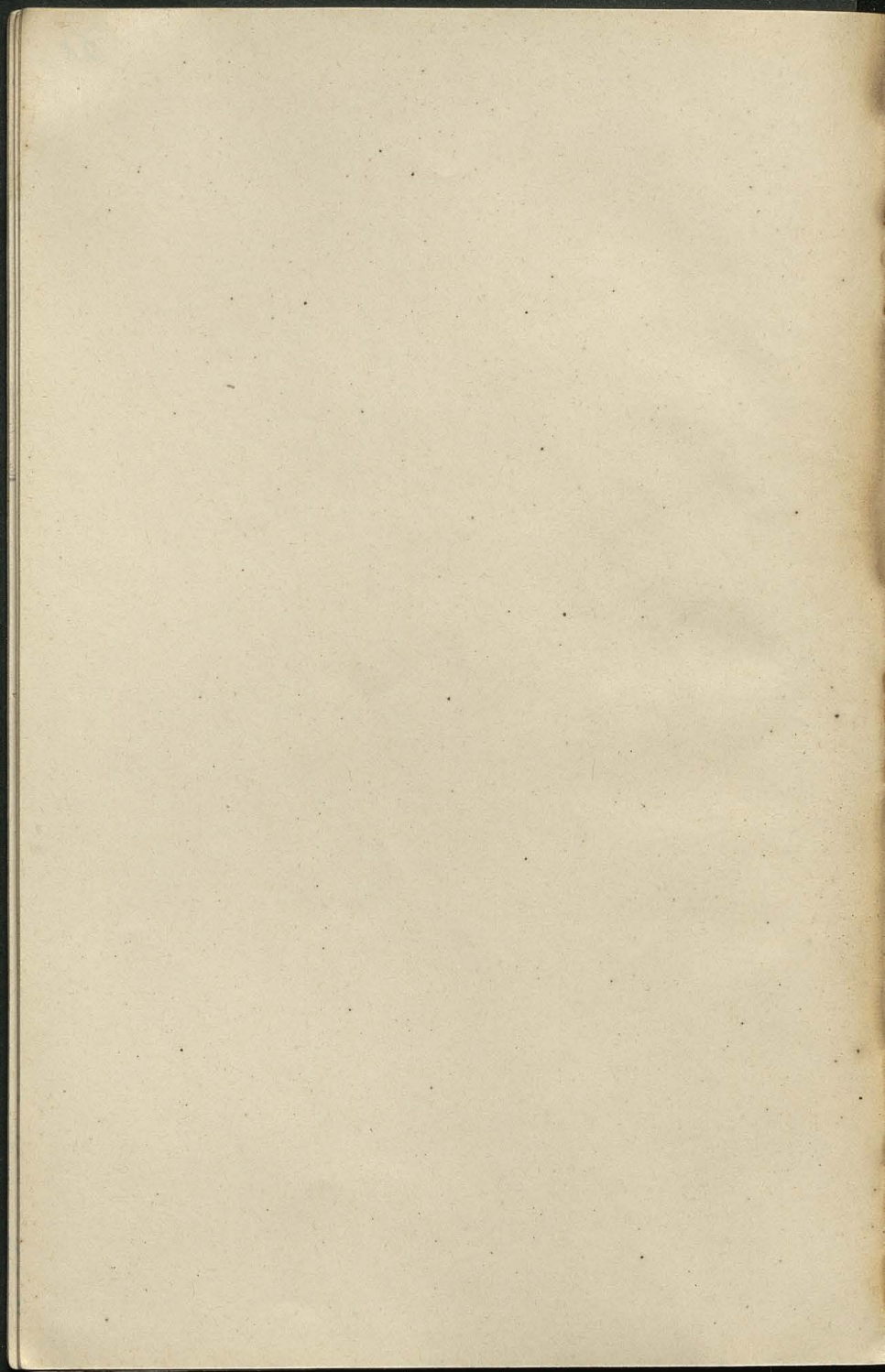


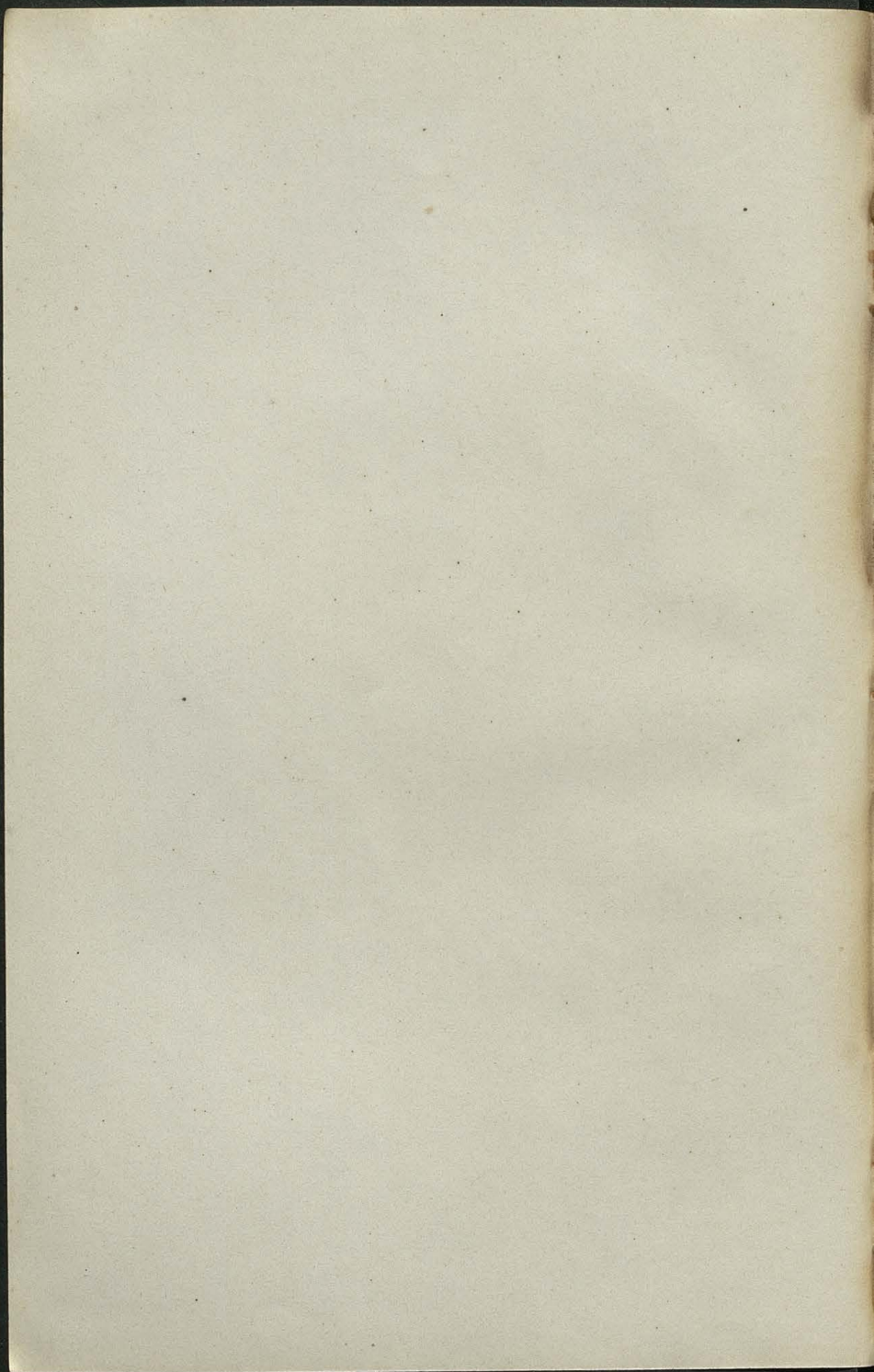


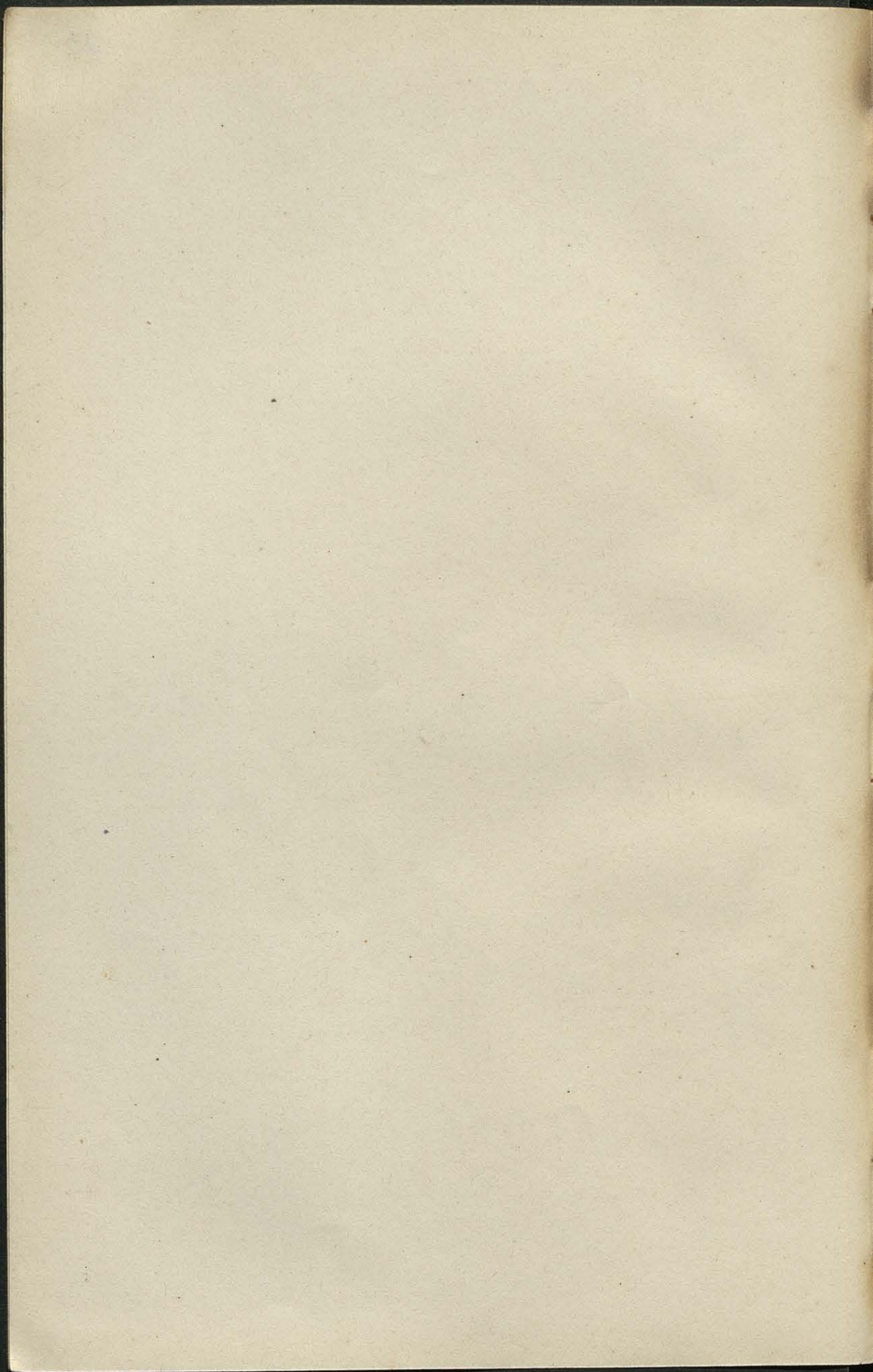


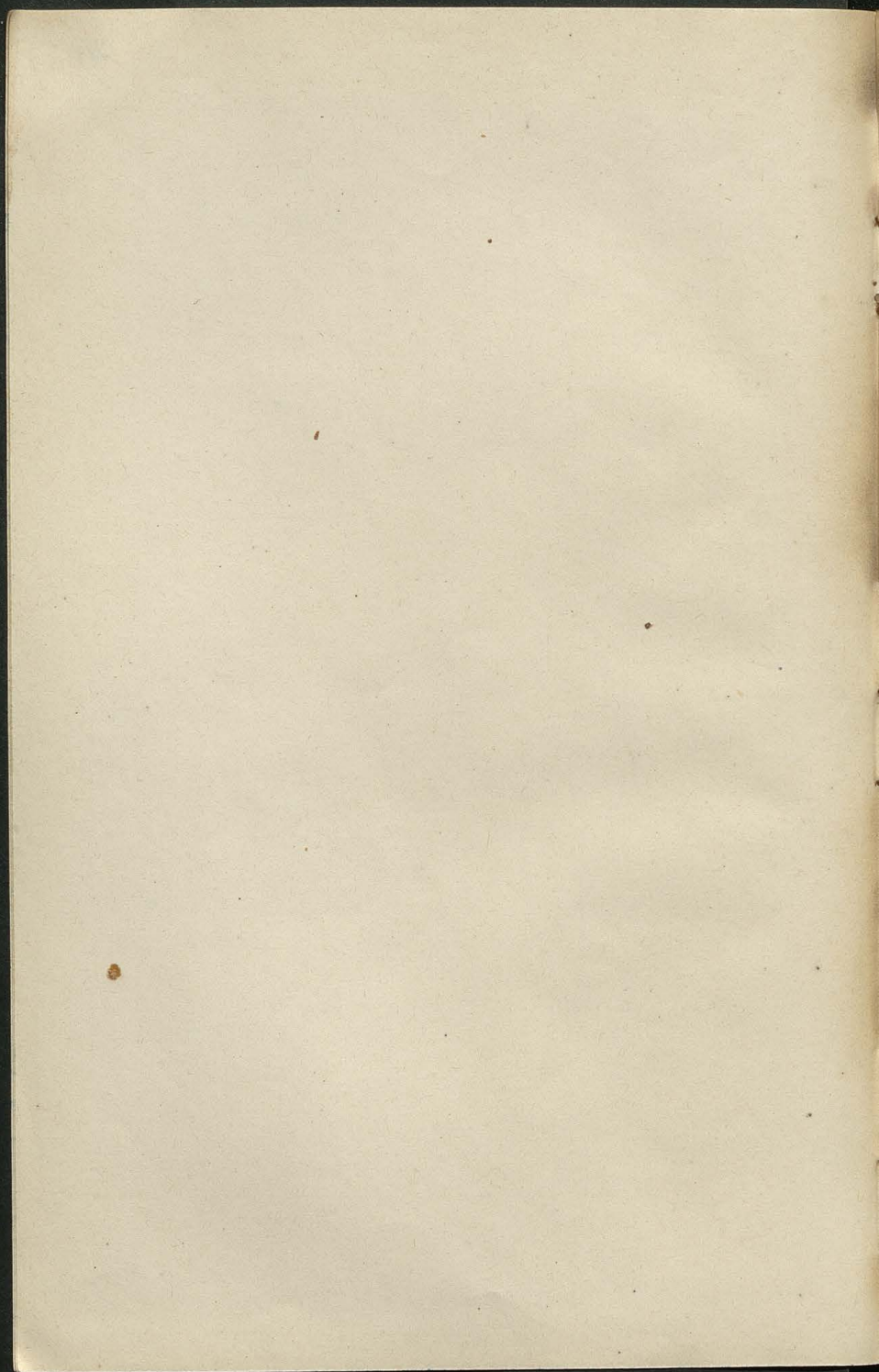


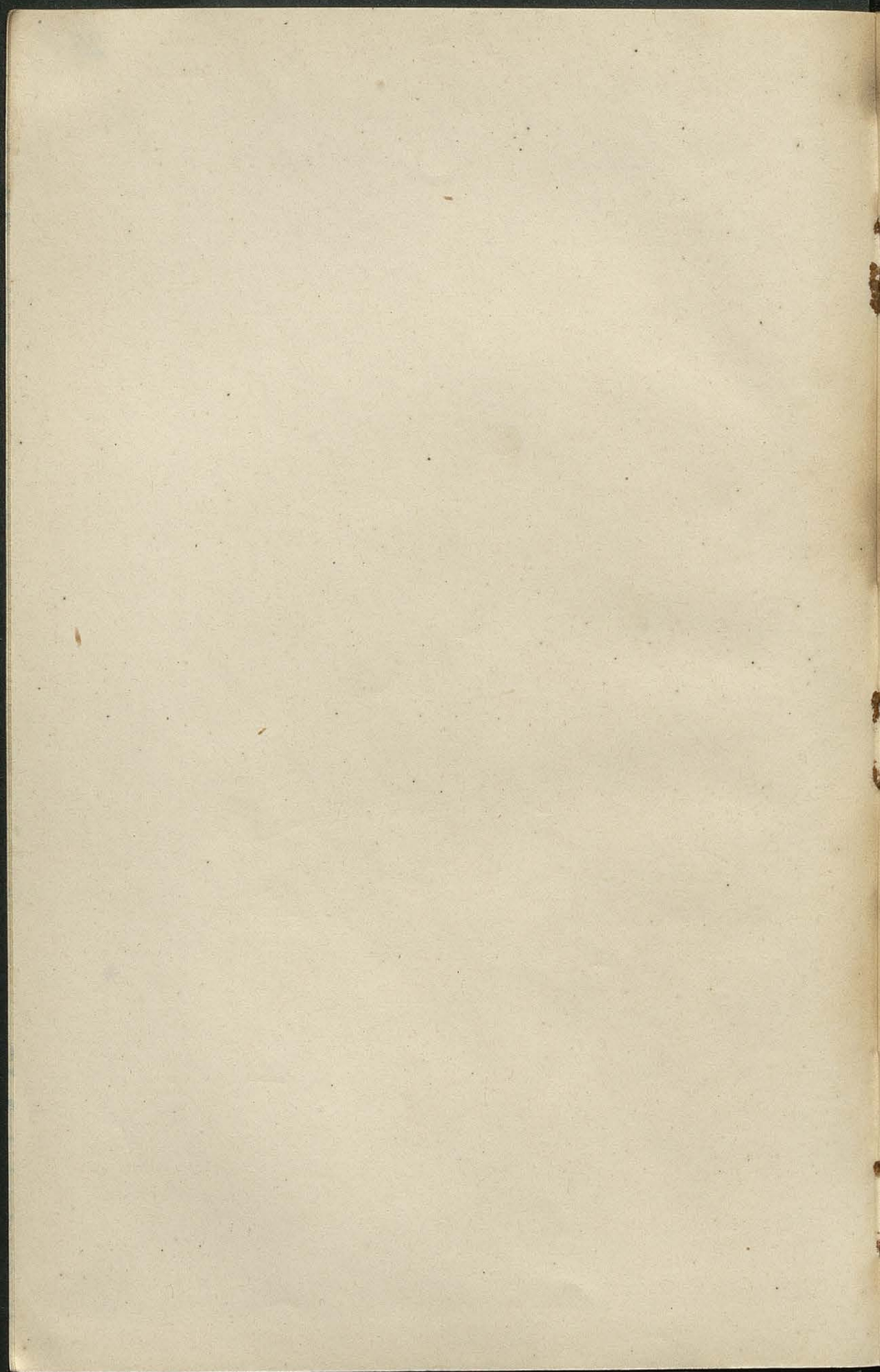












$$\varphi(u+u) = \varphi(u)$$

$$\varphi(u+u') = e^{-\frac{u'}{u} u} \varphi(u)$$

при $u=0$

$$u' = 2u - u = u$$

$$u = u' - u$$

$$\varphi(u'+u) = \varphi(u'-u)$$

$$\varphi(u+u') = e^{-\frac{u'}{u}(u-u)} \varphi(u-u')$$

$$-\varphi(u-u) = -\varphi(u'-u)$$

$$\varphi(u+u') = \varphi(u-u) = 0$$

и т.д.

и т.д.

и т.д.

$\varphi(u) = e^{-\frac{u}{a}}$

$$-\varphi(u) = e^{-\frac{u}{a}} \quad \text{Answer } + C$$

$$C(u-a)$$

$$\varphi(w) = \sum_{n=0}^{k-1} \left[a_n e^{\frac{i\pi}{w}(n-\beta)} + \frac{i\pi}{w} \theta_{00} \frac{2=(n-\beta)k\pi}{2w} (n-\theta) \frac{k\pi}{2w} k\pi \right]$$

$$\sigma = \log \rho = \frac{u}{w}$$

$$r = r'$$

$\varphi(w) \sim \text{log of } \theta_{00} \text{ of } \theta$

n. $k=1$, $n=0$, ~~$k=2$~~

~~$$\tau + \log \rho = 0$$~~

$$\rho = e^{-\tau} = \frac{1}{p} \quad \beta=0$$

$$z = \frac{u i \pi}{2w}$$

~~$$\varphi\left(\frac{2u z}{i\pi}\right) = a_0 e^{2z + \frac{i\pi}{w}} \theta_{00}(z, \tau)$$~~

1 or θ_k

~~$$f(z) = \dots$$~~

$$z_0 \sum \varphi e^{2z} \# e^{p 2z}$$

$$= a_0 \sum_{p=0}^{\infty} e^{p^2 \tau + 2p z}$$

$$\varphi(w) = a_0 \sum_{p=0}^{\infty} e^{p^2 \tau + \frac{i\pi u p}{w}}$$

$$= 1 + \sum_{p=1}^{\infty} e^{p^2 \tau} \left[e^{\frac{i\pi u p}{w}} + e^{-\frac{i\pi u p}{w}} \right]$$

$2 \cos \frac{\pi u p}{w}$

$$e^{i\pi u - p^2 \tau} \theta_{00} \sim \varphi(-u) = \varphi(u)$$

$$\mu^p a_{n-p} = \rho^{2p} + \rho^{p-1} \kappa$$

$$p(n) = \sum_{n=0}^{k-1} \left[a_n \sum_{-\infty}^{+\infty} \frac{1}{\mu^p} \rho^{2p} e^{(2p+1) \frac{\rho n}{\omega}} \right]$$

$$\psi(n) = \sum_{-\infty}^{+\infty} \frac{1}{\mu^p} e^{2p n \tau + \rho^2 k - \rho k \tau + (\rho + \nu) \frac{\rho n}{\omega}}$$

$$u = \alpha v + \beta$$

$e^{\rho \sigma} \delta \sigma \dots$

$$\psi(r, \alpha v + \beta) = \sum_{-\infty}^{+\infty} e^{2p r \tau + \rho^2 k \tau - \rho k \tau +}$$

$$+ (\alpha k \rho + \alpha r) \frac{i \nu v}{\omega} + \frac{i \nu}{\omega} (k \rho \beta + \nu \rho) - \rho \log \mu$$

$$\cancel{2p r \tau - \rho k \tau +}$$

$$2 \rho r \tau + \cancel{\rho^2 k \tau} - k \tau + \frac{i \nu v \alpha k}{\omega} - \frac{i \nu k \rho}{\omega} - \log \mu$$

$$\psi(r, \alpha v + \beta) = \sum_{-\infty}^{+\infty} e^{2p^2 k \tau + \alpha k \rho \frac{i \nu v}{\omega}} \cdot e^{\alpha \rho \frac{i \nu v}{\omega} + \frac{i \nu \rho}{\omega} \rho}$$

$$= e^{\frac{\alpha r i \nu v}{\omega} + \frac{i \nu}{\omega} \nu \rho} \sum_{-\infty}^{+\infty} e^{\rho^2 \tau + 2p \tau}$$

$$= \cancel{\psi(r)}$$

$$\theta_{00} (2r, \tau')$$

10. 2=1

$$\mu^p a_{pk} = a_0 \rho \quad k(p-0)\rho$$

$$\downarrow$$
$$\mu^p a_{pk} = a_{-k} \rho \quad m=-k \quad 2k$$

$$\mu a_0 = a_{-k} \rho$$

$$\mu e_{-k} = a_{-k} \rho \quad -4k$$

$$\mu e_{-pk} = a_{-pk} \rho \quad -2pk$$

$$\mu^p a_0 = a_{-pk} \rho \quad k \rho (0^+)$$

$$\begin{aligned} \sum a_m e^{im\omega} &= \mu e^{\frac{ikir}{\omega}} \sum a_m e^{\frac{m\omega}{\omega}} \quad 28 \\ &= \mu \sum a_m e^{im\omega} \\ &= \mu \sum a_m e^{(m-k)\omega} \\ &= \mu \sum a_m e^{\frac{m\omega}{\omega} + \frac{-k\omega}{\omega}} \end{aligned}$$

$$e^{\frac{+ikir}{\omega}} = \rho$$

$$\begin{aligned} &= \rho \sum a_m e^{im\omega} \\ &= \sum \rho a_{m+k} e^{im\omega} \end{aligned}$$

$$\rho a_{m+k} = a_m \rho^{2m}$$

and so on for $k=1, 2, \dots$
 ρ^{2k}

~~ρ^{2k}~~

$$\rho^{2k} a_{m+k} = a_m \rho^{2m} \quad m=0$$

$$\rho^{2k} a_{2k} = a_0 \rho^{2k} \quad a_k = \frac{a_0}{\rho^k}$$

$$\rho^{4k} a_{4k} = a_0 \rho^{4k}$$

$$\rho^{6k} a_{6k} = a_0 \rho^{6k}$$

$$\left. \begin{aligned} \rho^{2(p-1)k} a_{2(p-1)k} &= a_0 \rho^{2(p-1)k} \\ \rho^{2p k} a_{2p k} &= a_0 \rho^{2p k} \end{aligned} \right\} [1+2+\dots+p-1]$$

$$q_{ev} = -2\gamma K$$

$$\omega R \frac{u}{c} = \frac{\pi a}{c}$$

$$a = -\frac{\gamma K}{2\omega}$$

$$\lambda = \frac{\lambda - \gamma' K}{2\omega l} = -\frac{\gamma K}{2\omega} \omega \omega l -$$

$$-\lambda l = K(\omega l \gamma K - \gamma' l) = \frac{\gamma K}{\omega} \frac{\pi a}{2}$$

$$\lambda = -\frac{K \pi a}{\omega l^2}$$

~~$\psi(x) = e^{\lambda x}$~~ $\psi(x) = 0$ velo $\frac{1}{2} K$
signe $\frac{1}{2} \omega$

$$\psi(x) = \psi(x) \quad \omega = K \frac{1}{2} \gamma \omega$$

$$\psi(x) = e^{2\lambda x} \psi(x)$$

$$\lambda = -\frac{K}{\omega} \frac{\pi a}{2}$$

~~$\psi(x) = \sum_{-\infty}^{+\infty} a_m e^{i \frac{m \pi x}{l}}$~~

$$\psi(x) = \sum_{-\infty}^{+\infty} a_m e^{i \frac{m \pi x}{l}} = \sum_{-\infty}^{+\infty} a_m \rho$$

$$\psi(x) = \sum_{-\infty}^{+\infty} a_m e^{i \frac{m \pi x}{l}} = \sum_{-\infty}^{+\infty} a_m \rho$$

$$e^{i \frac{\pi x}{l}} = \rho = \sum_{-\infty}^{+\infty} a_m \rho$$

$$= \rho \sum_{-\infty}^{+\infty} a_m \rho$$

$$= \rho \sum_{-\infty}^{+\infty} a_m \rho$$

$\psi(x)$

$$y(uv) = \sigma(u) \quad (u-\alpha)(-1)^k e^{2\eta'x(u+v)}$$

29

by method

$$-2\eta' \Sigma \alpha_n + P(u+2v)$$

$$= \frac{y(uv)}{\sigma(u)} e^{2\lambda u + e} \quad \frac{P(u)}{e}$$

$$2\eta' u(u+v) + 2\eta' \Sigma \alpha_n + P(u+2v) - P(u) + (2\rho'+1)u - 2\lambda u - h = 2m' u$$

$$\therefore \exists P'(u+2v) = P'(u)$$

$\therefore P'(u)$ is a degree 3 constant

$$\therefore P(u) = au^2 + bu + c$$

$\rho' + \beta \rightarrow a, b, c$

$$P(u+2v) - P(u) = 4auv' + 4u'^2 + 4bv'u + (2\rho'+1)u - 2\eta' u(u+v) + 2\eta' \Sigma \alpha_n + 2\lambda u + h$$

by method

$$4au' = -2\eta' u + 2\lambda$$

$$a = \frac{-\eta' u + \lambda}{2u'} \quad \text{and } b = \dots$$

$$= 4auv + 4u^2 + 4bv = (2\rho'+1)u - 2\eta' u(u+v) - 2\eta' \Sigma \alpha_n$$

$x \sim \rho \in \varphi$

$$\rho \varphi - \omega \ln [f \cdot v - \theta \rho]$$

$$e \varphi(u) \sim \dots$$

$$0 = k \theta \omega \sim \alpha_1 \alpha_2 \dots \alpha_k$$

$$= c \rho \theta \omega$$

$$\varphi(u) = \sigma(u-\alpha_1) \sigma(u-\alpha_2) \dots \sigma(u-\alpha_k) e^{\rho u}$$

$$\sim e \varphi : \dots$$

$$\dots = \dots$$

$$= e k$$

$$= \rho(u)$$

$$\varphi(u+2\omega) = \sigma(u-\alpha_1) \dots \sigma(u-\alpha_k) (-1)^k e^{-2\eta \sum \alpha_k} + \rho(u+2\omega)$$

$$= \varphi(u) = \sigma(u-\alpha_1) \dots e^{\rho u}$$

$$(-1)^k = e^{k i \pi} = e^{2\pi i n}$$

$$\rho''(u+2\omega) = \rho''(u) + 2\eta \sum \alpha_k + \rho(u+2\omega) - \rho(u) + 2\pi i n$$

$$\rho''(u+2\omega) = \rho''(u)$$

$$\varphi(u+2\omega) = \sigma u -$$

$$F(u+w) = \frac{1}{h} e^{a(u+w) + at} + b - au^2 - 4awt -$$

$$4aw^2 - 2fu - 4pw$$

$F(u)$

$$e^u \approx F(u+w) \approx F(u) \approx$$

$$f' : a + b - 4a^2u^2 - 4a^2w^2 - 4pw \approx \frac{1}{2} \sin u$$

$e \ln u \approx 0$

$$a - 4a^2w = 0$$

$$a = \frac{a}{4w}$$

$$b - 4a^2w^2 - 4pw = \frac{1}{2} \sin u$$

$e^u \approx f(u) + \frac{1}{2} \sin u; e^{u^2} \approx f(u^2)$

$$F(u+w) = F(u)$$

$$F(u+w) = e^{A(u+w) + B} F(u)$$

$$e^{A(u+w)} \approx e^A = 0$$

$$e^{A(u+w)} \approx f(u) + \frac{1}{2} \sin u$$

$$f(u+w) = e^{2g(u+w)} f(u)$$

$$g = \frac{f'(u)}{f(u)}$$

für $z = 2i\pi$

$$\varphi(z + 2i\pi) = \varphi(z)$$

$$\varphi(z + 2\pi) = e^{2z} \varphi(z)$$

$\sqrt{2} \varphi(z)$

$$\varphi(z + 2\pi) = \varphi(z)$$

$$\varphi(z + 2\pi) = e^{2z} \varphi(z)$$

$\sqrt{2} \varphi(z)$

$\varphi(z + 2\pi) = \varphi(z)$

$$\varphi(n\pi) = \varphi(\pi) \quad \text{für } n \text{ gerade}$$

$$\varphi(n\pi) = e^{2n} \varphi(\pi)$$

$\varphi(z + 2\pi) = \varphi(z)$

$$\varphi(n\pi) = A e^{2n} \varphi(\pi)$$

$$\varphi(n\pi) = A e^{2n} \varphi(\pi)$$

$$\varphi(n\pi) = \varphi(\pi)$$

$$\varphi(n\pi) = e^{2n} \varphi(\pi)$$

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$$\varphi(n\pi) = e^{2n} \varphi(\pi)$$

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$$\varphi(n\pi) = e^{2n} \varphi(\pi)$$

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$$\varphi(n\pi) = e^{2n} \varphi(\pi)$$

$$D_{0,1}(2+\frac{1}{2}i\pi) = \sum_{n=1}^{\infty} e^{n\tau + 2n\pi + n i \pi} \quad 31$$

$$= D_{0,0}(2)$$

$$D_{1,0}(2+\frac{1}{2}i\pi) = \sum_{n=1}^{\infty} e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}(2\pi + \pi)$$

$$\leq \sum_{n=1}^{\infty} e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}\tau + \frac{2n+1}{2}2\pi$$

$$= e^{-\frac{\tau}{4}-2} \sum_{n=1}^{\infty} e^{\left[\frac{(2n+1)}{2}n + \frac{1}{2}\right]\tau + \frac{2n+1}{2}2\pi}$$

$$+ \left[\frac{2n+1}{2} + 1\right]2\pi$$

$$= e^{-\frac{\tau}{4}-2} \sum_{n=1}^{\infty} e^{[n+1]\tau + n\pi} 2$$

$$= e^{-\frac{\tau}{4}-2} D_{0,0}(2)$$

$$D_{1,0}(2+\frac{1}{2}i\pi) = \sum_{n=1}^{\infty} e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}(2\pi + i\pi) =$$

$$= \sum_{n=1}^{\infty} e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}2\pi + \frac{2n+1}{2}i\pi$$

$$= e^{\frac{(2n+1)\tau}{2}} +$$

$$= \frac{1}{i} D_{1,1}(2)$$

$$D_{1,1}(2+\frac{1}{2}i\pi) = i \sum_{n=1}^{\infty} (-1)^n e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}2\pi + \frac{2n+1}{2}i\pi$$

$$= i e^{-\frac{\tau}{4}-2} \sum_{n=1}^{\infty} (-1)^n e^{(n+1)\tau + (n+1)2\pi}$$

$$= i e^{-\frac{\tau}{4}-2} D_{0,1}(2)$$

$$D_{1,1}(2+\frac{1}{2}i\pi) = i \sum_{n=1}^{\infty} (-1)^n e^{\frac{(2n+1)\tau}{2}} + \frac{2n+1}{2}2\pi + \frac{2n+1}{2}i\pi$$

$$= \frac{-D_{1,0}(2)}{i}$$

$$= A e^{2 + \frac{1}{2}ni + \frac{1}{2}\tau} (1 - e^{-2\tau}) \prod_{n=1}^{\infty} (1 - 2q^{-2n} \cos^2 \tau + q^{4n})$$

$$= A q^{\frac{1}{2}} 2 \sin(i\tau) \prod_{n=1}^{\infty} (1 - q^{2n} \cos^2 \tau + q^{4n})$$

alignement

$$V_{1,1}(z) = \sum_{-\infty}^{+\infty} e^{\frac{(2n+\frac{1}{2})^2}{2} \tau + (2n+\frac{1}{2})(2\tau + p + i\tau)}$$

~~1 + \frac{1}{2}~~ ω ν ρ σ τ ϵ δ γ β α

$$p = 2\nu + \rho \quad \rho = 2\nu + \rho_1 \quad n + g_j = E_j g_j$$

$$V_{2\nu + \rho_1, 2\nu + \rho_1}(z) = \sum_{-\infty}^{+\infty} e^{\frac{(2n+\frac{1}{2})^2}{2} \tau + \frac{2n+\frac{1}{2}}{2} (2\tau + \rho_1 + i\tau) + 2\nu i\tau}$$

$$e^{\frac{2n+\frac{1}{2}}{2} \tau + i\nu \tau} = e^{i\nu \tau} = (-1)^{i\nu}$$

$$= (-1)^{i\nu} V_{1,1}(z)$$

ν ω γ δ ϵ β α

$$V_{1,1}(z) = \sum_{-\infty}^{\infty} e^{n^2 \tau + n p \tau + \frac{n^2 \tau}{4} + 2n\tau + p\tau + \frac{1}{2} n p i \tau + \frac{1}{2} p^2 \tau}$$

$$V_{0,0}(z) = \sum_{-\infty}^{\infty} e^{n^2 \tau + p\tau + \frac{1}{2} p p i \tau} \sum_{-\infty}^{\infty} e^{n\tau + n p \tau + 2n\tau + n p i \tau}$$

$$2 - 2 + \frac{1}{2}$$

$$v_{10} = e^{2 + \frac{1}{2}\tau} v_{00}(2 + \frac{1}{2}\tau)$$

$$e^{2 + \frac{1}{2}\tau} \sum e^{n\tau} + 2\tau$$

$$= e^{2 + \frac{1}{2}\tau} \sum e^{-\tau(n + \frac{1}{2})} f(n + \frac{1}{2})$$

$$= \sum e^{-\tau(\frac{2n+1}{2})} + \frac{2n+1}{2} 2\tau$$

$$v_{10}(z) = A e^{2 + \frac{1}{2}\tau} \prod_0^{\infty} \left(1 + \rho \frac{e^{2n\tau}}{e^{2n\tau}}\right) \left(1 + \rho \frac{e^{2n\tau - 2\tau}}{e^{2n\tau - 2\tau}}\right)$$

$$= A e^{2 + \frac{1}{2}\tau} \prod_0^{\infty} \left(1 + \rho \frac{e^{2n\tau + 2\tau}}{e^{2n\tau}}\right) \left(1 + \rho \frac{e^{2n\tau - 2\tau}}{e^{2n\tau}}\right)$$

$$= A e^{2 + \frac{1}{2}\tau} (1 + \rho e^{-2\tau}) \prod_1^{\infty} \left(1 + \rho \frac{e^{2n\tau}}{e^{2n\tau}} + \rho \frac{e^{4n\tau}}{e^{4n\tau}}\right)$$

$$= A \rho^{\frac{1}{2}} \text{Li} \text{ with } \prod (1 + \rho \frac{e^{2n\tau}}{e^{2n\tau}} + \rho \frac{e^{4n\tau}}{e^{4n\tau}})$$

$$v_{11}(z) \sim v_{10} \sim 0 \quad z - 2 + \frac{i\pi}{2}$$

$$v_{11} = v_{10} \left(2 + \frac{i\pi}{2}\right) = \sum_{-\infty}^{\infty} e^{\frac{(2n+1)\tau}{2} + \frac{2n+1}{2} 2\tau} \left(2 + \frac{i\pi}{2}\right)$$

+

$$= A e^{2 + \frac{1}{2}i\pi + \frac{1}{2}\tau} \prod_0^{\infty} \left[1 + \rho \frac{e^{2n\tau + 2\tau}}{e^{2n\tau}}\right] \left[1 + \rho \frac{e^{2n\tau - 2\tau}}{e^{2n\tau}}\right]$$

$$u(z) = \prod_{n=1}^{\infty} (1 - q^{2n})$$

$$f(z) = u(z), \quad w\left(\frac{z}{2}\right) = A_0 \sum_{n=-\infty}^{\infty} (-1)^n q^{\frac{1}{2}n(n-1)} 2^n$$

$$q^{\frac{1}{2}n(n-1)} = q^{n^2}$$

$$w\left(\frac{z}{2}\right) = A_0 \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2} \left(\frac{z}{2}\right)^n$$

$$\frac{z}{2} = -e^{2z}$$

by $q = e^{-2z}$
= modul

$$w\left(\frac{z}{2}\right) = \sum_{n=-\infty}^{\infty} q^{n^2} e^{2nz} =$$

$$= \sum_{n=-\infty}^{\infty} e^{n^2 + 2nz}$$

$$D_{01}(z) = \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2} e^{2nz} = \prod_{n=1}^{\infty} (1 + e^{2nz})$$

$$D_{00} = A \prod_{n=1}^{\infty} (1 + q^{2n-1} e^{2z}) (1 + q^{2n} e^{-2z})$$

$$w(z) = \prod_{n=1}^{\infty} (1 + q^{2n-1} e^{2z})$$

$$w\left(\frac{z}{2}\right) = \prod_{n=1}^{\infty} (1 + q^n)$$

4

$$= A \prod_{n=1}^{\infty} \left[1 + q^{2n-1} \cos 2iz + q^{2n} \cos 4iz \right]$$

$$A_{01} = 1 - 2(n+1)e$$

$$= A \prod_{n=1}^{\infty} \left[1 - q^{2n-1} \cos 2iz + q^{2n} \cos 4iz \right]$$

$$f(z) = \sum_{-\infty}^{+\infty} e^{nz + 2n^2} = A \pi^{-1/2} (pe^{2z})^{-1/2} \eta(-z/p)$$

$$f_0(z) = \sum_{-\infty}^{+\infty} (-1)^n e^{nz + 2n^2} = A \pi^{-1/2} (pe^{2z})^{-1/2} \eta(z/p)$$

$$V = \prod_0^{\infty} (1 + \rho e^{2nz + 2n^2})$$

$$\prod_0^{\infty} (1 + \rho e^{2nz - 2n^2})$$

$$1 + \rho e^{2n+1} (e^{2z} + e^{-2z}) + \rho e^{4n+2}$$

$$= \prod (1 + \rho e^{2nz} (e^{2z} + e^{-2z}) + \rho e^{4n+2})$$

$$\Rightarrow = \prod (1 - 2\rho e^{2nz} \cos 2z + \rho e^{4n+2})$$

$$f_0(z) \sim \prod_{n=1}^{\infty} (2 - 2 + 2 + \frac{\pi^2}{4} n^2) \times e^{2z + \frac{\pi^2}{4}}$$

$$\sim \prod_{n=1}^{\infty} (2 + \frac{\pi^2}{4} n^2) \times e^{2z + \frac{\pi^2}{4}}$$

for $f_0(z)$

$$b_n 2^{-n} + b_0 + b_1 2^{-1} + \dots = \sum_{k=0}^{\infty} b_k 2^{-k} + b_0 2^0 + \dots$$

34

$$b_1 = -b_0$$

$$b_2 = -b_1 2$$

$$b_{-n} = -b_{-n+1} 2^{-n+1}$$

$$b_{-n-1} = -b_{-n-2} 2^{-n-2}$$

$$b_{-1} = b_{-2} 2^{-2}$$

$$b_n = -b_{n-1} 2^{n-1}$$

$$b_n = (-1)^n 2^{n(n-1)/2} b_0$$

$$b_n = (-1)^n 2^{n(n-1)/2} b_0$$

~~$$b_0 = (-1)^0 b_0$$~~

$$b_n = (-1)^n 2^{n(n-1)/2} b_0$$

$$f(x) = b_0 \sum_{n=-\infty}^{+\infty} (-1)^n 2^{n(n-1)/2} x^n = x(2) x^{\left(\frac{p}{2}\right)}$$

$$f(x^2) = b_0 \sum_{n=-\infty}^{\infty} (-1)^n 2^{n(n-1)/2} x^{2n} = x(2) x^{\left(\frac{p}{2}\right)}$$

$$= b_0 \sum_{n=-\infty}^{\infty} (-1)^n 2^{n^2} \left(\frac{x}{2}\right)^n$$

$$\left. \begin{array}{l} \frac{x}{2} = e \\ \frac{x}{2} = e \\ 2yq = \tau \\ q = e \end{array} \right\} \text{m.}$$

$$100 \text{ c/f - k } 23 \quad 2 - p2$$

$$f(p2) = w(p2) w\left(\frac{1}{2}\right) = \frac{w(p2)}{1-2} w\left(\frac{1}{2}\right) = \frac{w(p2)}{1-2} \frac{1-1}{2} w\left(\frac{1}{2}\right)$$

$$= -\frac{f(p2)}{2}$$

$$f(z) = -2f(p2) \quad \text{c} \quad f(z) \text{ r } w(z) \text{ w } d \text{ f}$$

$$f(z) = A + \dots + A_n z^n + \dots$$

dypp 3 Ad

$$f(z+w) = \mu e^{2az} f(z)$$

$$f(z+w) = \mu' e^{2az} f(z)$$

d) \sim k

(Part w) = z

$$f(z+w) = e^{az+az}$$

5 ~~(1-2) w(p2)~~

$$f(z) = w(z) w\left(\frac{p}{2}\right) = [a_0 + a_1 z + \dots] [a_0 + a_1 \frac{p}{2} + \dots]$$

$$f(z) = w(z) w\left(\frac{1}{2}\right)$$

$$\frac{w(z)}{1-2} \frac{1-1}{2} w\left(\frac{1}{2}\right) = -\frac{f(z)}{2}$$

$$f(z) = -2f(p2)$$

$$f(z) = \dots + \frac{b_{-1}}{2} + b_0 + b_1 z + b_2 z^2 + \dots$$

$$f(z) = \dots - \frac{b_{-n}}{2^n p^{-n}} + \dots + b_0 + b_1 p z + b_2 p^2 z^2 + \dots$$

$$r^0 p^2 - p^2 r$$

$$w(p^2) = (1-p^2) w(p^2) \quad w(0) = (1-2) w(p^2)$$

$$= A_0 + A_1 r +$$

$$A_0 + A_1 r + A_2 r^2 + \dots = (1-2) [A_0 + A_1 p^2 + A_2 p^4 + \dots]$$

$$= A_0 + A_1 p^2 + A_2 p^4 + A_3 p^6 + \dots$$

$$- A_0 2 - A_1 p^2 - A_2 p^4 - A_3 p^6 - \dots$$

$$A_0 = A_0$$

$$A_1 = A_1 p - A_0$$

$$A_1 = \frac{A_0}{p-1}$$

$$A_2 = A_2 p^2 - A_1 p$$

$$A_2 = \frac{A_1 p}{p^2-1} = \frac{A_0 p}{(p^2-1)(p-1)}$$

$$A_3 = A_3 p^3 - A_2 p^2$$

$$A_3 = \frac{A_2 p^2}{p^3-1} = \frac{A_0 p^3}{(p^3-1)(p^2-1)(p-1)}$$

$$A_n = A_n p^n - A_{n-1} p^{n-1}$$

$$A_n = \frac{A_0 p^n}{(p-1)(p^2-1) \dots (p^n-1)} \quad p^{2n}$$

$$A_0 = 1$$

$$A_n = \frac{p^{2n}}{(p^n-1)(p-1)}$$

if $c=10$ for $w(0) = w(\frac{1}{2})$

$$b_2(u)^2 - b_3(u)^2 + [e_1 - e_3] b(u)^2 = 0$$

$$\frac{b_2}{b_3} = 1 - g_0$$

$$\left(\frac{b_2(u)}{b_3(u)}\right)^2 - 1 + \frac{[e_1 - e_3] \sqrt{u}}{e_1 - e_3} \sin \text{am} \sqrt{u}$$

$$\sin \text{am} \sqrt{e_1 - e_3}$$

36

$$\left(\frac{b_2(u)}{b_3(u)}\right)^2 = 1 - \sin^2 \text{am} \sqrt{u}$$

$$= \cos^2 \text{am}(\sqrt{e_1 - e_3} u, k)$$

$$b_1(u)^2 - b_3(u)^2 + (e_2 - e_3) b(u)^2 = 0$$

$$\left(\frac{b_1(u)}{b_3(u)}\right)^2 - 1 + \frac{(e_2 - e_3)}{e_2 - e_3} \sin^2 \text{am}(\sqrt{e_2 - e_3} u, k) = 0$$

$$\left[\frac{b_1(u)}{b_3(u)}\right]^2 = 1 - k^2 \sin^2 \text{am}(\sqrt{e_2 - e_3} u, k)$$

$$\frac{b_1(u)}{b_3(u)} = \Delta \text{am}$$

$$K = \int_0^1 \frac{dt}{\sqrt{1-t^2} \sqrt{1-kt^2}} \quad \text{Periode} \quad \sim \sqrt{e_1 - e_2} c$$

$$\left[\frac{b_2(u)}{b_3(u)}\right]^2 = p(u) - e_2$$

$$\left[\frac{b_3(u)}{b_1(u)}\right]^2 = p(u) - e_3$$

$$\text{oder } w_1 = \frac{w}{2} \quad w$$

$$\frac{b_3(u)}{b_1(u)} = \sqrt{e_1 - e_3}$$

$$\left(\frac{dx}{dt}\right)^2 = [1 - (e_p - e_a) z^2] [1 - (e_f - e_d) z^2]$$

$$X_{0a}$$

$$\alpha = 3$$

$$\rho = 2$$

$$j = 1$$

$$\left(\frac{dz}{dt}\right)^2 = [1 - (e_2 - e_1) z^2] [1 - (e_3 - e_1) z^2]$$

$$\sqrt{e_3 - e_1} z = \int$$

$$\frac{1}{e_3 - e_1} \left(\frac{dz}{dt}\right)^2 = \left[1 - \frac{e_2 - e_1}{e_3 - e_1} z^2\right] [1 - z^2]$$

$$\frac{e_2 - e_1}{e_3 - e_1} = k^2$$

$$\frac{dz}{\sqrt{e_3 - e_1} z} = [1 - k^2 z^2] [1 - z^2]$$

$$\left(\frac{dz}{dt}\right)^2 = [1 - k^2 z^2] [1 - z^2]$$

$$z = \sin am v$$

$$= \sqrt{e_3 - e_1} X_{0a}$$

am v = arcsin z

$$\frac{\partial \omega}{\partial z} = \frac{1}{\sqrt{e_3 - e_1}} \sin am v = \frac{1}{\sqrt{e_3 - e_1}} \sin am (\sqrt{e_3 - e_1} X_{0a})$$

$$\int \frac{dx}{\sqrt{(2-x)(2kx-y)}} = \text{am } u \quad z = p(u)$$

$$= \int \frac{dy}{\sqrt{(1-y)(2kxy)}} = \int \frac{dy}{\sqrt{(1-y)(2kxy)}}$$

$$y = \text{am } u$$

$$\sigma_r(-\omega) = \sigma_r(\omega) \quad \text{if } \rho$$

$$\sigma_1(\omega) \sigma_2(\omega) = \rho \sigma^2$$

$$\sigma_1(\omega) \rho \approx \frac{\sigma_1(\omega)}{\omega} \approx \frac{\sigma_2(\omega)}{\omega} \quad \sigma_3(\omega) \approx \frac{\omega + \omega'}{\omega}$$

$$\sigma_1^2(\omega) = \mu(\omega) - \mu\left(\frac{\omega}{2}\right)$$

$$\mu_1 \mu_2 / 3 = 0$$

$$\frac{\sigma_1(\omega)}{\sigma(\omega)} = \sqrt{\mu(\omega) - \mu_1} \quad \mu_1$$

$$\frac{\sigma_2(\omega)}{\sigma(\omega)} = \sqrt{\mu(\omega) - \mu_2}$$

$$\frac{\sigma_3(\omega)}{\sigma(\omega)} = \sqrt{\mu(\omega) - \mu_3}$$

$$\mu(\omega) = \frac{1}{3} \frac{\sigma_1^2(\omega) + \sigma_2^2(\omega) + \sigma_3^2(\omega)}{\sigma^2(\omega)}$$

$$\left[\frac{\sigma_1(\omega) \sigma_2(\omega) \sigma_3(\omega)}{\sigma^3(\omega)} \right]^2 = \mu(\omega)^2$$

~~+~~

$$\mu'(\omega) = \pm \frac{2 \sigma_1(\omega) \sigma_2(\omega) \sigma_3(\omega)}{\sigma^3(\omega)}$$

x -

$$\sigma_1(\omega) = \omega +$$

$$\sigma_2(\omega) = 1 + \omega^2$$

Period

$$\sigma_1(\omega) = e^{2\pi i \omega} \sigma_2(\omega)$$

σ^2 ref. per ω

$\omega < \omega' < \omega''$ / 2 Per. f. $\sigma(\omega) = \pm \sigma_1(\omega) + \sigma_2(\omega)$

$$\sigma_1'(\omega) = \mp e^{2\pi i \omega} \left[\sigma_1'(\omega) + \sigma_2'(\omega) \right]$$

1. $\sigma(u)$

$$\frac{\sigma(u)}{u} = \prod_{\lambda \in \Lambda} \left[1 - \frac{u}{\lambda} e^{\frac{u}{\lambda} + \frac{u^2}{2\lambda^2} + \dots} \right]$$

$$w = \mu u + \nu u^2$$

$$\frac{u}{h} \frac{\lambda u + \lambda' u'}{\lambda u}$$

$$\sigma(\lambda u, \lambda u', \lambda u'') = \prod (1 - \lambda u)$$

$$\sigma(\lambda u, \lambda u') = \lambda \sigma(u) \prod \dots$$

10) $\lambda = \frac{1}{\omega} \quad \sigma = \sigma\left(\frac{u}{\omega}, \frac{u'}{\omega}\right)$

zu f. $\mu, \nu, \gamma, \rho, h, e$

$$\sigma\left(u + \frac{u}{\omega}\right) = \sigma\left(u - \frac{u}{\omega} + u\right) = -e^{\gamma \left(u - \frac{u}{\omega} + \frac{u}{\omega}\right)} \sigma\left(u + \frac{u}{\omega}\right)$$

$$\frac{\sigma\left(u - \frac{u}{\omega}\right)}{\sigma\left(u + \frac{u}{\omega}\right)} = e^{\gamma u} \frac{\sigma\left(\frac{u}{\omega} - u\right)}{\sigma\left(\frac{u}{\omega}\right)}$$

$$\sigma(-u)$$

$$\sigma(u - \rho h)$$

$$e^{\gamma \rho h} \sigma(u - \rho h)$$

$\sigma(u, u')$...

$$F(u+v) = F(u)$$

$$F(u+iv) = e^{g(u)} \cdot F(u)$$

Prinzip der Summe

Wichtig für die Laplace-Transformation

$$F(u+iv) = F(u+iv) = e^{g(u+iv)} F(u+iv)$$

$$e^{g(u)} F(u) = \uparrow F(u)$$

$$F(u) = e^{g(u)} \cdot e^{-g(u)} = 1$$

$$g(u+iv) - g(u) = 2k\pi$$

erst $g(u)$ linear ~

in v $g(u)$ & linear $g(u)$

die 2 Funktionen $g(u)$ & $g(u+iv)$ sind stetig

✓ 6. v. eff. - e. h. v. Adygen

$$b_1(u) = e^{-\eta \frac{u}{2}} \frac{b(u + \frac{u}{2})}{b(\frac{u}{2})}$$

$$b_2(u) = e^{-\eta \frac{(u+iv)}{2}} \frac{b(u + \frac{u+iv}{2})}{b(\frac{u+iv}{2})}$$

$$b_3(u) = e^{-\eta \frac{u}{2}} \frac{b(u + \frac{u}{2})}{b(\frac{u}{2})}$$

} ob h.

$$m, \varphi(u) = \sum_{k=0}^n A_k u^k$$

$$\chi(u) = \sum_{k=0}^{n+1} B_k u^k \quad \text{with } B_0 = 1$$

$$\varphi(u) + \chi(u) - \chi(u+w) = 2\pi i n$$

$$\sum_{k=0}^n A_k u^k + \sum_{k=0}^{n+1} B_k [u^k - (u+w)^k] = 2\pi i n$$

~~is~~

$$\chi(u+w) = B_{n+1} [u^{n+1} + \binom{n+1}{1} u^n w + \binom{n+1}{2} u^{n-1} w^2 + \dots + \binom{n+1}{n} u w^n + w^{n+1}]$$

$$+ B_n [u^n + \binom{n}{1} u^{n-1} w + \dots + \binom{n}{n-1} u w^{n-1} + w^n]$$

+ ...

$$+ B_1 [u + w]$$

+ ~~B_0~~

$$- B_{n+1} u^{n+1} - B_n u^n = \dots \equiv \varphi(u) + 2\pi i n$$

Is $\chi(u) = B_0 + B_1 u + B_2 u^2 + \dots + B_{n+1} u^{n+1}$
 $\varphi(u) = 0$

$$a u - a' u + (\lambda - \nu) u = \dots$$

$$A u'' + 2B u' + (\lambda - \nu) u = \dots + 3' u$$

$$0 u' - a' u = A u u' [u - u'] + (\mu - \nu) [i u (u' - u) + \dots + \frac{u u' (\eta - \eta')}{\nu}] - [\eta u' - \eta' u] u u'$$

$$A u = \frac{a}{\nu} - (\mu - \nu) \frac{\eta}{\nu}$$

$$\int \eta' \mu = \nu \quad \eta' \mu = \nu \quad \dots$$

$$A = 0$$

$$\text{Hermitite: } 0 u' - 0' u = \dots (\Sigma u - \Sigma') u u'$$

Lösungsgleichung

$\eta = \nu \sqrt{\mu h} \quad | \quad \eta' = \dots$

$$f(u) = e^{\varphi(u)}$$

$$f(u) = e^{\varphi(u)} f(u)$$

$$f(u) = e^{\chi(u)} F(u)$$

$$u \chi' = \varphi' + \dots$$

$$\chi' = \dots$$

$$F(u) = F(u)$$

$$F(u) = \dots$$

$$f(u) = \frac{e^{\varphi(u)}}{e^{\chi(u)}} = e^{\varphi(u) - \chi(u)} F(u)$$

$$F(u) = \dots$$

$$= e^{\varphi(u)} f(u) = e^{\varphi(u) + \chi(u)} F(u)$$

$$F(u) = e^{\varphi(u) + \chi(u) - \chi(u)} F(u)$$

$$\nu' \chi' = \dots = 10$$

C. Puffin \times dark Day \times Rant \times 5 eggs
 y constante in Wied. Ann. XXVIII 87-107
 1886
 Puffin. 475 $\mu = 0.458. 500 \times 2 \mu$

Rax h 58 M395
 tambunak 59
 Kishun 60

Zhu mens. Nr 2
 Cinn p. 40
 Reth 48
 Carg Zunkh.
 Zchunke 49
 Langk
 12/9 Cinn dran Zunkh
 3 NW
 13/9 P. dan men Wy
 mit Samsel & Dmai

$$1) \mu(u) = \ln \alpha = \frac{\ln^2(u)}{6(u)}$$

41

$$\mu'(u) = -2 \frac{\ln(u) \cdot \ln(u) \cdot \ln(u)}{6^2(u)}$$

$$\ln^2(u) - \ln(u)^2 + \ln(u) \cdot \ln(u) \cdot \ln(u) = 0$$

$$(\ln(u) - \ln(u)) \ln(u)^2 + (\ln(u) - \ln(u)) \ln(u) \cdot \ln(u) = 0$$

$$X_{\alpha 0} = \frac{\ln(u)}{6(u)} \quad X_{\rho \gamma} = \frac{\ln(u)}{6(u)} \quad X_{0\alpha} = \frac{\ln(u)}{6(u)}$$

s.D.

$$X_{\alpha 0}^2 = \mu(u) - \ln$$

$$2) \text{ diff. } 2 X_{\alpha 0} X'_{\alpha 0} = \mu'(u) = -2 X_{\alpha 0} X_{\rho 0} X_{\gamma 0}$$

dividi

$$I) X'_{\alpha 0} = - X_{\rho 0} X_{\gamma 0}$$

$$X_{\rho \gamma} = \frac{X_{\rho 0}}{X_{\gamma 0}} \quad X'_{\rho \gamma} = \frac{X_{\gamma 0} X'_{\rho 0} - X_{\rho 0} X'_{\gamma 0}}{X_{\gamma 0}^2}$$

$$= \frac{-X_{\alpha 0} X_{\gamma 0}^2 + X_{\rho 0}^2 X_{\alpha 0}}{X_{\gamma 0}^2}$$

$$= X_{\alpha 0} \left[\frac{X_{\rho 0}^2 - X_{\gamma 0}^2}{X_{\gamma 0}^2} \right]$$

$$\sigma_{p_0}^2 - \sigma_{y_0}^2 + (e_p - e_y) \sigma_{u^2} = 0$$

$$\text{III } X_{p_0}^2 - X_{y_0}^2 = -(e_p - e_y)$$

$$\text{II } X'_{py} = - (e_p - e_y) X_{xy} X_{oy}$$

$$X_{\alpha\alpha} = \frac{1}{X_{\alpha\alpha}}$$

$$\text{IV } X'_{\alpha\alpha} = \frac{-X_{\alpha\alpha}}{X_{\alpha\alpha}^2} = \frac{X_{p_0} X_{y_0}}{X_{\alpha\alpha}^2} = X_{p\alpha} X_{y\alpha}$$

- 2. 2. 1. $X_{\alpha\alpha}$ 1. 1.:

$$X_{\alpha\alpha}^2 = X_{p_0}^2 X_{y_0}^2 \quad \parallel \quad \sigma_{\alpha\alpha}^2 - \sigma_{p_0}^2 + [e_p - e_y] \sigma_{u^2} = 0$$

$$X_{p_0}^2 = X_{\alpha\alpha}^2 + (e_p - e_y)$$

$$X_{y_0}^2 = X_{\alpha\alpha}^2 + (e_y - e_p)$$

$$X'_{\alpha\alpha} = [e_p - e_y + X_{\alpha\alpha}^2] [e_y - e_p + X_{\alpha\alpha}^2]$$

$$X'_{py} = (e_p - e_y) X_{xy} X_{oy}$$

$$\sigma_p(u) - \sigma_y(u) + [\epsilon_p - \epsilon_y] \sigma(u) = 0$$

42

$$X_{py}^2 - 1 + [\epsilon_p - \epsilon_y] X_{oy}^2 = 0$$

$$\rightarrow [\epsilon_p - \epsilon_y] X_{oy}^2 = 1 - X_{py}^2$$

$$[\epsilon_\alpha - \epsilon_p] \sigma_y(u) + [\epsilon_p - \epsilon_y] \sigma_\alpha(u) + [\epsilon_y - \epsilon_\alpha] \sigma_p(u) = 0$$

$$[\epsilon_\alpha - \epsilon_p] + [\epsilon_p - \epsilon_y] X_{\alpha y}^2 + [\epsilon_y - \epsilon_\alpha] X_{py}^2 = 0$$

$$[\epsilon_p - \epsilon_y] X_{py}^2 = [\epsilon_p - \epsilon_\alpha] + [\epsilon_\alpha - \epsilon_y] X_{py}^2$$

$$X'_{py}^2 = [1 - X_{py}^2] [\epsilon_p - \epsilon_\alpha] + (\epsilon_\alpha - \epsilon_y) X_{py}^2$$

$$X'_{\alpha\alpha} = X_{p\alpha} X_{y\alpha}$$

$$X'_{\alpha\alpha} = X_{p\alpha} X_{y\alpha}$$

$$\sigma_\alpha(u) - \sigma_p(u) + (\epsilon_\alpha - \epsilon_p) \sigma(u) = 0$$

$$1 - X_{p\alpha}^2 + [\epsilon_\alpha - \epsilon_p] X_{\alpha\alpha}^2 = 0$$

$$X_{p\alpha}^2 = 1 + (\epsilon_\alpha - \epsilon_p) X_{\alpha\alpha}^2$$

$$X_{y\alpha}^2 = 1 + (\epsilon_\alpha - \epsilon_y) X_{\alpha\alpha}^2$$

$$X'_{\alpha\alpha} = [1 - (\epsilon_p - \epsilon_\alpha) X_{\alpha\alpha}^2] [1 - (\epsilon_y - \epsilon_\alpha) X_{\alpha\alpha}^2]$$

1. 8/2 e f. 323 & 294 17

$$\left(\frac{dx}{dt}\right)^2 = [1 - (ep - ea)z^2][1 - (ey - ea)z^2]$$

20. / 2:

$$X_{\alpha 0} = \sqrt{ep - ea} \sqrt{ey - ea} z$$

$$(ea - ep) + X_{\alpha 0}^2 = ea - ep + (ep - ea)(ey - ea)z^2$$

$$\pi = -(ep - ea) [1 - (ey - ea)z^2]$$

$$(ea - ep) + X_{\alpha 0}^2 = -(ey - ea) [1 - (ep - ea)z^2]$$

∴ (ep - ea) z

$$z = \frac{X_{\alpha 0}}{\sqrt{ep - ea} \sqrt{ey - ea}} \quad \left. \begin{array}{l} \text{" = } \sqrt{ey - ea} \\ \text{" = } X_{\alpha 0} \end{array} \right\}$$

$$\text{II } X_{\beta 1} = \sqrt{ep - ea} z \quad \left. \begin{array}{l} \text{E h 27 = 1} \\ \text{wog : } z = \frac{X_{\beta 1}}{\sqrt{ep - ea}} \end{array} \right\}$$

28/ 29/ 30/

∴ 17 f. 6 Quat. f. 296.

f. B.H. = < , B.H. coll. fu.

Long B.H. = $6^m u = 0$

$$b(0) = 0 \quad X_{0\alpha}(0) = 0$$

$$b_2(0) = 1$$

$$\text{en } X_{12} = \left. \frac{\partial p(u)}{\partial y(u)} \right|_{u=0} = 1$$

$$X_{20} \Big|_0 = \infty$$

$r^0 p/r \quad \alpha = 3, \beta = 2, \gamma = 1$

$$\left(\frac{dr}{du}\right)^2 = [1 - (l_2 - l_1) r^2] [1 - (e_1 - l_3) r^2]$$

$$2 = X_{03}$$

$$2 = \frac{X_{30}}{\sqrt{e_1 - l_3} \sqrt{l_2 - l_1}}$$

$$2 = \frac{X_{21}}{\sqrt{e_1 - l_3}}$$

Transformation

$$y = \sqrt{e_1 - l_3} r$$

$$= k^2$$

~~$$\frac{dr}{du} (l_2 - l_1) =$$~~

$$\frac{1}{e_1 - l_3} \left(\frac{dy}{du}\right)^2 = [1 - y^2] \left[1 - \frac{(l_2 - l_1)}{e_1 - l_3} y^2\right]$$

f = 1/1000000

$$\frac{1}{(s_1 - s_2)} \left(\frac{ds}{dt} \right)^2 = [1 - \beta^2] [1 - u^2 \beta^2]$$

in a m = $\frac{6(u)}{63(u)}$ $\beta = \frac{1}{1000000}$

BJ

$$q = e^{i\alpha \frac{u}{\omega}}$$

$$f(u) = \sum a_n e^{\frac{m i \alpha n}{\omega}}$$

$$m a_{n+k} = a_n \rho^{2m}$$

$$m^{\rho} a_{\rho k} = a_0 \rho^{k \rho - 1}$$

$$\alpha < k_{m=\alpha}$$

$$m^{\rho} a_{\rho k}$$

$$m a_{2k+\alpha} = a_{k+\alpha} \rho^{2k+\alpha}$$

$$m a_{\rho k + \alpha} = a_{(\rho-1)k + \alpha} \rho^{2\rho-1 k + \alpha}$$

$$m^{\rho} a_{\rho k + \alpha} = a_{\alpha} \rho^{2\alpha \rho + \rho \rho - 1 k}$$

$$f(u) = \sum_{n=0}^{k-1} \left(a_n \sum_{m=-\infty}^{+\infty} \rho^{\frac{2m\rho + \rho \rho - 1 k}{m}} e^{\frac{(k+n) i \alpha n}{\omega}} \right)$$

$$u \in \Sigma \theta_{\rho}$$

$$\theta_{\rho}(u) = \sum e^{\frac{(2m+\rho)u}{2} + \frac{(2n+\rho)}{2}(2u+\rho i \pi)}$$

$$= \sum e^{m^2 \tau + \frac{\rho^2}{4} \tau + 2m\rho \tau + 2nu + \rho u + i \pi n + \frac{\rho i \pi}{2}}$$

$$= \sum_{-\infty}^{+\infty} \cancel{2\rho} e^{2\rho r \frac{i\pi d'}{\omega} + \rho^2 \kappa i\pi \frac{d'}{\omega} - \kappa \rho i\pi \frac{d'}{\omega}}$$

$$\theta_{pp^{(2)}} = \frac{i\pi}{\omega} n = 2z \quad + 2z\rho\kappa + 2z^2$$

$$i\pi \frac{d'}{\omega} = \tau$$

$$= \sum e^{2\rho n\tau + \rho^2 \kappa \tau - \kappa \rho \tau + 2\rho \kappa z + 2z^2}$$

$$p \ll z \ll \sqrt{\tau} \ll \kappa$$

$$(-\rho^2 \gamma \mu)$$

$$\tau' = \kappa \tau$$

$$= \sum e^{n^2 \tau' + \frac{2\rho n}{\kappa} \tau' - n\tau' + 2z(\rho\kappa + \tau) - n^2 \gamma \mu}$$

$$= \sum e^{[n^2 + \frac{2\rho n}{\kappa} - n] \tau' + (n\kappa + z) 2z}$$

$$\sum e^{(n^2 + \frac{2\rho n}{\kappa} + \frac{\rho^2}{\kappa^2}) \tau' - \frac{\rho^2}{\kappa^2} \tau' +}$$

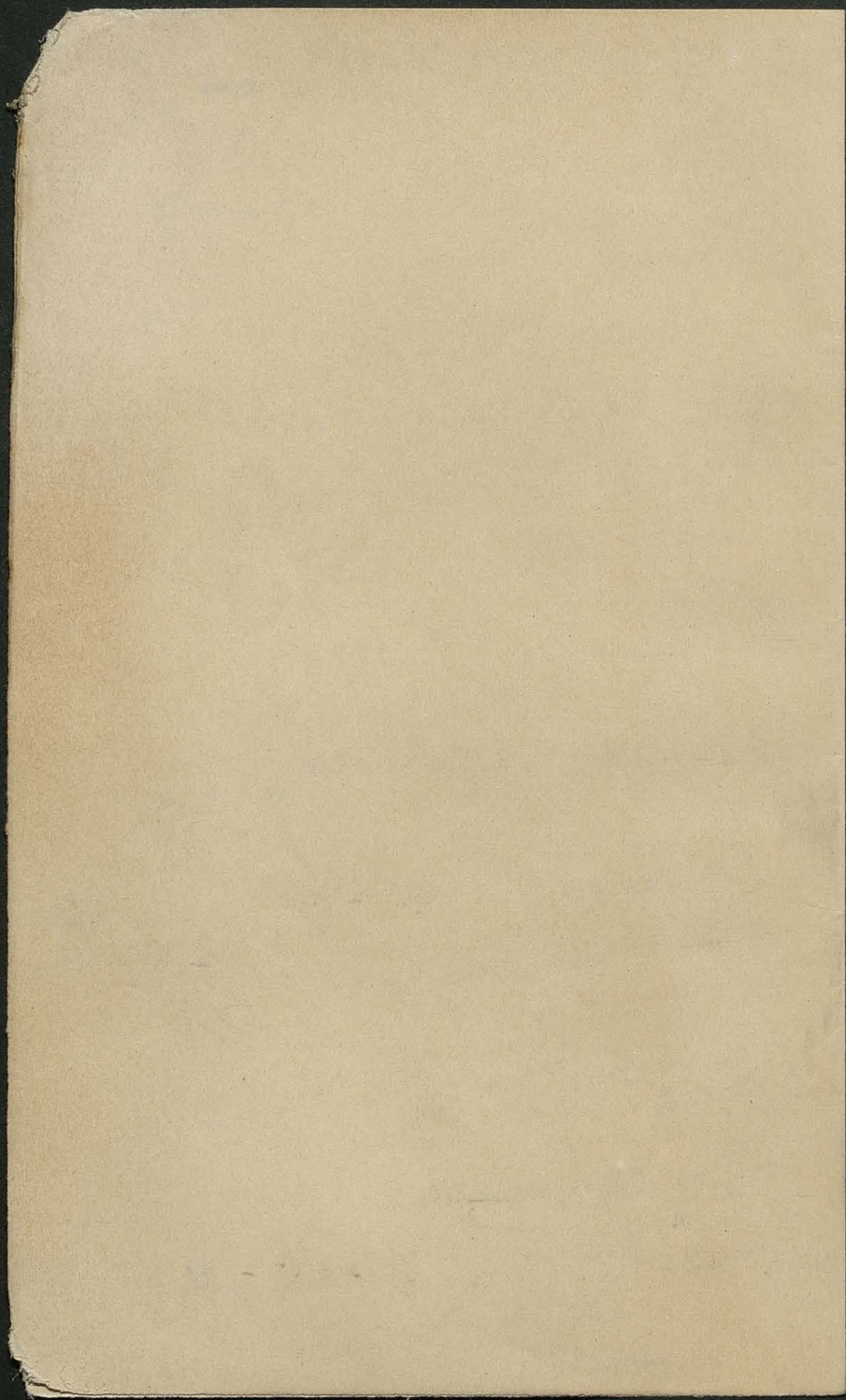
$$= \sum e^{-\frac{\rho^2}{\kappa^2} \tau' + \frac{\rho^2 \tau'}{\kappa} \sum e^{(n + \frac{\rho}{\kappa})^2 \tau' + 2z(\frac{n\rho}{\kappa} + z)}$$

-log p

$$\frac{5'}{k'} = \frac{5''}{4}$$

$$2k+2 = 2m+r$$

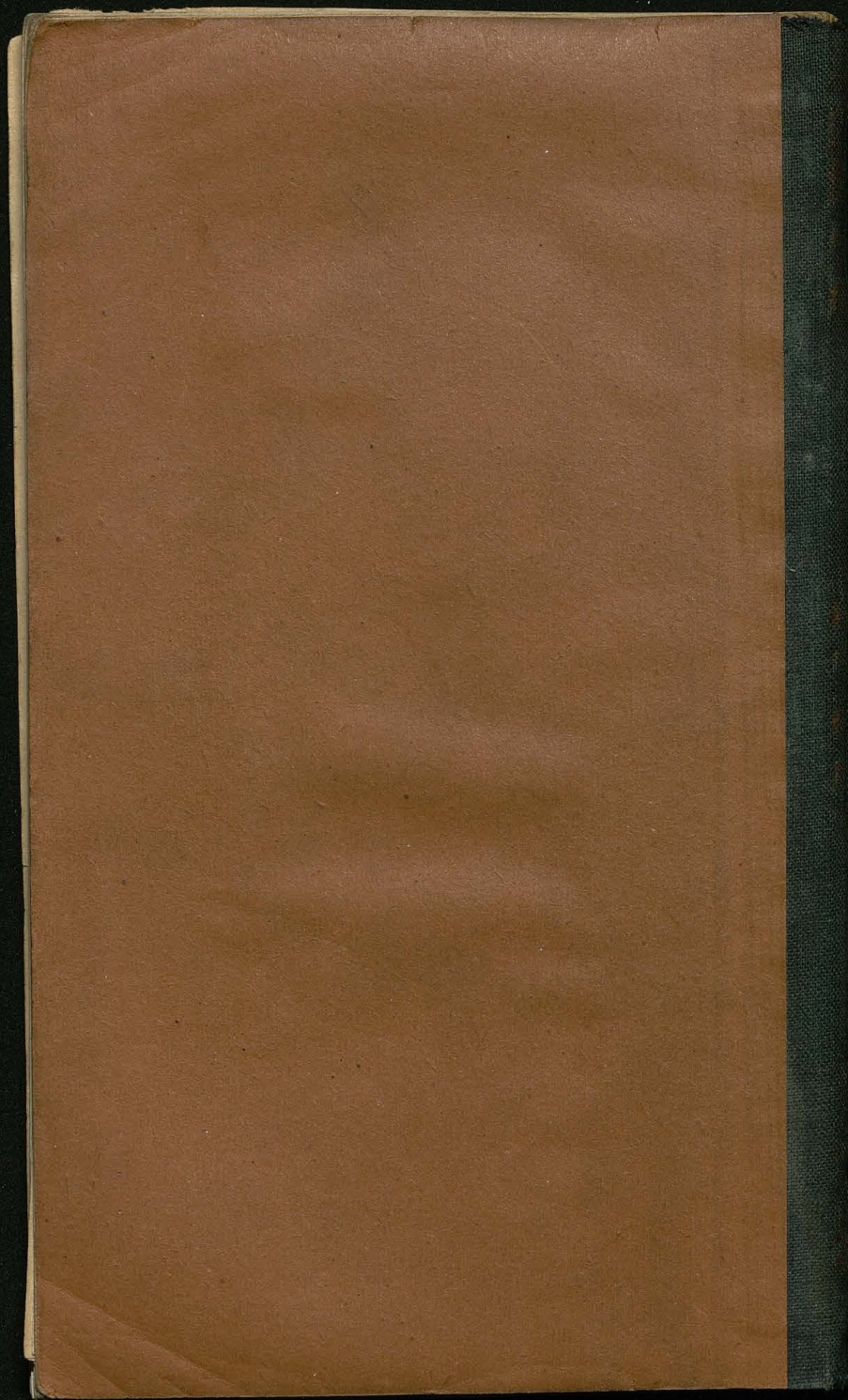
46



A. Reynolds On rolling friction Phil. Trans. 47
155-175 CL XVII

Phil. mag. (5) I. 75-77

Forbes & Co. m. J. 1876 p. 200





9451

148

$f(x,y,z) = 0$ and $f_p dx + f_q dy + f_r dz = 0$

BJ

Path - Curve

$f_1 dx + f_2 dy + f_3 dz = 0$

$q_1 dx + q_2 dy + q_3 dz = 0$

$\frac{dx}{ds} = \cos \alpha$

$f_1 \cos \alpha + f_2 \sin \alpha + f_3 \sin \alpha = 0$

or $f_1 \cos \alpha + f_2 \sin \alpha + f_3 \sin \alpha = 0$

$q_1 \cos \alpha + q_2 \sin \alpha + q_3 \sin \alpha = 0$

or $q_1 \cos \alpha + q_2 \sin \alpha + q_3 \sin \alpha = 0$ on C.O.

and: $\cos \alpha : \sin \alpha : \sin \alpha = q_1 - x : q_2 - y : q_3 - z$

$f_1(q_2 - y) + f_2(q_3 - z) + f_3(q_1 - x) = 0$

By 1 there is a unique tangent plane at (x,y,z)

= Tangent plane

1. The normal to the plane is ∇f .

2. $f(x,y,z) = 0$ is the equation of the plane.

or ∇f is normal to the plane.

or $\nabla f \cdot \nabla f = 0$

$$z = f(x, y)$$

$$= f(x) + f_1 x^2 + 2 f_1 x y + 2 f_1 y^2$$

$$+ (\quad)^3 + - (\quad)^4$$

$$\omega z = 0 \text{ ist die XY}$$

1. f_1 ist die erste Ableitung von f nach x oder y .

2. f_2 ist die zweite Ableitung von f nach x oder y .

3. f_{xy} ist die gemischte Ableitung.

$$f_{xx} - f_1 f_{xy} > 0 \text{ ist ein lok. Min.}$$

$$< 0 \text{ ist ein lok. Max.}$$

$$f(x, y) - z = 0$$

$$\frac{\partial z}{\partial x} = f \quad \frac{\partial z}{\partial y} = g$$

$$\frac{\partial^2 z}{\partial x^2} = f_{xx} \quad \frac{\partial^2 z}{\partial x \partial y} = f_{xy} \quad \frac{\partial^2 z}{\partial y^2} = g_{yy}$$

$$f(f_{xx} - f_{xy}) + g(f_{xy} - g_{yy}) - (f - z) = 0$$

$$\beta x: \frac{f}{\sqrt{f_{xx}^2 + f_{xy}^2}} \quad \frac{g}{\sqrt{f_{xy}^2 + g_{yy}^2}} \quad \frac{-1}{\sqrt{f_{xx}^2 + f_{xy}^2}}$$

$$dx = f dx + g dy$$

375 2b. $\Delta \cos \theta = N \sqrt{R \cdot N \cos \theta} = \theta$ 51

$$\cos \theta = \frac{p \cos \alpha + q \cos \beta}{\sqrt{p^2 + q^2}}$$

$$d \cos \theta = p d \cos \alpha + q d \cos \beta + d p \cos \alpha + d q \cos \beta$$

$$d \cos \theta = d(p \cos \alpha + q \cos \beta) + d p \cos \alpha + d q \cos \beta$$

$$\neq d(p \cos \alpha + q \cos \beta) = -d p \cos \alpha + d q \cos \beta$$

$$d \cos \theta = - \frac{d p \cos \alpha + d q \cos \beta}{\sqrt{p^2 + q^2}}$$

$$d p = r d \alpha + s d \beta = d s (r \cos \alpha + s \cos \beta)$$

$$d q = s d \alpha + t d \beta =$$

$$= d s \cos \alpha + d t \cos \beta$$

$$\frac{d \cos \theta}{d s} = - \frac{r \cos \alpha + 2 s \cos \alpha \cos \beta + t \cos \beta}{\sqrt{1 + r^2 + s^2}}$$

$\frac{1}{R}$

$$\frac{\cos \theta}{R} = \frac{r \cos \alpha + 2 s \cos \alpha \cos \beta + t \cos \beta}{\sqrt{1 + r^2 + s^2}}$$

f. No. - by Radi. d. 1375 2. cos²

$$\sin \alpha + \cos \beta = 1$$

$$\cos \beta = \sin \alpha$$

$$\frac{1}{R} = \frac{r \sin^2 \alpha + 2s \sin \alpha \cos \alpha + t \cos^2 \alpha}{\sqrt{\dots}}$$

Extremum Werte: $\frac{d}{d\alpha} \dots = 0$

$$-2r \sin \alpha \cos \alpha + 2s \cos 2\alpha + 2t \sin \alpha \cos \alpha$$

$$-(r-t) \sin 2\alpha + 2s = 0$$

$$\sin 2\alpha = \frac{2s}{r-t}$$

Wahl $\frac{1}{2}$ oder $\frac{3}{2}$ Rad.

... $\frac{1}{2}$ oder $\frac{3}{2}$ Rad. ...

... $\frac{1}{2}$ oder $\frac{3}{2}$ Rad. ...

... $\frac{1}{2}$ oder $\frac{3}{2}$ Rad. ...

... $\frac{1}{2}$ oder $\frac{3}{2}$ Rad. ...

f. N. ... Haupt N. ...

Haupt N. ...

... $\frac{1}{2}$ oder $\frac{3}{2}$ Rad. ...

$$\frac{1}{R} = r_1 \cos^2 \alpha + r_2 \sin^2 \alpha \quad | \quad \alpha = 0$$

$$\alpha = 0 \quad \dots \quad \frac{1}{R_1} = r_1 \quad \left| \quad \frac{d}{d\alpha} = 0 \right. \quad \frac{1}{R_2} = r_2$$

$$\frac{1}{R} = \frac{\cos^2 \alpha}{R_1} + \frac{\sin^2 \alpha}{R_2} \quad \text{Euler'sche Regel}$$

$$\hookrightarrow \rho \sim f \text{ an } \perp \text{ s } \sim r_1$$

$$\frac{1}{R_1} = \frac{\sin^2 \alpha}{R_1} + \frac{\cos^2 \alpha}{R_2}$$

$$\frac{1}{R} + \frac{1}{R_1} = \frac{1}{R_1} + \frac{1}{R_2} \quad \text{const.}$$

($\hookrightarrow \rho \sim f$ an der Schnittstelle Rad. R_1 & R_2
 $\hookrightarrow \frac{1}{R_1} = \text{Rad. V. } \rho \text{ Wp. } R \text{ Regel } (M_1 \text{ an } \rho)$)

$$\frac{1}{R_1} = \frac{1}{R_1} + \underbrace{\cos^2 \alpha \left(\frac{1}{R_2} - \frac{1}{R_1} \right)}_{\omega f > 0 \cdot 0 \sim \frac{1}{R_1} > \frac{1}{R_2}}$$

$$\frac{1}{R_2} - \frac{1}{R_1} > 0 \quad \frac{1}{R_1} > \frac{1}{R_2}$$

$$\cancel{R_1 - R_2} < 0 \quad R_1 \text{ Max.}$$

$$\frac{1}{R} = \frac{1}{R_2} + \cos^2 \alpha \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

$$R_1 - R_2 > 0$$

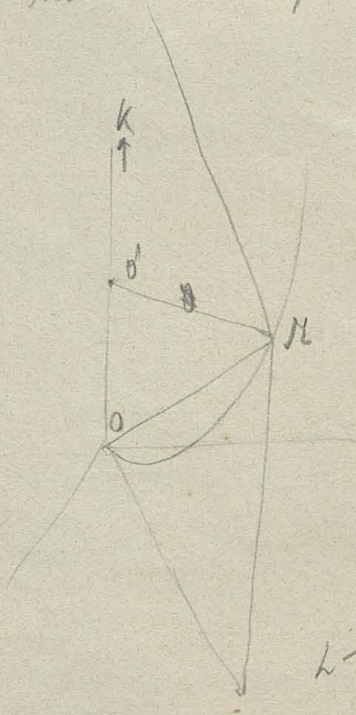
R_1 Max.

$$\frac{1}{R_1} < \frac{1}{R_2}$$

R_2 Min.

$$\frac{1}{R_2} > \frac{1}{R_1}$$

Théorème de Dupin



$$MO' \approx OO', OK$$

$$OK = OO' + O'K = OO' + \frac{rO'^2}{OO'}$$

$$R = \frac{1}{2} \lim \frac{MO'^2}{OO'}$$

$$= \frac{1}{2} \lim \frac{u^2}{2}$$

$h = 2 \text{ } \dots \text{ } \dots$

$$= \frac{1}{2} h \lim \frac{u^2}{2}$$

$$\lim \frac{u^2}{2} = \rho$$

$$R = \frac{\rho^2}{2h}$$

$\rho = \dots$

1 H. 10 γ_0 λ $\approx \gamma_0$ R Hk.

1 H. T₂ γ_1 \approx p dasymp γ_1

$$\frac{d\psi}{R} = \frac{a \sin \alpha + b \sin \alpha \cos \beta + t \sin \beta}{\sqrt{1 + p^2 + q^2}} = \frac{1}{R} \text{ dasymp}$$

$$\sin \beta = p \sin \alpha + q \cos \alpha$$

$$\sin \alpha + \cos \beta + (p \sin \alpha + q \cos \alpha) = 1$$

$$\sin \alpha (1 + p) + \dots = 1$$

$$\frac{R}{\sqrt{1 + p^2 + q^2}} = \frac{(1 + p) \sin \alpha + 2pq \cos \alpha + t \cos \beta}{\sin \alpha + \cos \beta + (p \sin \alpha + q \cos \alpha)}$$

$$\frac{\sin \alpha}{\cos \beta} = \lambda$$

$$\lambda^2 R \left\{ \frac{R}{\sqrt{1 + p^2 + q^2}} - (1 + p) \right\} +$$

$$+ 2\lambda \left\{ \frac{Rq}{\sqrt{1 + p^2 + q^2}} - pq \right\} + \left\{ \frac{Rt}{\sqrt{1 + p^2 + q^2}} - (1 + q^2) \right\} = 0$$

or $\lambda^2 \sqrt{1 + p^2 + q^2} + 2\lambda pq + R = 0$

for a \forall λ p, q, R $\lambda^2 + 2\lambda \frac{pq}{R} + 1 = 0$

$$\text{or } \lambda^2 (R + 2R) + 2\lambda (C + DR) + E + FR = 0$$

$$R(\lambda^2 + 2\lambda + 1) = 0$$

$$\begin{aligned}
 \frac{\partial R}{\partial \alpha} &= -\frac{\partial p}{\partial \alpha} = 0 & \frac{\partial p}{\partial x} = 2\sqrt{x} \\
 & \text{I} & \text{II} & \text{III}
 \end{aligned}$$

$$\left[\frac{R_1}{\sqrt{1+p^2+q^2}} - 1 \right] - \left[\frac{R_2}{\sqrt{\quad}} - (1+p^2) \left[\frac{R_1}{\sqrt{\quad}} - (1+p^2) \right] \right] = 0$$

Wp HKR

$$\begin{aligned}
 \frac{1}{2} \frac{R^2}{1+p^2+q^2} (s=rt) + \frac{R}{\sqrt{\quad}} [-2pqs + (1+p^2)t + (1+q^2)r] \\
 + p\dot{p} - (1+p^2)(1+q^2) = 0 \\
 -1+p^2+q^2
 \end{aligned}$$

Wp HKR

$$I \cos \alpha + II \sin \alpha \cos \beta + III \sin \alpha \sin \beta = 0$$

$$\frac{\partial}{\partial \alpha} I \cos \alpha + II \sin \alpha \cos \beta = 0$$

$$\frac{\partial}{\partial \beta} II \sin \alpha \cos \beta + III \sin \alpha \sin \beta = 0$$

$$\frac{\sin \beta}{\cos \beta} = -\frac{I}{II} = -\frac{II}{III}$$

sec

Wp HKR

$$(1+p^2) \cos \alpha + pq \sin \beta - \frac{R}{\sqrt{\quad}} [s \cos \alpha + t \sin \beta] = 0$$

$$pq \cos \alpha + (1+q^2) \sin \beta - \frac{R}{\sqrt{\quad}} [s \cos \alpha + t \sin \beta] = 0$$

I $\frac{1}{R_1} + \frac{1}{R_2} = \frac{1}{R}$ $\Rightarrow R = \frac{R_1 R_2}{R_1 + R_2}$

II $\frac{1}{R_1} + \frac{1}{R_2} = \frac{1}{R}$ $\Rightarrow R = \frac{R_1 R_2}{R_1 + R_2}$
 $\Rightarrow R = \frac{1}{\frac{1}{R_1} + \frac{1}{R_2}}$

$$R_1 + R_2 = \frac{(1+\mu)t + (1+\mu)r - 2ps}{\mu t - s^2}$$

$$\frac{1}{R_1 R_2 (R_1 + R_2)}$$

$$\frac{1}{R_1} + \frac{1}{R_2} = \frac{(1+\mu)t + (1+\mu)r - 2ps}{[\mu t + s^2]^{\frac{1}{2}}}$$

für $\mu = 0$ gilt

Wird $\mu = 0$ $\Rightarrow s = 0$ \Rightarrow $\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$
 = Mittelwert $\mu = 0$ \Rightarrow $\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$
 Minimalfläche $\Rightarrow s = 0$ \Rightarrow $\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$
 $\Rightarrow \mu = 0$

Capillartät

116 $\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$ \Rightarrow $\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$



$0 = dp \cdot \eta + dq \cdot \gamma$ η, γ are the direction cosines

$dp \cos \alpha' + dq \sin \alpha' = 0$

$dp = -r \cos \alpha \, ds$

$dq = t \sin \alpha \, ds$

$r \cos \alpha \cos \alpha' + t \sin \alpha \sin \alpha' = 0$ = direction cosines
 $\gamma \sim \sin \alpha' \cos \alpha'$
 $\eta \sim \cos \alpha' \sin \alpha'$

line of intersection

are the lines $P \sim \eta, \gamma$ of the plane T_1 and T_2 respectively

are also lines.

the η, γ are the direction cosines of the lines P, R, Q, S etc.

are orthogonal lines in involutions.

the lines P, R, Q, S are the lines of intersection of the planes T_1, T_2

of the lines P, R, Q, S are the lines of intersection of the planes T_1, T_2

η, γ



the lines P, R, Q, S are the lines of intersection of the planes T_1, T_2

$\frac{d}{dt} \rho = \text{Conver. constant} \cdot \rho$ in $\frac{1}{2}$ range $\frac{1}{2} \rho \sim \frac{1}{2} \rho$
 C. 2. 58

$$r dx dx' + t dy dy' = 0$$

$$\frac{r \cos^2 \alpha + 2s \sin \alpha \cos \alpha + t \sin^2 \alpha}{\sqrt{1+t^2}} = \frac{r}{\sqrt{1+t^2}}$$

$$r dx dx' + s(dx dy' + dx' dy) + t dy dy' = 0$$

$\frac{dx}{dy} = \frac{dx'}{dy'}$
 $\frac{dx}{dy} = \frac{dx'}{dy'}$

$\cos \alpha = \frac{r}{\sqrt{1+t^2}}$

conj. Curve

$$d\rho = \omega \rho \left(\frac{1}{r} + \frac{1}{\rho} - \frac{1}{\rho} \right) \text{ conj. } \rho$$

$\frac{1}{\rho} \frac{d\rho}{dt} = f(\rho) \text{ Hard } T_{\rho} = \text{Hard } T_{\rho} \text{ C.}$

$\frac{1}{\rho} \frac{d\rho}{dt} = f(\rho)$

$$r \cos^2 \alpha + 2s \sin \alpha \cos \alpha + t \sin^2 \alpha = 0$$

$$r dx^2 + 2s dx dy + t dy^2 = 0$$

$\frac{dx}{dy} = \frac{dx'}{dy'}$

$\frac{dx}{dy} = \frac{dx'}{dy'}$

et de l'esp = 2 8 C. et de 2 H 7 1/2 W.

2 Systeme L. Hays + 1/2.

2nd ~ 1/2 1/2 1/2 1/2 ~ H 7 1/2 W

et de 2 ~ 1/2 1/2 1/2 1/2 ~ H 7 1/2 W

et de 2 ~ 1/2 1/2 

et de 2 ~ 1/2 1/2 1/2 1/2 ~ H 7 1/2 W

et de 2 ~ 1/2 1/2 ~ H 7 1/2 W

et de 2 ~ 1/2 1/2, la ampl, et de 2 ~ 1/2 1/2

$\frac{x}{a+b} + \frac{y}{c+d} = 1$ et de 2 ~ 1/2 1/2

et de 2 ~ 1/2 1/2 ~ H 7 1/2 W

• all of the Norm of \vec{r} in $\omega = \omega^0 C_1 \vec{e}_1 + \dots$
 • all of the \vec{e}_i of $\omega = \omega^0 C_1 \vec{e}_1 + \dots$
 • $\vec{e}_1, \dots, \vec{e}_n$ are \vec{e}_i in $\omega = \omega^0 C_1 \vec{e}_1 + \dots$
 • $\vec{e}_1, \dots, \vec{e}_n$ are \vec{e}_i in $\omega = \omega^0 C_1 \vec{e}_1 + \dots$

$\frac{1}{2} \omega^0 \omega^1 \omega^2 \omega^3 \omega^4 \omega^5 \omega^6 \omega^7 \omega^8 \omega^9 \omega^{10} \omega^{11}$

$$\xi = x + l u$$

$$\eta = y + l v$$

$$\zeta = z + l w$$

$$l^2 \xi \eta \zeta = l^2 (x + l u)(y + l v)(z + l w)$$

W.D.:

$$\xi = x + l u + dx + l du \quad \parallel \quad \xi \eta \zeta = \omega^0 \omega^1 \omega^2 \omega^3 \omega^4 \omega^5 \omega^6 \omega^7 \omega^8 \omega^9 \omega^{10} \omega^{11}$$

$$0 = (l' - l) u + dx + l' du$$

$$0 = (l' - l) v + dy + l' dv$$

$$0 = (l' - l) w + dz + l' dw$$

3×3 / 2×2 $\vec{v}_1, \vec{v}_2, \vec{v}_3$ Δ $\omega^0 \omega^1 \omega^2 \omega^3 \omega^4 \omega^5 \omega^6 \omega^7 \omega^8 \omega^9 \omega^{10} \omega^{11}$

$$\begin{vmatrix} dx & dy & dz \\ u & v & w \\ du & dv & dw \end{vmatrix} = 0$$

= " of - Diff. $\omega^0 \omega^1 \omega^2 \omega^3 \omega^4 \omega^5 \omega^6 \omega^7 \omega^8 \omega^9 \omega^{10} \omega^{11}$
 = Diff. $\omega^0 \omega^1 \omega^2 \omega^3 \omega^4 \omega^5 \omega^6 \omega^7 \omega^8 \omega^9 \omega^{10} \omega^{11}$



$u_2 = v_2 + w_2$
 $= x_2 \frac{dx}{x} + v_2 \sim u_2$

$$du = u_1 dx + u_2 dy + u_3 dz$$

$$0 = \frac{dx}{x} \left(\frac{l-l'}{l'} u_1 \right) + dy u_2 + dz u_3 + \frac{(l'-l)}{l'} u$$

$$0 = dx v_1 + dy \left(\frac{l'}{l} + v_2 \right) + dz v_3 + \frac{l'-l}{l'} v$$

$$0 = dx w_1 + dy w_2 + \left(\frac{l'}{l} + w_3 \right) dz + \frac{l'-l}{l'} w$$

$u < 2 - x_2$: u, v, w $\rightarrow dx dy dz$ $\sim \frac{l'-l}{l}$

$$0 = dx u + dy v + dz w$$

$\rightarrow 4$ homog. eq, $dx dy dz \frac{l'-l}{l} \sim \frac{l'-l}{l}$

$$\Delta = 0 \sim$$

1000 a c u o e f Curve r p Hkr Rad / 19.0

v e w l e p Δ p Hkr Rad. 6.

$$u v w = p q - 1$$

$$\begin{vmatrix} dx & dy & dz \\ p & q & -1 \\ dp & dq & 0 \end{vmatrix} = dq [p dz + dx] + dp [q dz + dy] = 0$$

$$\frac{dq}{q dz + dy} = \frac{dp}{p dz + dx} \quad dz = p dx + q dy$$

$$\frac{s dx + t dy}{(q + pr) dy + p q dx} = \frac{r dx + s dy}{(1 + r^2) dx + p q dy}$$

$$(q + pr) dy + p q dx \quad (1 + r^2) dx + p q dy$$

dx, dy n p r q e c. a. d. Hkr "

$$\frac{s dx + t dy}{(1 + q^2) dx + p q dy} = \dots$$

pr, Be q e Hkr e

u s. v r. u a p

Hkr. q "

u e Hkr Curve u e = p v e p - devel.

Olinda Rodrigues:



Rockwell C. e α

f.c. > f.d. < e.h. < f

s.g. α e N $a, b, c = \alpha, \cos, \sin$

Cond. ec. $x + Ra$

α e Curva

$$\frac{dx + d(Ra)}{a} = ds =$$

$$= \frac{dy + d(Rb)}{b} = \frac{dz + d(Rc)}{c}$$

$$\frac{dx + R da}{a} + dR$$

$$ds - dR = \frac{dx + R da}{a} = \frac{dy + R db}{b} = \frac{dz + R dc}{c}$$

$$dx + R da = a(ds - dR) \quad \left| \begin{array}{l} a \\ b \\ c \end{array} \right. \quad \begin{array}{l} = 0 \\ = 0 \\ = 0 \end{array}$$

$$dy + R db = b(ds - dR)$$

$$dz + R dc = c(ds - dR)$$

$$0 + 0 = ds - dR$$

$$ds = dR$$

f 3 Formulas no
F. e Rodrigues

$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \frac{1}{2} \int \frac{1}{r^3} dr + \dots$

$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$

$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$

$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$

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$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$

$\omega = \frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$

$f-x$	$w-y$	$r-z$	$\frac{1}{2} \int \frac{1}{r} dr + \frac{1}{2} \int \frac{1}{r^2} dr + \dots$
dx	dy	dz	
d^2x	d^2y	d^2z	

$= 0$

- double Fl

$f(x) = \frac{1}{2} T_2(x) = \frac{1}{2} (2x^2 - 1)$

$f(x) = \frac{1}{2} (2x^2 - 1)$

$e^{-\frac{1}{2}} \cos \frac{\pi}{2} x$, involutions \times ~~the~~ $\sin C$

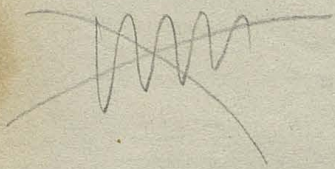
$[\frac{1}{2} \cos \frac{\pi}{2} x \sim \frac{1}{2} (2x^2 - 1) \text{ center } \frac{\pi}{2}]$

$\frac{1}{2} T_2(x) = \frac{1}{2} (2x^2 - 1)$

D.H. $\ln \sim \frac{1}{2} \ln 2 \cos \frac{\pi}{2} x$ und, $\frac{1}{2} T_2$ Curve

Thieren & Joachimsthal

M.H.



e.H. $\ln \frac{1}{2} \cos \frac{\pi}{2} x \sim \frac{1}{2} \ln 2$

$\frac{1}{2} \cos \frac{\pi}{2} x \sim \frac{1}{2} (2x^2 - 1)$ = invol. \wedge quad. Curve

$\frac{1}{2} \cos \frac{\pi}{2} x$ center $\frac{\pi}{2}$

$\frac{1}{2} \cos \frac{\pi}{2} x$

es

Tej. & involutions $\frac{1}{2} \cos \frac{\pi}{2} x$

z_3 zu z_2 } z_1 \in H.N. 64

$e \notin$ H.N. e ist $z_2 \in$ H.N. = $\tau + g$

$\nabla g|_f = 0 = \text{Tot. d. } \tau + g$

$\omega = z_3$ ∇ ω ∇ ω ∇ ω

ω ∇ ω ∇ ω ∇ ω ∇ ω

$\tau + g = 0$

$dt = 0$

Schnitt von 3 Flächen

$f(x,y,z) = \lambda$

$\varphi(x,y,z) = \mu$

$\psi(x,y,z) = \nu$

ω ∇ μ ∇ ω ∇ ω ∇ ω

ω ∇ ν ∇ ω ∇ ω ∇ ω

$x = f(\lambda, \mu, \nu)$

$y =$

$z =$



ω ∇ μ ∇ ω ∇ ω ∇ ω ∇ ω ∇ ω ∇ ω ∇ ω ∇ ω

$\frac{\partial x}{\partial \lambda}$ $\frac{\partial y}{\partial \lambda}$ $\frac{\partial z}{\partial \lambda}$

$$\frac{\partial x}{\partial \lambda} \frac{\partial x}{\partial v} + \frac{\partial y}{\partial \lambda} \frac{\partial y}{\partial v} + \frac{\partial z}{\partial \lambda} \frac{\partial z}{\partial v} = 0 \quad \text{or } \perp \text{ to } \vec{r}$$

or:

$$\frac{\partial x}{\partial \lambda} \frac{\partial x}{\partial v} + \frac{\partial x}{\partial \mu} \frac{\partial x}{\partial v} = 0$$

$$\frac{\partial x}{\partial \mu} \frac{\partial x}{\partial v} + \frac{\partial y}{\partial \mu} \frac{\partial y}{\partial v} = 0$$

$$\frac{\partial x}{\partial \lambda} \frac{\partial y}{\partial v} = \rho \left(\frac{\partial y}{\partial v} \frac{\partial z}{\partial \mu} - \frac{\partial z}{\partial v} \frac{\partial y}{\partial \mu} \right)$$

$$\frac{\partial y}{\partial \lambda} \frac{\partial x}{\partial v} = \rho \left(\frac{\partial z}{\partial v} \frac{\partial x}{\partial \mu} - \frac{\partial x}{\partial v} \frac{\partial z}{\partial \mu} \right)$$

$$\frac{\partial z}{\partial \lambda} \frac{\partial x}{\partial v} = \rho \left(\frac{\partial x}{\partial v} \frac{\partial y}{\partial \mu} - \frac{\partial y}{\partial v} \frac{\partial x}{\partial \mu} \right)$$

I $\frac{\partial x}{\partial \lambda}$ II $\frac{\partial y}{\partial \lambda}$ III $\frac{\partial z}{\partial \lambda}$:

$$\sum \frac{\partial x}{\partial \lambda} \frac{\partial x}{\partial v} + \frac{\partial x}{\partial \lambda} \frac{\partial z}{\partial v} = 0$$

$$\sum \frac{\partial x}{\partial \lambda} \frac{\partial x}{\partial v} + \frac{\partial x}{\partial \lambda} \frac{\partial y}{\partial v} = 0$$

$$\sum \frac{\partial x}{\partial \lambda} \frac{\partial x}{\partial v} + \frac{\partial x}{\partial \lambda} \frac{\partial z}{\partial v} = 0$$

$$1 + 2 - 3 = 0$$

$$\sum \frac{\partial x}{\partial \lambda \partial \mu} \frac{\partial x}{\partial \nu} = 0$$

$$\sum \frac{\partial x}{\partial \lambda \partial \nu} \frac{\partial x}{\partial \mu} = 0$$

$$\sum \frac{\partial x}{\partial \mu \partial \nu} \frac{\partial x}{\partial \lambda} = 0$$

orthogonale Vektoren
 x_1, x_2, x_3
 65

orthogonale Vektoren x_1, x_2, x_3

orthogonale Vektoren x_1, x_2, x_3

$$\begin{vmatrix} dx & dy & dz \\ u & v & w \\ du & dv & dw \end{vmatrix} = 0$$

Polare \sim Vektoren μ

$$\begin{vmatrix} \frac{\partial x}{\partial \mu} & \frac{\partial y}{\partial \mu} & \frac{\partial z}{\partial \mu} \\ \frac{\partial x}{\partial \nu} & \frac{\partial y}{\partial \nu} & \frac{\partial z}{\partial \nu} \\ \frac{\partial x}{\partial \mu \partial \nu} & \frac{\partial y}{\partial \mu \partial \nu} & \frac{\partial z}{\partial \mu \partial \nu} \end{vmatrix} = 0$$

$$\begin{cases} du = u \, d\lambda \\ dv = v \, d\lambda \\ dw = w \, d\lambda \end{cases}$$

orthogonale Vektoren x_1, x_2, x_3 : orthogonale Vektoren

orthogonale Vektoren

orthogonale Vektoren

orthogonale Vektoren : orthogonale Vektoren

Ab.

x, y, z - Rotations

3/Orthog. e^2



Meridian & Parallel e^2, y, z

$$\frac{x^2}{a^2 + \lambda} + \frac{y^2}{b^2 + \lambda} + \frac{z^2}{c^2 + \lambda} = 1$$

λ (constant) $2\lambda - \lambda$

$$\rightarrow \lambda \rightarrow +\infty$$

$\omega \sim \omega^2$, x, y, z same as $\omega, \omega^2, \omega^3$

$\times 3/4$ of ω^2 ; ω^2 of ω of ω^2

$\frac{1}{2}$ in. etc. ω^2 μ

$$\frac{x}{a + \lambda} \quad \frac{y}{b + \lambda} \quad \frac{z}{c + \lambda}$$

$$\frac{x}{a + \mu} \quad \frac{y}{b + \mu} \quad \frac{z}{c + \mu}$$

$$\frac{x^2}{(a^2-\lambda)(c^2-\lambda)} + \frac{y^2}{(b^2-\lambda)(c^2-\lambda)} + \frac{z^2}{(c^2-\lambda)(c^2-\lambda)} = 0$$

$$\text{resp: } \frac{x^2}{a^2-\lambda} + \frac{y^2}{b^2-\lambda} + \frac{z^2}{c^2-\lambda} = 1$$

$$(a^2-\lambda) \left[\frac{x^2}{(a^2-\lambda)(c^2-\lambda)} + \frac{y^2}{(b^2-\lambda)(c^2-\lambda)} + \frac{z^2}{(c^2-\lambda)(c^2-\lambda)} \right] = 0$$

$$\therefore P \perp p$$

$f(\lambda) =$

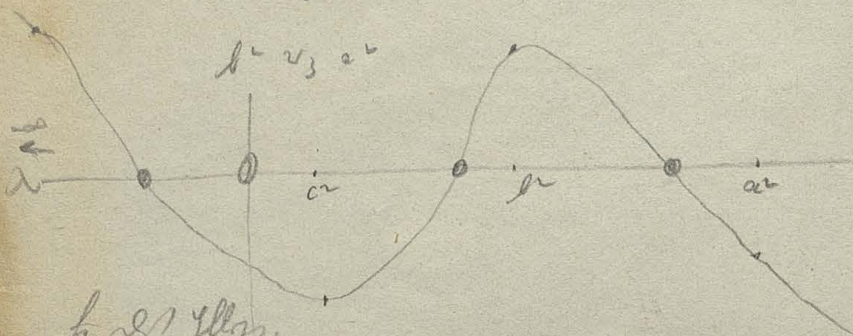
$$(a^2-\lambda)(b^2-\lambda)(c^2-\lambda) - (b^2-\lambda)(c^2-\lambda)x^2 - (a^2-\lambda)(c^2-\lambda)y^2 - (a^2-\lambda)(b^2-\lambda)z^2 = 0$$

put f as λ^3 zero:

$$-\infty \quad v_1 \quad c^2$$

$$a^2 \quad v_2 \quad b^2$$

$$b^2 \quad v_3 \quad a^2$$



f is λ^3 zero.

$$\frac{x^2}{a^2-k} + \frac{y^2}{b^2-k} + \frac{z^2}{c^2-k} = 1$$

\hookrightarrow 36 d 22 Ge Inter. 9 10 10 36 d

λ_1

λ_2

λ_3

λ_4

λ_1 - Ellipsoide

λ_2 einscheligs Hyperboloid

λ_3 zweiseheligs Hyperboloid

[$\lambda_2 > \lambda_3$ oder $\lambda_2 < \lambda_3$ null. Platte]

Wp $\lambda_1 = c^2$

\hookrightarrow 21 p $\lambda_1 = c^2$

λ_1 Ellipsoide $\hookrightarrow \lambda = c^2$ Wp $c^2 = 0$

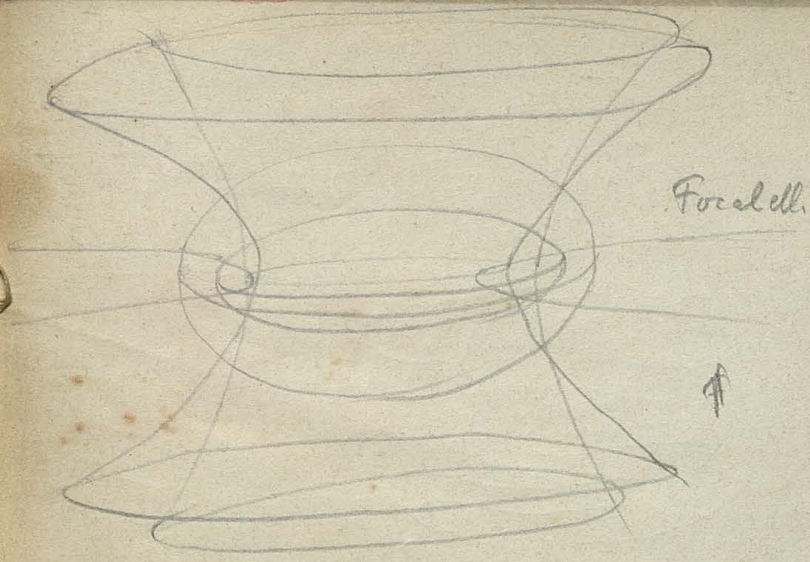
= XY Ebene

$$\frac{x^2}{a^2-c^2} + \frac{y^2}{b^2-c^2} = 1 \quad \text{---} \quad \text{Ellipsoid 16}$$

Focal-Ellipse

\hookrightarrow en $\lambda_2 = c^2$ null

\hookrightarrow $\frac{x^2}{a^2-c^2} + \frac{y^2}{b^2-c^2} = 1$ \hookrightarrow $\lambda = c^2$ Ellipsoid

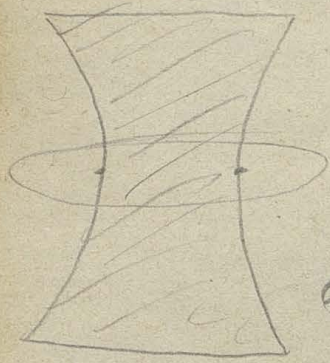


Focalell.

$\angle \lambda_2 = n b^2 d$

$\approx \frac{x^2}{a^2 - b^2} + \frac{z^2}{c^2 - b^2} = 1$

= Hyperbol. 1. sheet



$\angle \text{em } \lambda_3 = b^2 = 0$

ess of Hyper

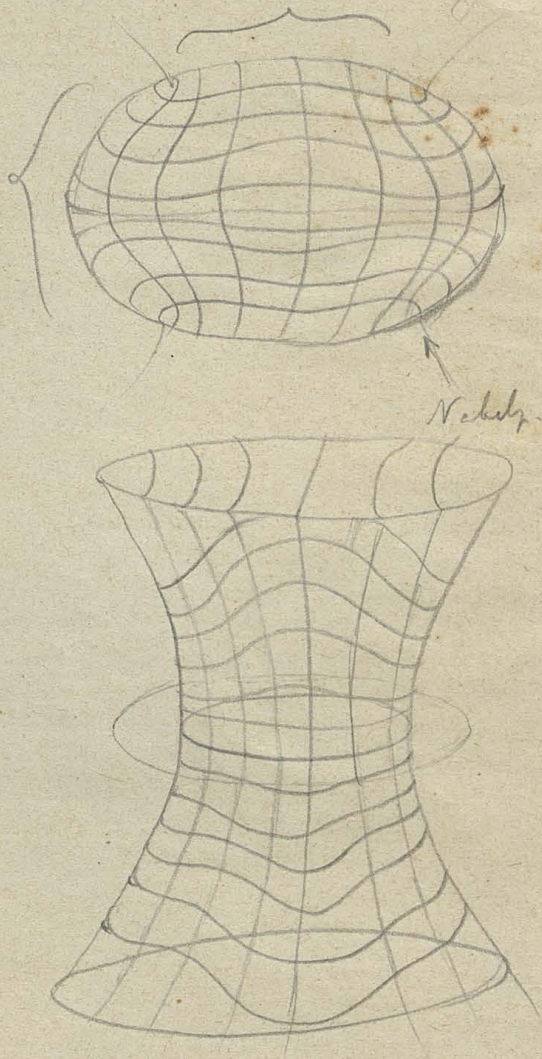
Focally hyperbol

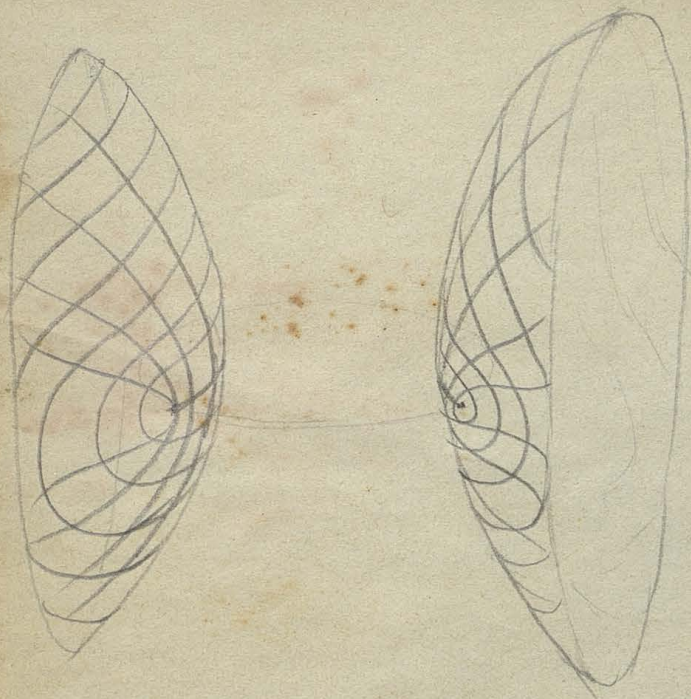
$\omega \lambda_3 n a^2 = 0$ on \mathbb{R}^3 \rightarrow image

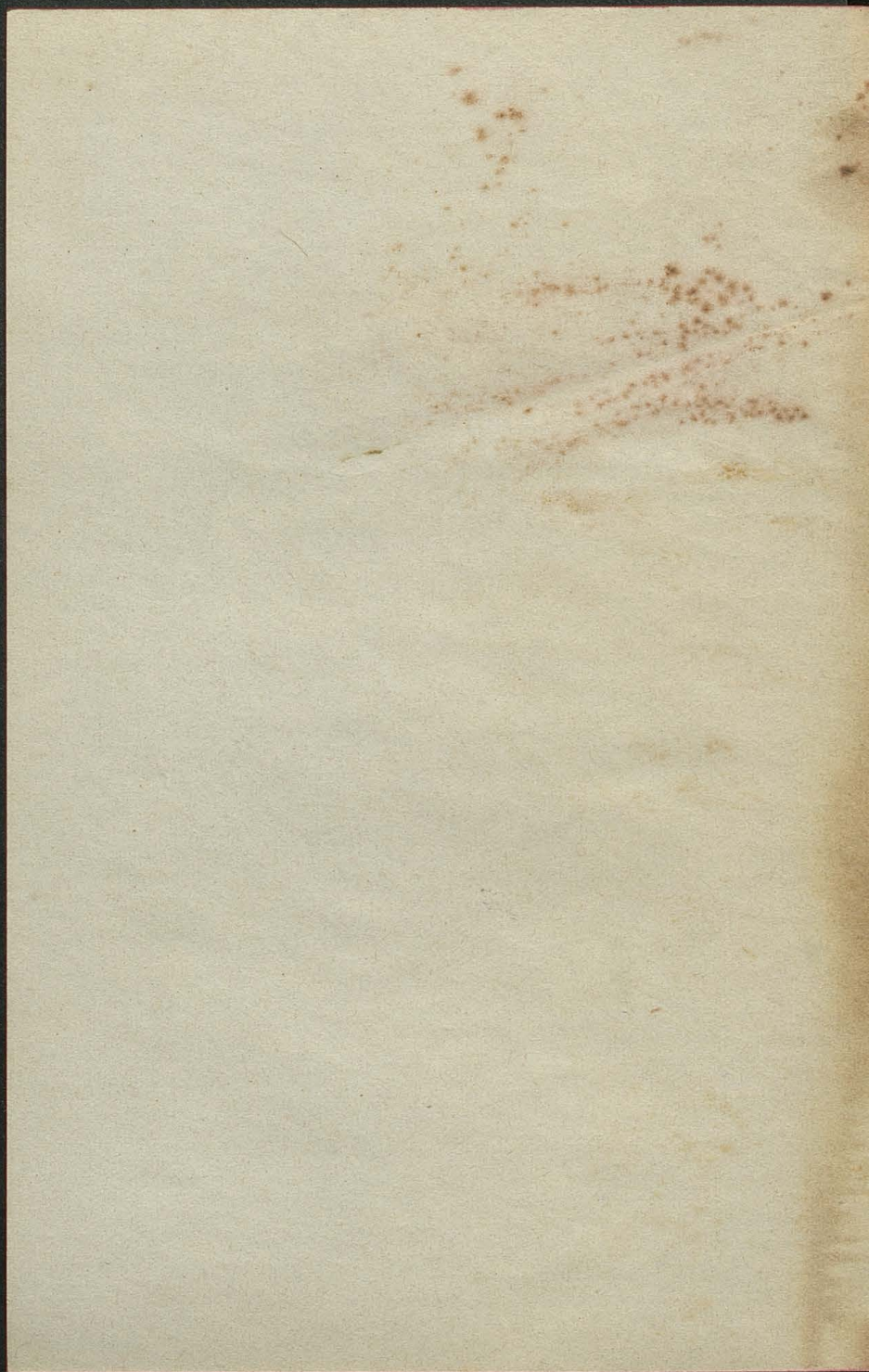
$\vee \lambda_2 = n b^2 d$

$\lambda_1 = 0$ on \mathbb{R}^3

$1 \text{ ell. } \sim \angle \text{ re } 1/2 \text{ } : \text{ ce } \text{ anta } \text{ } \sim$
 $< \text{ u ell. } \text{ c } 1/2 \text{ } \sim 2 \varphi \text{ sp. } \sim 2 \varphi \checkmark$







69

C. Neumann:

[Faint handwritten notes, partially obscured by a large dark stain]

$$+ (\delta H + \sum P \delta p) dt = 0 \quad \frac{\partial H}{\partial t}$$

$$\therefore \frac{\partial}{\partial t} \left(\frac{\partial H}{\partial t} + P_i \right) \text{ all. Lagrange's } \checkmark$$

no. e.g. $\frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{y}^2$ $U = mgy + \text{const}$

$$H = T - U \quad H = \text{kinet. Potential}$$

$$\therefore \frac{\partial}{\partial t} (m\dot{x}) = X \quad \frac{\partial}{\partial t} (m\dot{y}) = -mgy + Y$$

Master, des Principes de l'énergie

$$j_i \frac{\partial}{\partial t} \left(\frac{\partial H}{\partial j_i} \right) - j_i \frac{\partial H}{\partial j_i} = j_i P_i \quad \text{sur } P_i \text{ et } U_i$$

$$\sum \left[\frac{d j_i}{dt} \frac{d}{dt} \left(\frac{\partial H}{\partial j_i} \right) - \frac{d j_i}{dt} \frac{\partial H}{\partial j_i} \right] = \sum$$

$$\frac{dH}{dt} = \sum \frac{\partial H}{\partial j_i} \frac{\partial j_i}{\partial t} + \sum \frac{\partial H}{\partial p_i} \frac{\partial p_i}{\partial t}$$

$$\sum \frac{d j_i}{dt} \frac{d}{dt} \left(\frac{\partial H}{\partial j_i} \right) + \frac{\partial H}{\partial t} = \sum j_i P_i$$

$$\sum \left[\frac{d}{dt} \left(j_i \frac{\partial H}{\partial j_i} \right) - \frac{dH}{dt} \right] =$$

$$= \frac{d}{dt} \left[\sum j_i \frac{\partial H}{\partial j_i} - H \right] = \sum P_i d p_i$$

= \mathcal{E} ✓ e w

$$T = \text{homog. } \mathcal{L} \text{ de } f_i \quad \therefore \sum x \frac{\partial T}{\partial x} = 2T$$

ou N et q_i stat.

$$E = 2T - H \approx 2T - T + U = T + U$$

= \sum énergie pot. Energie

1. E est une fonction de H , e est une fonction de H

2. E est une fonction de H , e est une fonction de H

t. Gleichgewichtszustand

$$\int (T - Ut \sum P \delta p) dt = 0$$

$$to = e \text{ etc } = 0 \text{ etc } \delta T = 0 \quad \delta U = \sum P \delta p \text{ [Quadr.]}$$

Einführung einer neuen Variable in Legendre Gl.

$$t' = const - t$$

$$\therefore dt' = - dt \quad p' = - p$$

$$H' = H \text{ const } \text{ etc}$$

$$+ \frac{d}{dt'} \left(\frac{dH}{dt'} \right) = \frac{dH}{dp} \quad \text{etc}$$

$$\text{etc} \sim \text{etc}$$

$$\text{etc } A' = t_0 t_1 - t \quad \text{etc}$$

etc reversible in etc

etc in etc

etc etc [etc, etc]

etc etc

$$\text{etc} \frac{d^2}{dt^2} \text{ etc}$$

$$< 0 \text{ etc} \frac{d}{dt} \text{ etc}$$

etc etc

f was μ , $\rho \delta^2$ in $[c \sqrt{v}]$

$v \sim \sqrt{gpc} / \omega \sim \sqrt{P} \cdot \sqrt{g} / \omega$

$\frac{d}{dt} \left(\frac{\partial H}{\partial \dot{x}_i} \right) = \frac{\partial H}{\partial x_i} + P$ f. r. l. e. v. a. H. d. d. l. e.
P Dispersion 100%

$c^2 \sim \mu \sim \rho \omega^2 \sim \rho \omega^2 \cdot \lambda^2$

$v \sim \omega \sim 10^8 \text{ rad/s}$

f. in. $\omega \sim \mu \omega^2 \sim \rho \omega^2 \cdot \lambda^2$

Elastische Vorgänge

Ableitung der Bewegungsgl. / Kristall

$\omega \sim v \sim \text{contin. Bedingung}$, $\omega \sim \text{Anzahl der } \omega$

$\oint \left(\delta H + \sum_i P \delta p_i \right) dt = 0$

weil δp_i , δp_i v. d. e.

$H = T - U$ $T = \sum \mu \dot{x}_i^2 / 2$

$\begin{matrix} x \\ y \\ z \end{matrix} \quad \begin{matrix} u \\ v \\ w \end{matrix} \quad \left. \begin{matrix} \text{displacement} \\ \text{components} \end{matrix} \right\} \quad \begin{matrix} i \\ j \\ k \end{matrix} \quad \left. \begin{matrix} \text{direction} \\ \text{components} \end{matrix} \right\}$

$T = \frac{\rho}{2} \int (\dot{u}^2 + \dot{v}^2 + \dot{w}^2) dt$

$U = \text{pot. En.}$, $\mu \omega^2 \sim \rho \omega^2 \cdot \lambda^2$

P kinematische f. l. e. $\sim \rho \omega^2 \cdot \lambda^2$

$$\sum P_i \delta r_i = \int ds (X_x \delta u + Y \delta v + Z \delta w)$$

$$\int_{t_0}^{t_1} dt \left\{ \int ds h (i \delta u + \ddot{v} \delta v + \ddot{w} \delta w) - \int ds \delta f + \dots \right.$$

$$\left. \begin{aligned} \delta (d \dot{u}) &= d \dot{u} \delta t + \dot{u} \delta ds + \dots \\ &\dots \end{aligned} \right\}$$

$$+ \int ds \{ X_u \delta u + Y_v \delta v + Z_w \delta w \} = 0$$

in die part. int.

$$\int_{t_0}^{t_1} dt \left\{ \int ds (i \delta u + \ddot{v} \delta v + \ddot{w} \delta w) - \int ds \left[\frac{\partial f}{\partial x_x} \delta x_x + \dots \right] \right.$$

$x_x = \frac{\partial u}{\partial x}$ zur Abkürzung: $-\frac{\partial f}{\partial x_x} \text{ etc.} = X_x \text{ etc.}$

$$\int ds \left[X_x \delta \frac{\partial u}{\partial x} + \dots \right] = \int ds \left[\frac{\partial X_x}{\partial x} \delta u + \dots \right]$$

$$+ \left(\frac{\partial X_y}{\partial y} \delta u + \frac{\partial X_x}{\partial x} \delta v \right) + (\dots) + (\dots) \int ds$$

$$= \int ds \left[X_x \cos \alpha_x \delta u + Y_y \cos \alpha_y \delta v + Z_z \cos \alpha_z \delta w + \dots \right]$$

$k \ddot{u} = \dots$

$$k \ddot{u} + \frac{\partial X_x}{\partial x} + \frac{\partial X_y}{\partial y} + \frac{\partial X_z}{\partial z} = 0$$

$$k \ddot{v} + \frac{\partial Y_x}{\partial x} + \frac{\partial Y_y}{\partial y} + \frac{\partial Y_z}{\partial z} = 0$$

$$k \ddot{w} + \dots = 0$$

$e^{t \dots}$
 Elong

$$X_v = X_x \cos \alpha_x + X_y \cos \alpha_y + X_z \cos \alpha_z$$

$$Y_v = Y_x \cos \alpha_x + Y_y \cos \alpha_y + Y_z \cos \alpha_z$$

$$Z_v = \dots$$

$f(x, y, z) = \dots$

$\partial X_v \dots$

$f(x, y, z) \dots$

Sem.

$f(x, y, z) \dots$

t_0, x_0, y_0
 t_1, x_1, y_1

$$H = \frac{m}{2} (\dot{x}^2 + \dot{y}^2 - mgy)$$

$\int dt H \dots$

$$\begin{aligned}
 t_0 = 0 & \quad t_1 = 1 & \quad x_0 = y_0 = 0 \\
 & & \quad x_1 = 1 \quad y_1 = 0
 \end{aligned}$$

$p, r \sim \omega_2 \sim p$ Drehung; $y = r \sim e^{i\theta}$

$x = t \quad y = \frac{1}{2}gt(1-t)$

v am 90° ω_2 für $0 \leq t \leq 1$ ω_2 ω_1 ω_2

od. $x = vt$

$y = vt^2$

v ω_2 ω_1 ω_2

$\left\{ \begin{array}{l} \omega_1 \sim \omega_2 \sim \omega_1 \\ \text{at } \text{rot. } \omega_1 \sim \omega_2 \sim \omega_1 \end{array} \right.$

Beugung eines Rotations Körpers

od. = Cardan'scher Suspension

1 ω vertical α

1 ω horizontal β

of ω_1, ω_2 - horiz. ω_3 ω_1, ω_2



$\omega_3 = 0 = 2\omega_2$ ω_1

$H = T = \frac{1}{2} [P\dot{\alpha}^2 + Q\dot{\beta}^2 + R\dot{\gamma}^2]$ $\omega_1, \omega_2, \omega_3$

$P = \sin\alpha \sin\beta \dot{\alpha} + \cos\beta \dot{\beta}$ comp. e $(\sin\alpha \dot{\alpha} \dot{\beta} \dot{\gamma})$

$Q = \sin\beta \cos\alpha \dot{\alpha} - \sin\alpha \dot{\beta}$

$R = \cos\beta \dot{\alpha} + \dot{\gamma}$

$e^{i\theta} \sim \omega_1 \sim \omega_2 \sim \omega_1$ $\omega_1, \omega_2, \omega_3$

$H = \frac{1}{2} Q (\sin\beta \dot{\alpha}^2 + \dot{\beta}^2) + R (\cos\beta \dot{\alpha}^2 + \dot{\gamma}^2 + 2\cos\beta \dot{\alpha} \dot{\gamma})$

$$\frac{d}{dt} \left(\frac{\partial H}{\partial \dot{q}_i} \right) - \frac{\partial H}{\partial q_i} = P$$

Asa $\sqrt{a^2 + b^2}$

$$\begin{cases} \frac{d}{dt} [P \sin^2 \beta \dot{\alpha} + R \cos \beta (\cos \beta \dot{\alpha} + \dot{\beta})] = A \\ \frac{d}{dt} [Q \dot{\beta}] - \{ \sin \beta \cos \beta Q \dot{\alpha}^2 - R (\cos \beta \dot{\alpha} + \dot{\beta}) \sin \beta \dot{\alpha} \} = B \\ \frac{d}{dt} [R (\cos \beta \dot{\alpha} + \dot{\beta})] = \Gamma \end{cases}$$

Euler'sche'sche $\beta, \alpha, \gamma, \delta < 1$ per δ

$\sqrt{a^2 + b^2} + \sqrt{c^2 + d^2} = A \sin \theta + B \cos \theta$
 $< \sqrt{a^2 + b^2} + \sqrt{c^2 + d^2} = \sqrt{a^2 + b^2 + c^2 + d^2}$

$$\therefore \Gamma = 0$$

$$R (\cos \beta \dot{\alpha} + \dot{\beta}) = \text{const} = c$$

$$\dot{\beta} = \frac{c}{R} - \cos \beta \dot{\alpha}$$

$$\begin{cases} \frac{d}{dt} [P \sin^2 \beta \dot{\alpha} + c \cos \beta] = A \\ \frac{d}{dt} [Q \dot{\beta}] - [Q \dot{\alpha}^2 \sin \beta \cos \beta - c \sin \beta \dot{\alpha}] = B \end{cases}$$

$$e^{\gamma} = P \sin^2 \beta \dot{\alpha} + \frac{c}{R} \cos \beta$$

$L < P \sin^2 \beta \dot{\alpha} + \frac{c}{R} \cos \beta < \sqrt{L^2 + c^2}$

$\therefore \sqrt{L^2 + c^2} \leq \sqrt{L^2 + c^2} \leq \sqrt{L^2 + c^2}$

$$A = \frac{d}{dt} [P \sin^2 \beta \dot{\alpha} + R \cos \beta \dot{\alpha}]$$

$$B = \frac{d}{dt} [Q \dot{\beta}] - \sin \beta \cos \beta [P \dot{\alpha}^2 + R] \dot{\alpha}$$

$\sigma_j - f_j(f)$ or $\pi_j \pi_j^2$

occure \rightarrow as $\pi_j \sim y / H$

$$\frac{\partial \mathcal{L}}{\partial \pi_j} = \frac{\partial H}{\partial \pi_j} + \sum_i \frac{\partial H}{\partial \pi_i} \frac{\partial \pi_i}{\partial \pi_j}$$

"c, etc."

$$= \frac{\partial}{\partial \pi_j} [H + \sum c_i \pi_i']$$

$$\frac{\partial \mathcal{L}}{\partial \pi_j} = \frac{\partial H}{\partial \pi_j} + \sum \frac{\partial H}{\partial \pi_i} \frac{\partial \pi_i}{\partial \pi_j}$$

$$= \frac{\partial}{\partial \pi_j} [H + \sum c_i \pi_i']$$

$$H - \sum c_i \pi_i' = K = f(\pi_j, \pi_i)$$

$$\left(\frac{\partial H}{\partial \pi_j} \right)_{\pi_i'} = \left(\frac{\partial K}{\partial \pi_j} \right)_{\pi_i'}$$

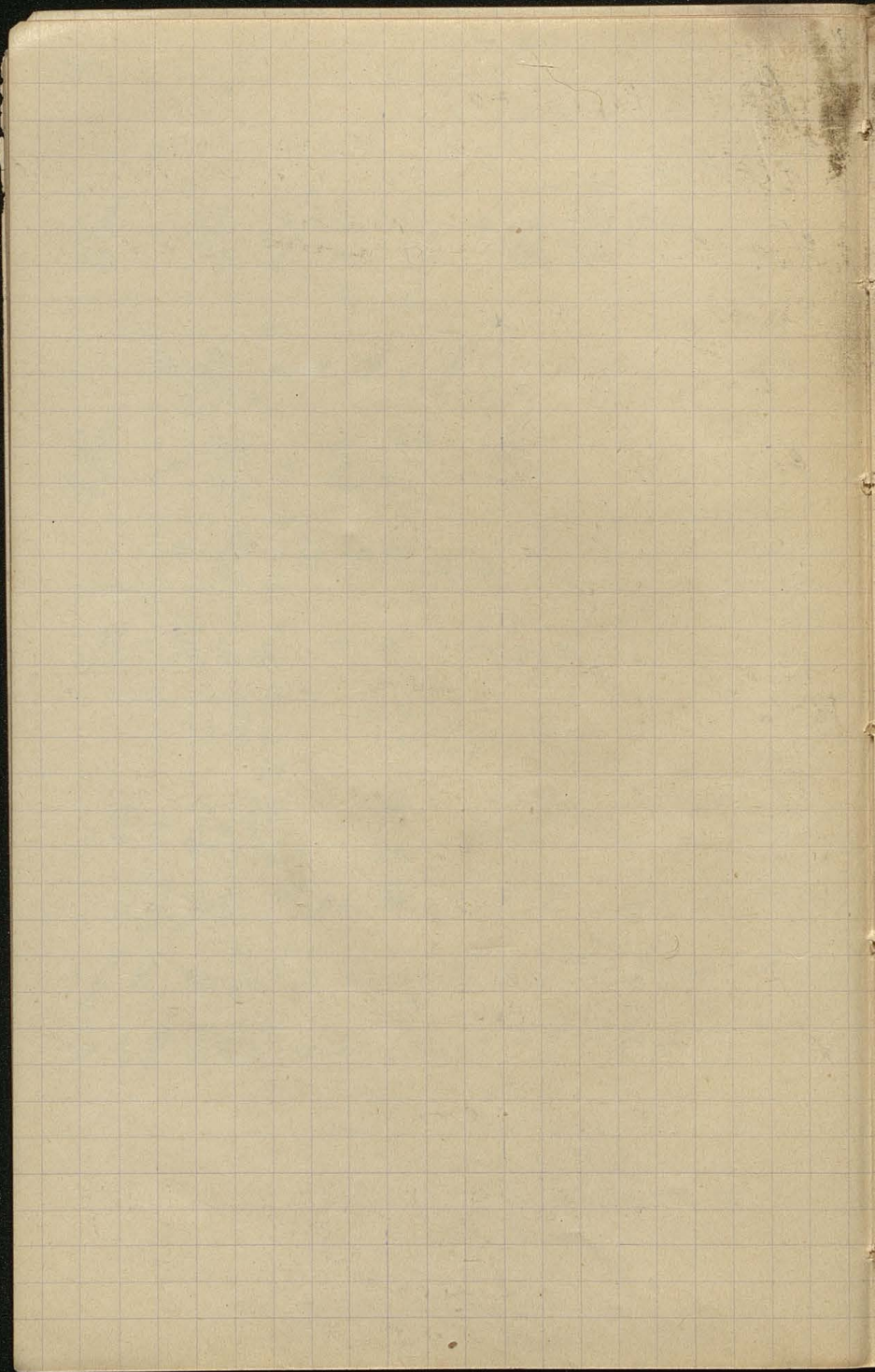
$$\therefore P = \frac{\partial}{\partial \pi_j} \left(\frac{\partial K}{\partial \pi_j} \right) - \frac{\partial K}{\partial \pi_j}$$

x effe... / ...

$\sim \int \pi_j \times \dots$

f K = a n / a honey ...

... [P ...]



$$\int (\delta H + \sum P \delta p) dt = 0$$

to 2. p i n f p - f l y o c a s t e s p h i n g d e l
of mlt. c - n s y u v l e s ~~s p h i~~ o e + L M N

o p u v r e d c n l ; c o y e d p r z v o b l o f =

20.:

$$L = u \quad M = v \quad N = w$$

$$X = \frac{c}{k} \left(\frac{\partial u}{\partial y} - \frac{\partial v}{\partial z} \right)$$

$$Y = \frac{c}{k} \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

$$Z = \dots$$

[6 ~ 10. 13 ~ 1]

$$\delta H = \delta T - \delta U \quad \sum P \delta p = \sum L \delta u + \sum V \delta v + \dots$$

u p c p e p h y - } f o w
v p c s r s i x c : =
u = u_0

$$\int_{t_0}^{t_1} \delta T dt = \int_{t_0}^{t_1} \frac{\mu}{4\pi} [L \delta L + M \delta M + N \delta N] =$$

$$= \int_{t_0}^{t_1} \frac{\mu}{4\pi} [u \delta u + \dots]$$

$$= - \int_{t_0}^{t_1} \frac{\mu}{4\pi} \left[\frac{\partial L}{\partial t} \delta u + \frac{\partial M}{\partial t} \delta v + \frac{\partial N}{\partial t} \delta w \right]$$

$$+ \int_{t_0}^{t_1} [L \delta u + M \delta v + \dots]$$

$$c \delta u = \delta v = \delta w = 0$$

$$\begin{aligned}
 \int_{t_0}^{t_1} \delta U dt &= - \int_{t_0}^{t_1} \frac{k}{4\pi} [X \delta X + \dots] dt \\
 &= - \frac{c}{4\pi} \int_{t_0}^{t_1} [X \delta \left(\frac{\partial v}{\partial y} \right) + \dots] dt \\
 &= - \frac{c}{4\pi} \int_{t_0}^{t_1} \left[\left(\frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x} \right) \delta w + \dots \right] dt \\
 &\quad + \frac{c}{4\pi} \int_{t_0}^{t_1} [X (\delta v \cos \gamma - \delta v \cos \alpha) + \dots] dt
 \end{aligned}$$

$$\begin{aligned}
 \therefore \frac{\partial L}{\partial t} &= c \left(\frac{\partial Y}{\partial x} - \frac{\partial Z}{\partial y} \right) \\
 \frac{\partial L}{\partial x} &= c \left(\frac{\partial Z}{\partial x} - \frac{\partial X}{\partial z} \right)
 \end{aligned}$$

in the presence of a medium ...
 the propagation of light ...
 the velocity of light ...

$$\begin{aligned}
 \text{the velocity of light in a medium} &= \frac{c}{n} = \frac{c}{\sqrt{\epsilon \mu}} \\
 \text{where } \epsilon &= \epsilon_0 + \epsilon_p \quad \mu = \mu_0 + \mu_p
 \end{aligned}$$

of medium ...

$$\begin{aligned}
 \text{Poynting vector } \mathbf{S} &= \mathbf{E} \times \mathbf{H} \\
 U_r &= - \frac{c}{4\pi} [Y \cos \alpha - Z \cos \gamma] \\
 V_r &= \\
 W_r &=
 \end{aligned}$$

o f UVW $\epsilon > c \alpha / \delta$; $c \epsilon \alpha + c \epsilon \beta \delta / \alpha \epsilon$

Energie $\int_0 (U_r \delta u + V_r \delta v + W_r \delta w) d\delta dt$ 77

$$\frac{c}{4\pi} d\delta dt [U_r \dot{u} + \dots]$$

$$= i (Z_{er} v_y - Y_{er} v_z)$$

$$+ \dot{v} (\dots) + \dot{w} (\dots)$$

$$= L (Z_{er} v_y - Y_{er} v_z) + \dots$$

$$= \text{Energiefluss} \quad (\alpha \epsilon \gamma \leq c)$$

V_0 α -spe. Integrale u, v, w δt ; $\alpha \epsilon \gamma$ δt / δt $\alpha \epsilon \gamma$ c

f. l. $\alpha \epsilon \gamma \frac{\partial u}{\partial y} - \frac{\partial v}{\partial z}$ etc. δt . $\beta \sim c \gamma \leq 3 \gamma \alpha \epsilon / \alpha \epsilon$

17/5 Spezielles Beispiel: Endliche Anzahl von linearen

Leetern, welche $\sim C_r \delta_{\gamma 0}^{\epsilon} \sim C_r \delta^{\epsilon}$ (δ, δ^2 etc.)

p. Var. $\delta, \delta^2, \delta^3$ \dots $\gamma_1, \gamma_2, \gamma_3$

$\delta \sim C_r \delta$ $\epsilon, \epsilon_2, \epsilon_3$ [var. E. w. $\delta \delta^2 \delta^3$ $\delta_{\gamma 0}^{\epsilon}$]

$$H = T - U = E$$

$$\frac{1}{2} [L_1 \dot{\epsilon}_1^2 + L_2 \dot{\epsilon}_2^2 + \dots] + L_{12} \dot{\epsilon}_1 \dot{\epsilon}_2 + \dots + T$$

$$L_i = \frac{1}{2c} \iint \frac{d\alpha d\alpha'}{n} \cos \theta$$

f. l. ρ δ ϵ γ δ

$$\mathcal{L} \sim P_1, P_2 \dots E, E_n \sim m ; U_p^c = 0$$

$$P = \frac{\partial}{\partial \dot{x}} \left(\frac{\partial H}{\partial \dot{x}} \right) - \frac{\partial H}{\partial x}$$

$$= \frac{\partial T}{\partial \dot{x}_i} - \frac{\partial T}{\partial x_i} - \frac{1}{2} \left[L_{i1} \dot{x}_1 + L_{i2} \dot{x}_2 + \dots \right]$$

$$\therefore \text{Induktionsterm}$$

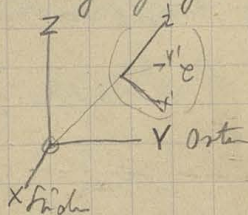
$$E_i = \frac{\partial E}{\partial \dot{x}_i} = \frac{\partial}{\partial \dot{x}_i} \left[L_{i1} \dot{x}_1 + L_{i2} \dot{x}_2 + \dots \right]$$

= Induktionsterm

$\sim E_i \sim U_p \text{ für } n/a \text{ } \epsilon_0^c \text{ } 2 \text{ } \text{in } \dot{x} \text{ } P$

el. mit. /

Bewegungsgleichg., $\sim \mathcal{L}$ \rightarrow rotierender Körper



$\sim \mathcal{L}$ - 4te Coord. $\dot{\varphi}$

ϵ_2 $\text{in } 2 \text{ } \mathcal{L}$

$$\frac{1}{2} \mathcal{L} \dot{\omega}^2 + \frac{m}{2} [\dot{x}'^2 + \dot{y}'^2 + \dot{z}'^2]$$

$$\dot{x}' = \dot{x} - y \sin \varphi \dot{\omega}$$

$$\dot{y}' = \dot{y} [R + r \sin \varphi + z \cos \varphi] \dot{\omega}$$

$$\dot{z}' = \dot{z} - y \cos \varphi \dot{\omega}$$

$$= \frac{1}{2} m$$

$$\Omega = \frac{d}{dt} \left(\frac{\partial H}{\partial \dot{u}} \right)$$

$f \cdot p / \cos \alpha \sin \alpha < \gamma \cos \alpha \tan \alpha$
const. \dot{u} ; $v \cdot p \tan \alpha / \cos$

oder $\cos \alpha \sin \alpha \cdot 2 \rho \sin \alpha \cdot \cos \alpha \cdot \dot{u} \cdot \sin \alpha$

$$\Omega = x'Y' - y'X' = (xY - yX) \sin \rho + [(R+Z)Y - yZ] \cos \rho$$

$$X = m \left[\ddot{x} - \frac{\partial}{\partial t} (y \sin \rho) \dot{u} \right] = m \ddot{x} - 2 \dot{y} \sin \rho \dot{u} - \rho \dot{u}^2 \sin \rho$$

$$Y = m \ddot{y} - \frac{\partial}{\partial t} (\rho) \cdot \dot{u} \sin \rho = m \ddot{y} - 2 \dot{x} \sin \rho \dot{u} + 2 \dot{z} \cos \rho \dot{u} - y \dot{u}^2$$

$$Z = \dots = m \ddot{z} - 2 \dot{y} \cos \rho \dot{u} - \rho \dot{u}^2 \cos \rho$$

oder $v \cdot p / \cos \alpha$ eliminieren $\sin \alpha \cdot \dot{u}$

$$\frac{\partial \Omega}{\partial t} = 0$$

$$\Omega = \text{const} \quad f \cdot \text{ob} \cdot \dot{u} \cdot \cos \alpha \cdot \sin \alpha / \cos$$

20. $\gamma \cdot g = Z = g$

$X = Y = 0$
und $\cos \alpha \cdot \rho \cdot g \cdot \sin \alpha = \sin \alpha \cdot m$

$$\ddot{x} - R \dot{u}^2 \sin \rho = 0$$

$$\ddot{y} - \dot{y} \dot{u}^2 = 0$$

$$\ddot{z} - R \dot{u}^2 \cos \rho = g$$

von f und \dot{u} $\rho = \frac{1}{2} \cos \alpha \cdot \sin \alpha \cdot 2 \rho \sin \alpha \cdot \dot{u} / \cos$

in $\dot{u} \cdot \sin \alpha$, $\dot{u} \cdot \cos \alpha$ $\dot{y} \cdot \sin \alpha$

$$g = \ddot{x} - 2 \dot{y} \sin \rho \dot{u}$$

$$0 = \ddot{y}$$

$$-g = \ddot{z}$$

$\dot{u} \cdot \sin \alpha \cdot \dot{u} \cdot \cos \alpha$

$$\therefore x = 0$$

$$z = \frac{1}{2} \frac{g t^2}{\sin \alpha}$$

$$y = \frac{1}{3} \dot{u} \sin \rho g t^3$$

$$dE = d\phi = P dr$$

$$\frac{d\phi}{\sigma} = dS$$

$$17 \quad P C = P H X \quad \sigma \quad \text{Vol} - G P Z \quad \sigma = \text{Vol}$$

$$\frac{Jt}{\epsilon} = N + N' \quad \text{if } J \text{ in } g \text{ or } e \text{ Concentr. } \epsilon A$$

a) $AgCl$ $\times 6 \times$ Conc. $\times 2 \text{ M}$; ϵNO_3 9×10^8
 b) $\times 2 \text{ e}$ $\times 6 \text{ e}$ $\times 8 \text{ e}$ $\times 1 \text{ e}$ $\times 1 \text{ e}$



$$\frac{N}{N+N'} \quad \epsilon \text{ in } \text{atm} = \text{atm} = n'$$

$U = 1 \text{ e}$ ϵ Joules

$P = \text{Joules}$ ϵ elect. Field

ϵ $\text{gr. mol. } NO_3$ $\times 1 \text{ cm}^3$ $\times 1 \text{ e}$ $\epsilon = 3 \times 10^8$

$$CU \times \phi = \frac{Jt}{\epsilon} n$$

$$CV \times \phi = \frac{Jt}{\epsilon} n'$$

$$U = \frac{Jn}{\epsilon \phi C}$$

$$V = \frac{Jn'}{\epsilon \phi C}$$

$$\epsilon = \frac{Jn}{V \phi C} = f \text{ } \epsilon \text{ gr. } \text{Joules}$$

20% 5% Lösung
 1 gr. - 20 cm³
 2615 gr. 730 cm
 $C = \frac{1}{730}$

m. HCl
 $\epsilon = 9659$
 $q = 10$
 $n = 0.81$

$f = 10^{-1}$

$U = 0.00061 \text{ cm}$

$V = 0.00014 \text{ cm}^3 \text{ pro}$

Potentialgefälle P, $\kappa = \sqrt{\rho}$

$U = \frac{\kappa P_m}{C \epsilon}$

$V = \frac{\kappa P_m'}{C \epsilon}$

$U+V = \frac{\kappa P}{C \epsilon}$

$\frac{\kappa}{C} = \text{molekulare } \lambda \text{ (A)}$

λ für $P = 0.01$ $\frac{1}{10} m$ λ $\approx 10^{-2}$ $\approx 10^{-2}$ $\approx 10^{-2}$
 (50 cm²)

oder λ'

$\approx \sqrt{\rho}$ λ $\approx \sqrt{\rho}$ $\approx \sqrt{\rho}$; \approx Potentialgefälle $\approx \sqrt{\rho}$

bei ϵ^2 Kationen \approx \approx \approx

$f = \frac{1}{\epsilon} \sqrt{\rho}$ (Kohlensäure: $\approx 10^{-2}$ $\approx 10^{-2}$)

$\lambda = u_0 + v_0$

$u_0 = \frac{U \epsilon}{P}$ $\approx 10^{-2}$ $\approx 10^{-2}$

$\approx 10^{-2}$ $\approx 10^{-2}$ $\approx 10^{-2}$ $\times 10^{-2}$

	K	Na	Li	Ag	NaN ₃	H		cca 18°C
$u_0 \cdot 10^7$	60	41	33	52	60	290		80

$v_0 \cdot 10^7$	63	63	58	52	54	165	44	33
	Cl	J	NO ₃	ClO ₃	ClO ₄	H ₂ O	CHO ₂	C ₂ H ₂ O ₂

dabei ist $\lambda = \frac{h}{m} / \text{yr}$; ρ in mol. of ; ρ in mol. 1 Liter
 $m = 1000 \cdot C$

[S f f e \rightarrow n f f f Ag]

$m, n, \text{AgNO}_3 \quad \frac{52}{52+58} = 0.473, \text{Ag}$

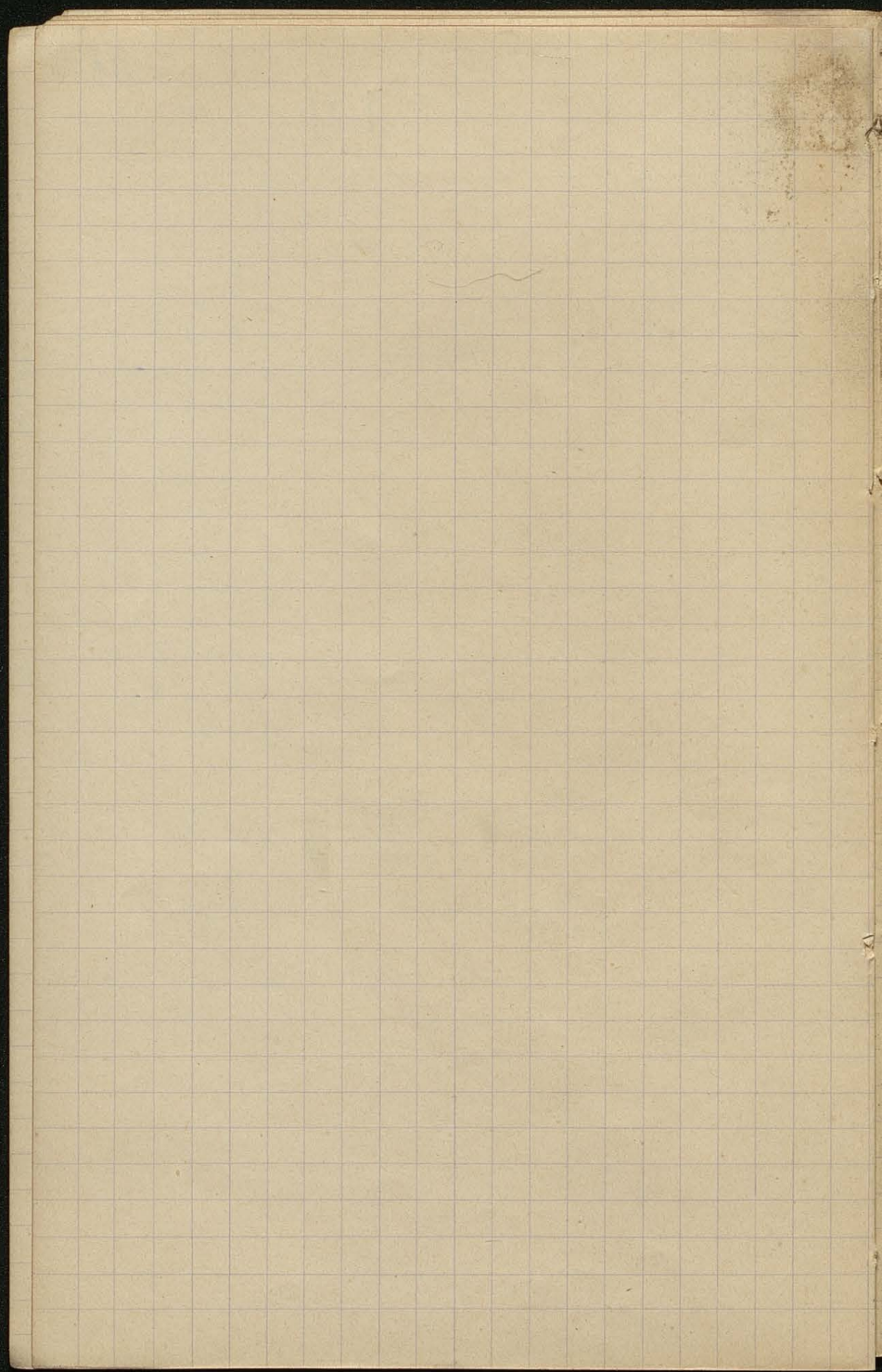
$m, \rho \text{ f f } \rightarrow \rho \text{ 1 Volt pro cm} \quad 1 \text{ km} = 1000 \text{ m} \rightarrow \rho = 109$

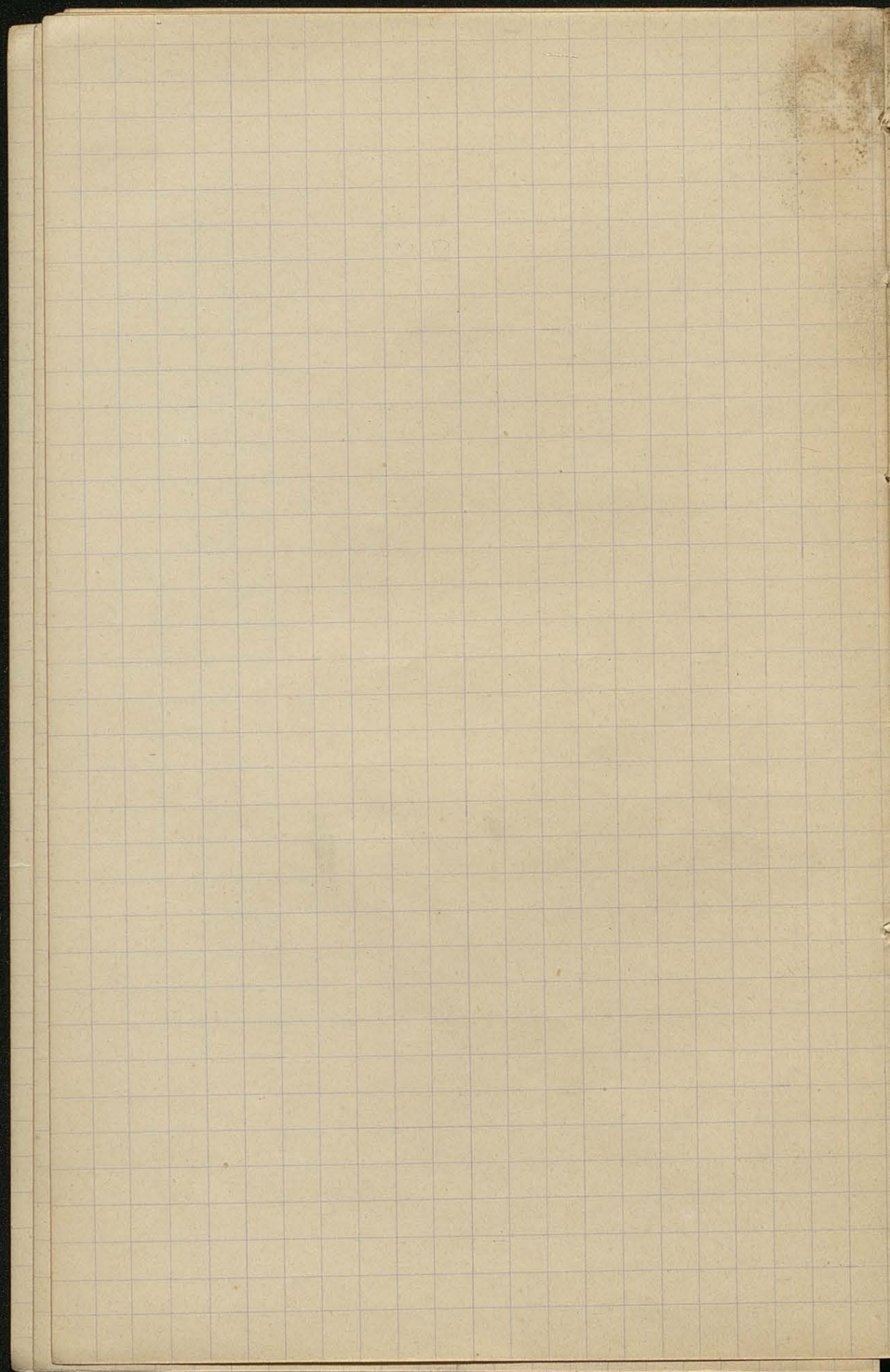
$P = 10^{+8} \quad \Sigma = 9653 \quad u_0(\text{obs}) = \rho \cdot 1000 \cdot u_0 \cdot \lambda_{\text{Ag}}$

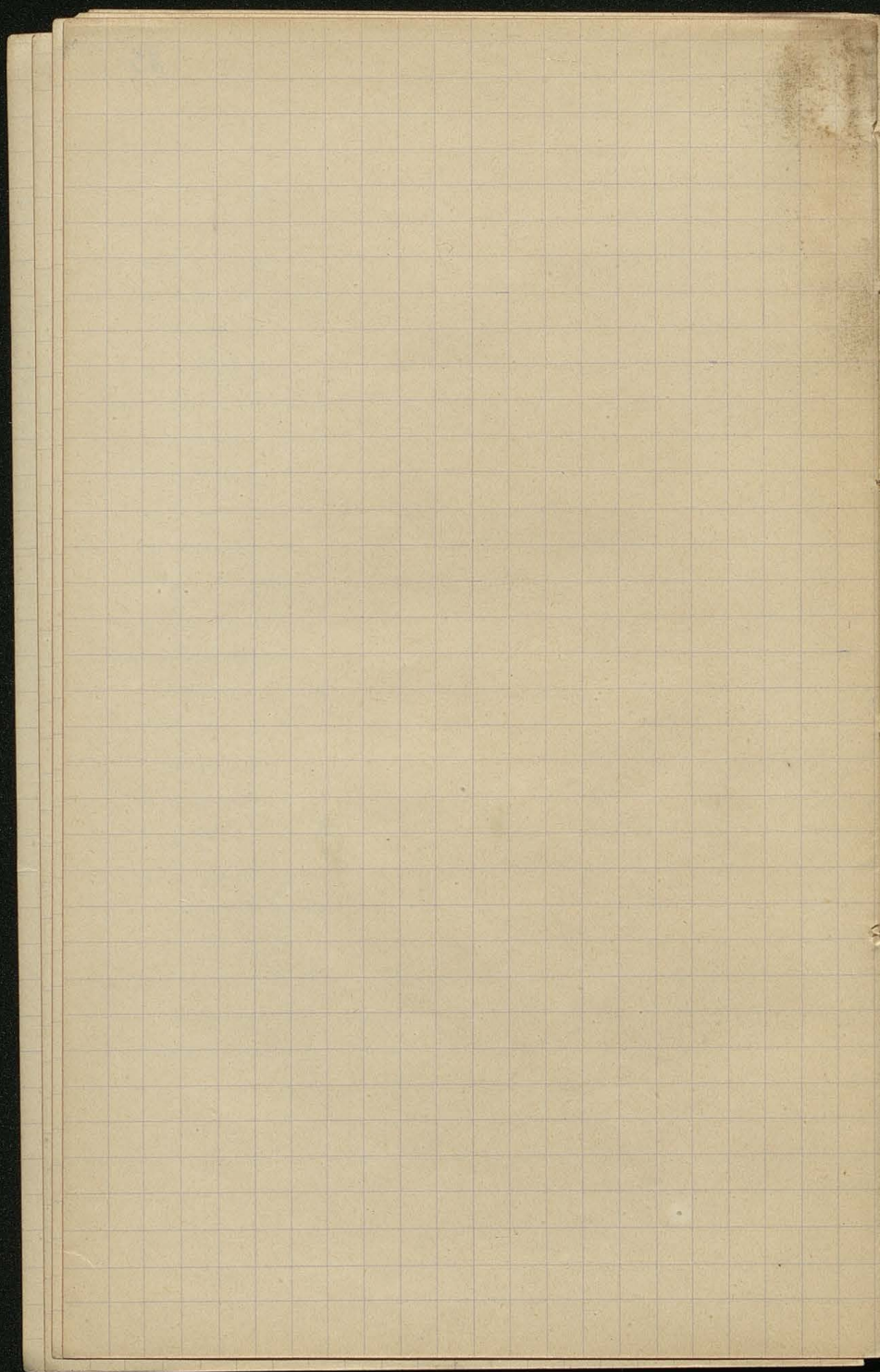
$U = \frac{10^8 \cdot 10^3 \cdot 10^{-7} \cdot 52 \cdot \cancel{10^3} \cdot \cancel{10^3} \cdot 1863 \cdot 10^3}{9653 \cdot \cancel{10^3} \cdot 109}$

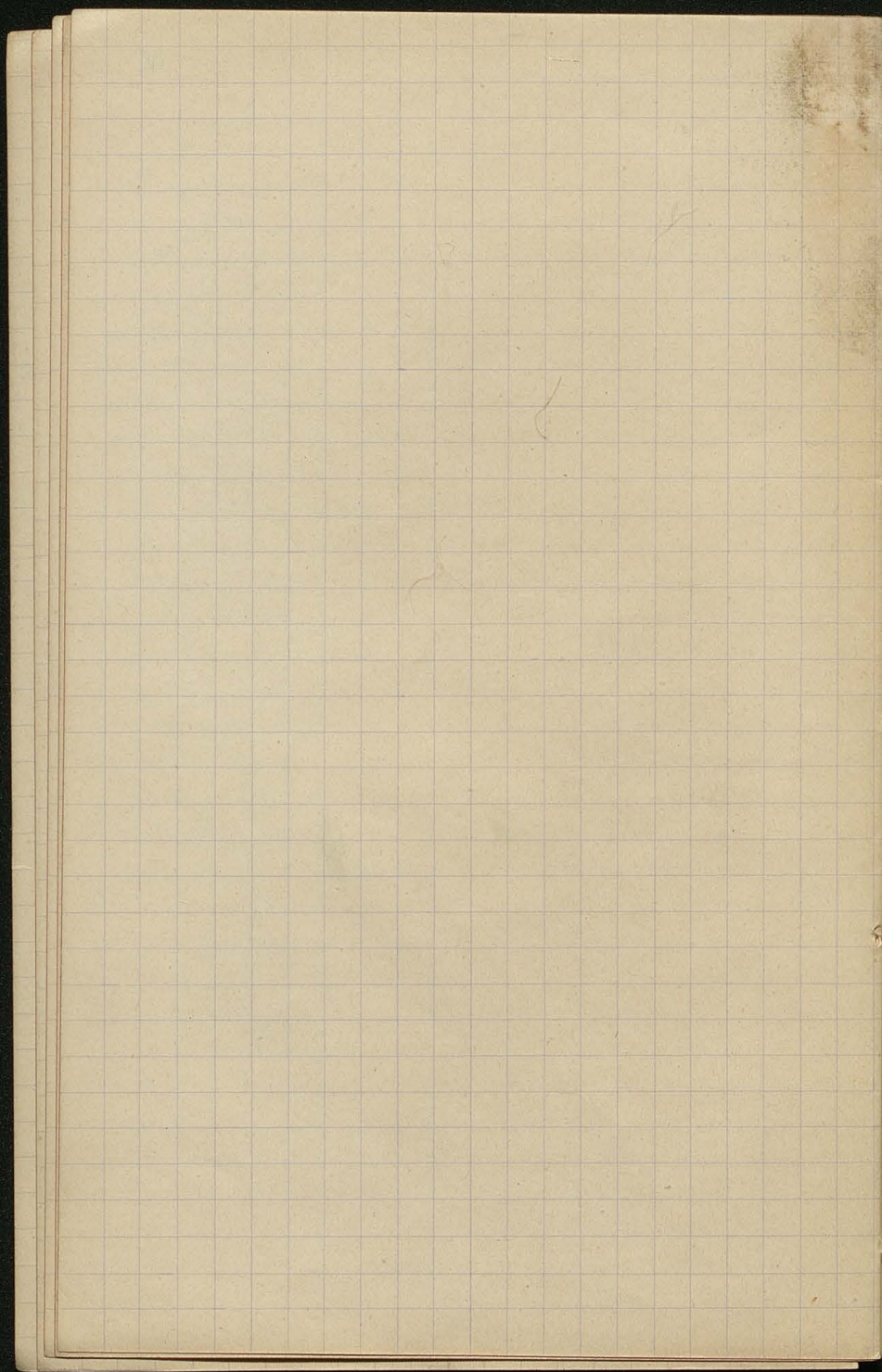
$= 52 \cdot 10^{-7} \cdot 10^2$

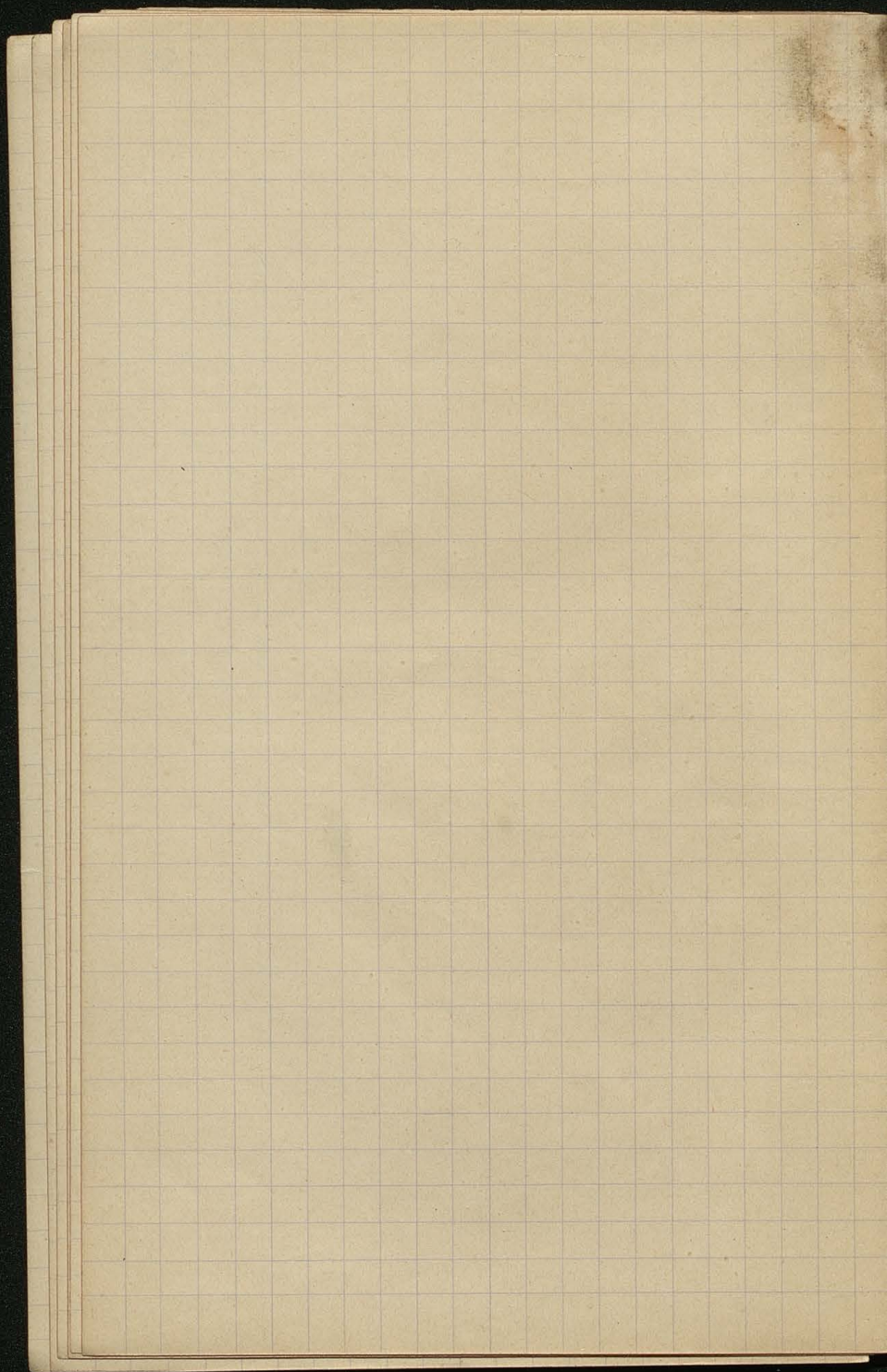
B.J

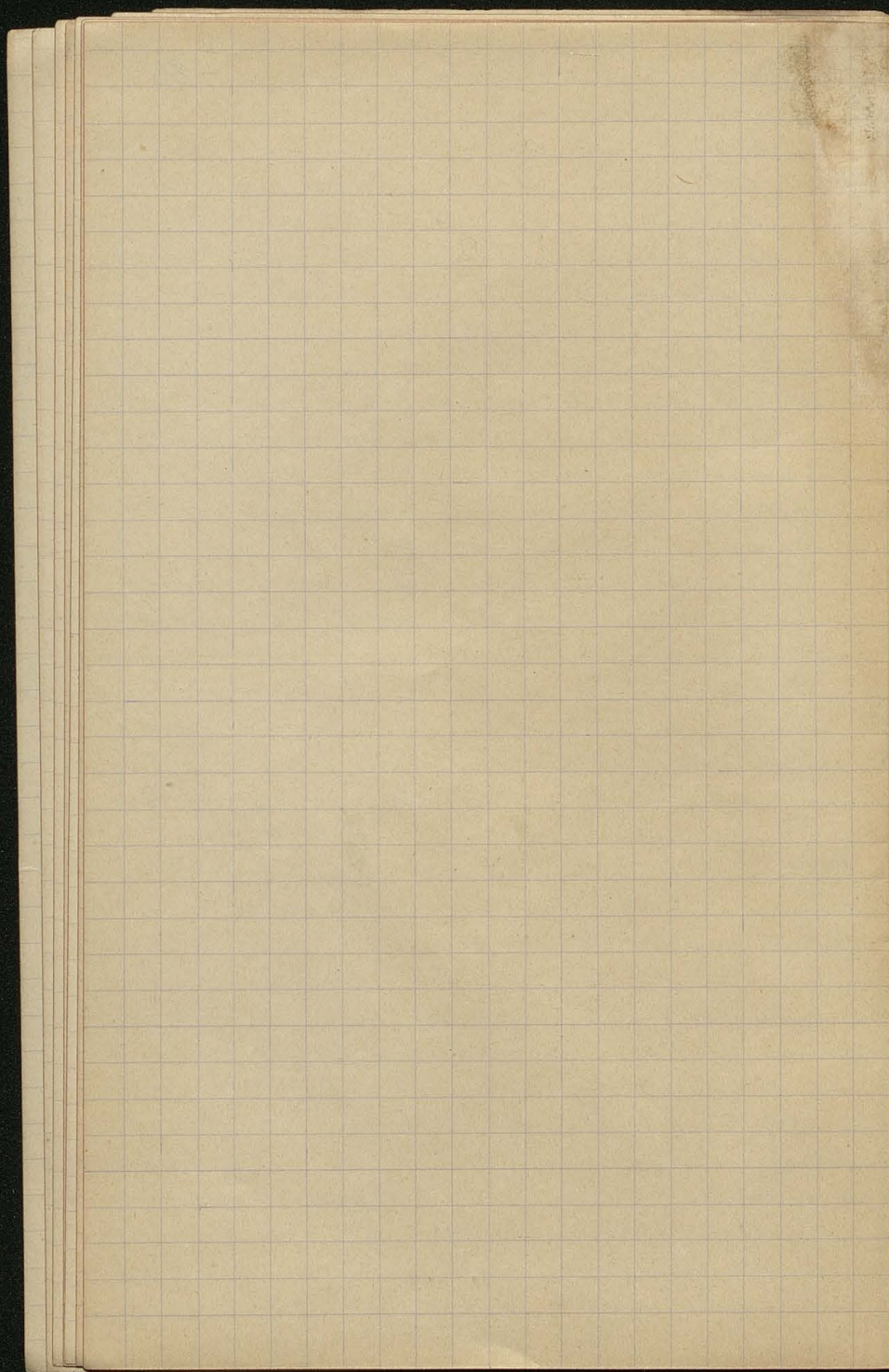


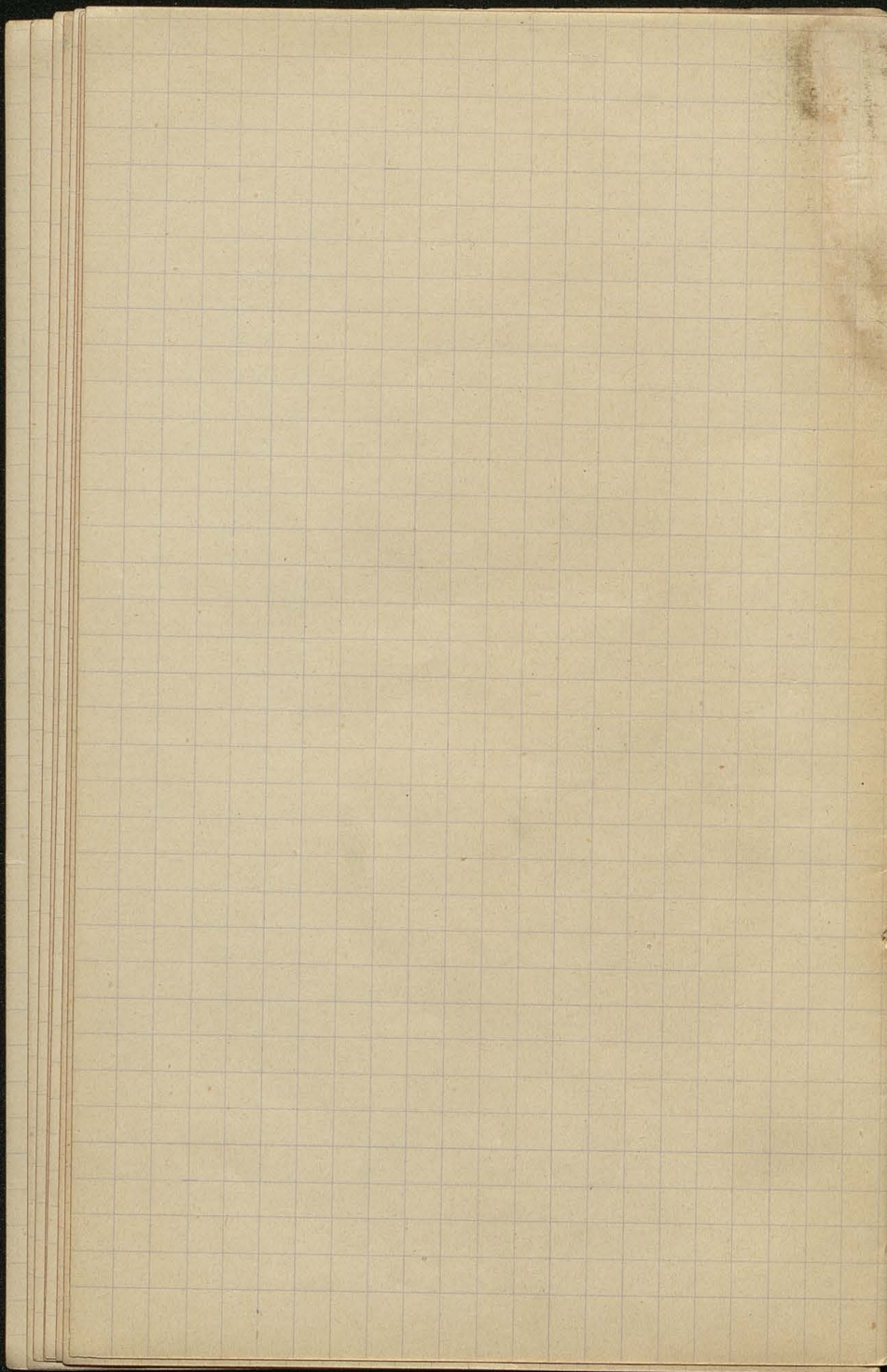


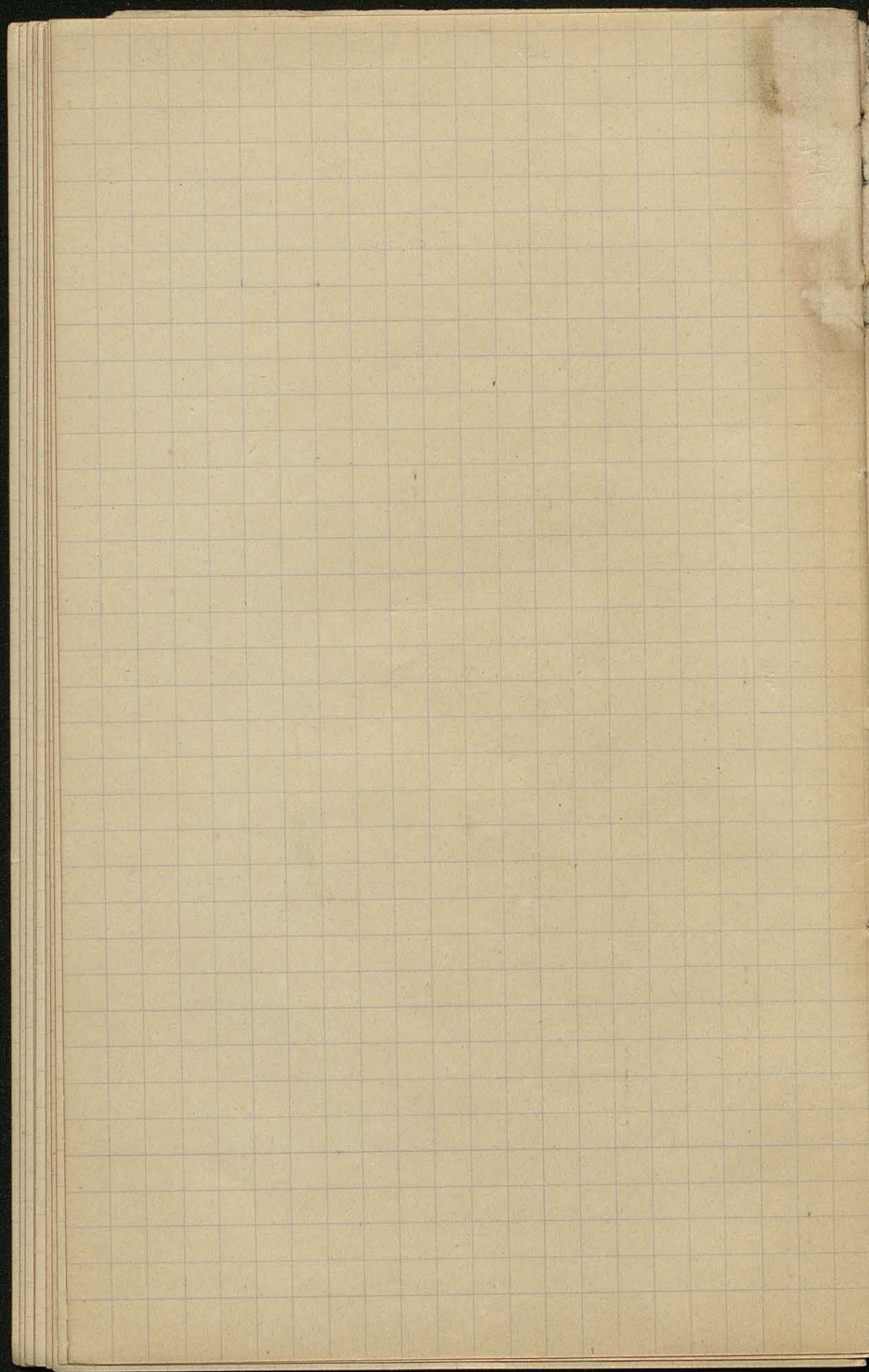


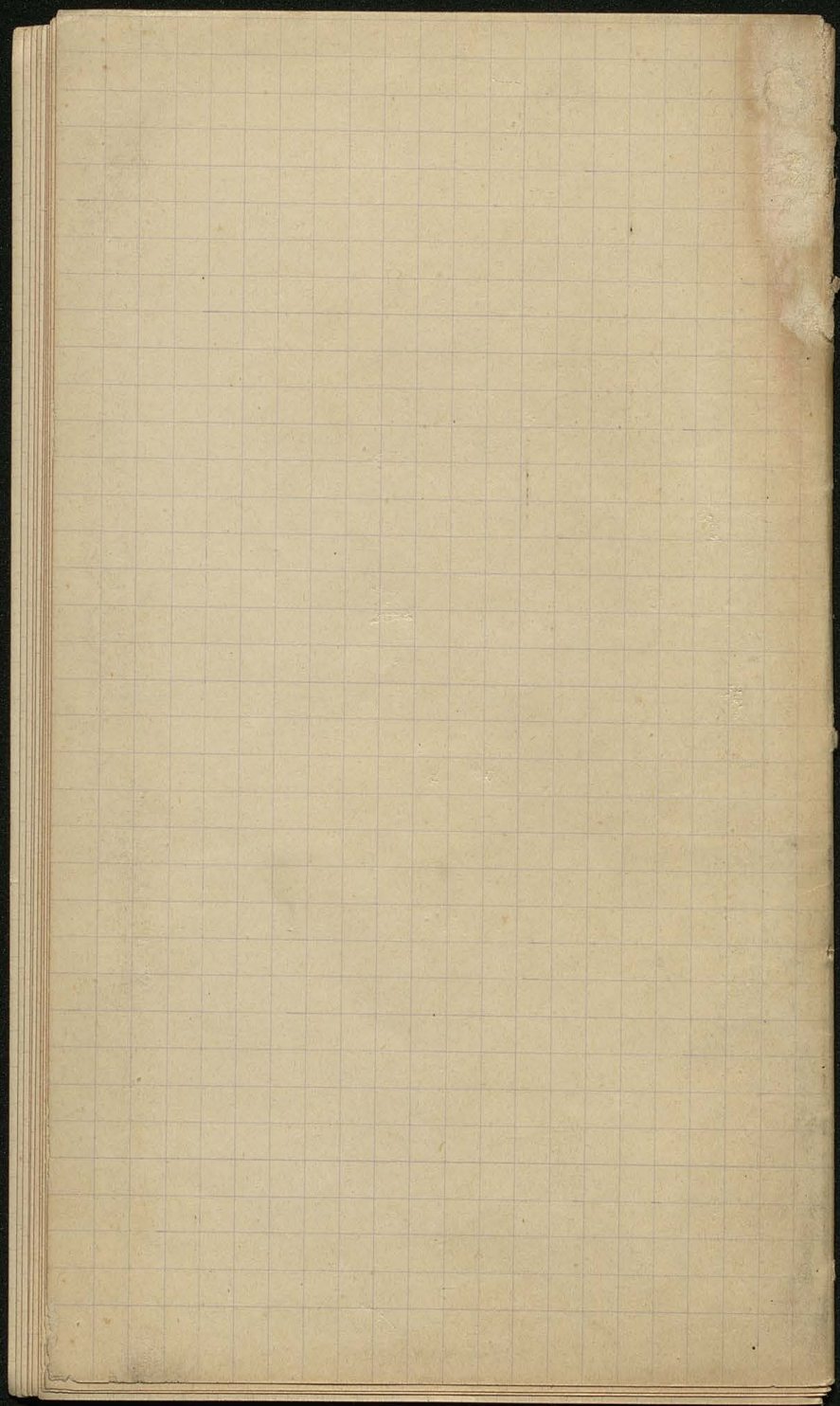


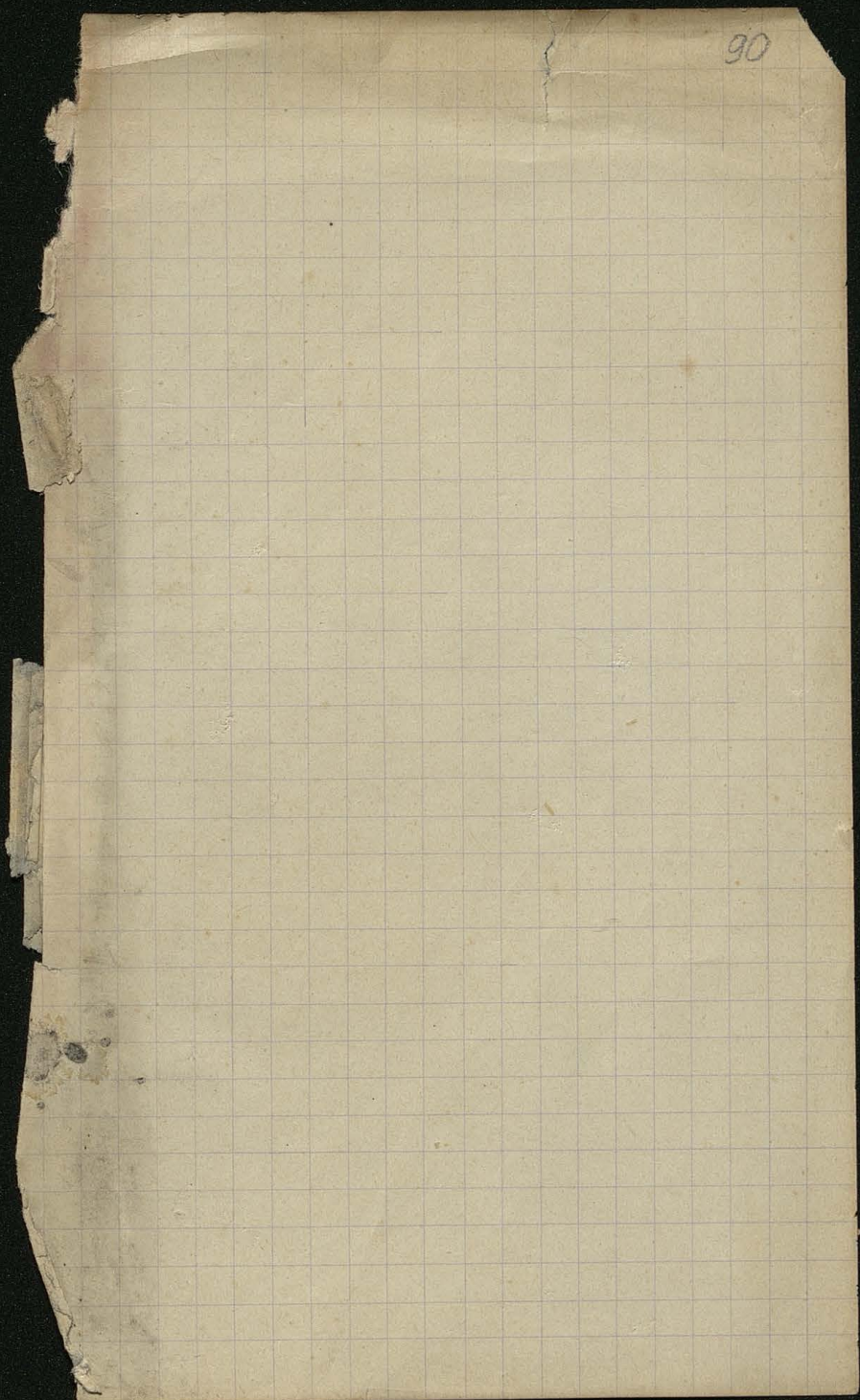


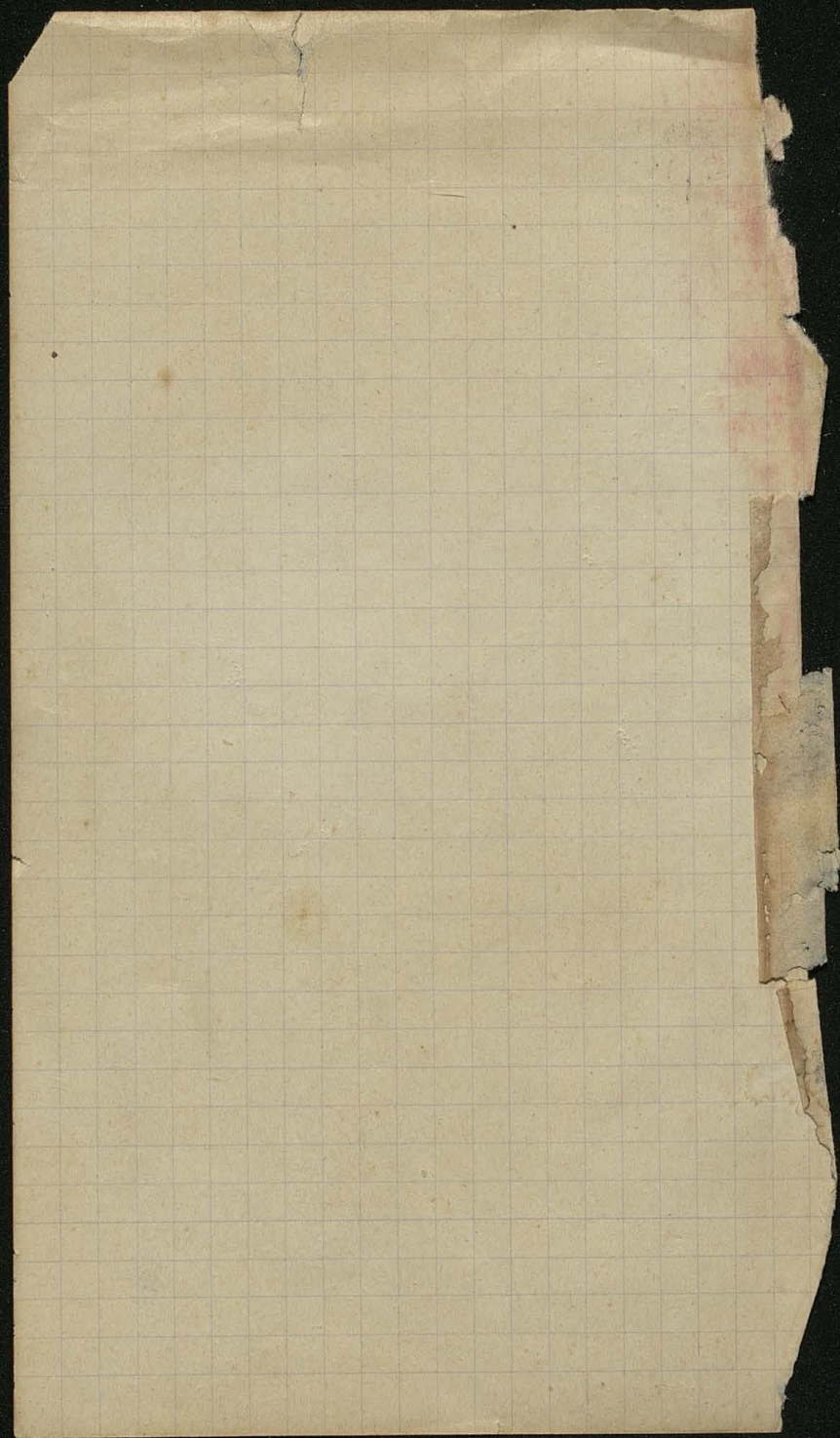






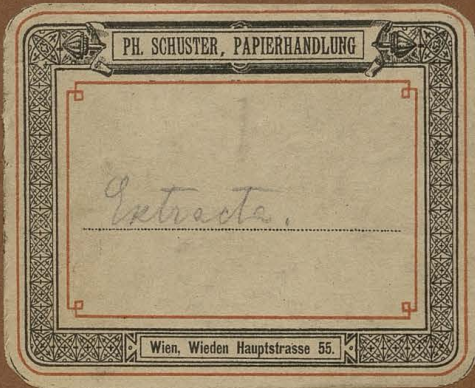


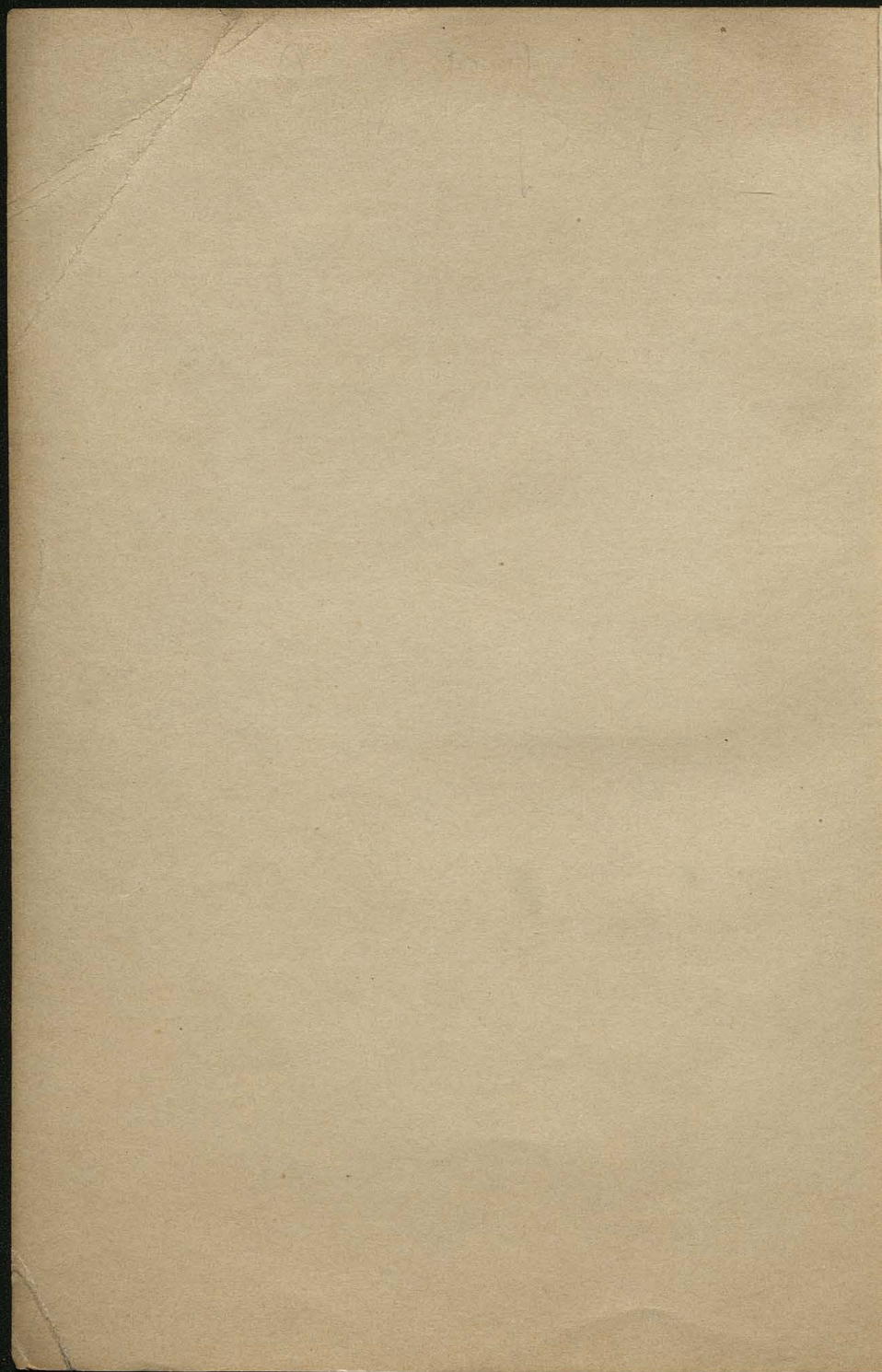




1545

91





100

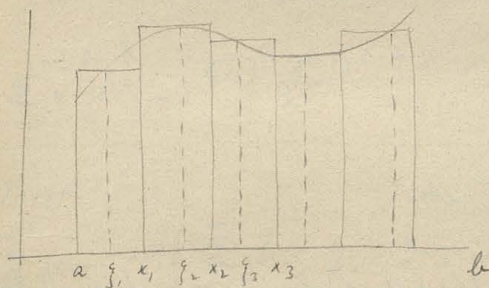
B.J.
 Riemann: Partielle Differential-ll.
 und Anwendungen auf Physik.

Braunschweig 1869
 Vieweg.

Newton's Principia philosophiæ nat. math. 1687
 Erste von $\int x^a$ p. 8. f. d. 2. Element [1706] 1747
 Fourier's $\int \sqrt{x}$ 1807

I Bestimmte \int

✓



$$S = (x_1 - a) f(\xi_1) + (x_2 - x_1) f(\xi_2) + \dots + (b - x_{n-1}) f(\xi_n)$$

2) $M \sim \gamma$, $m \sim \alpha$ $\int y^2$, bei \int $\alpha < \beta$ $\int \alpha < \beta$

3) $M = m$ $\int \alpha < \beta$

$$M = (x'_1 - a) f(\xi'_1) + (x'_2 - x'_1) f(\xi'_2) + \dots + (b - x'_{n-1}) f(\xi'_n)$$

$$m = (x''_1 - a) f(\xi''_1) + (x''_2 - x''_1) f(\xi''_2) + \dots + (b - x''_{n-1}) f(\xi''_n)$$

$a = m$ Δx / $f(b) = f(a) + \dots$ Δx : r_1, r_2, \dots, r_{n-1}

$$M = (r_1 - a) f(p_1') + (r_2 - r_1) f(p_2') + \dots + (b - r_{n-1}) f(p_{n-1}')$$

$$[\dots]$$

$$m = (r_1 - a) f(p_1'') + (r_2 - r_1) f(p_2'') + \dots + (b - r_{n-1}) f(p_{n-1}'')$$

$$M - m = (r_1 - a) [f(p_1') - f(p_1'')] + (r_2 - r_1) [f(p_2') - f(p_2'')] + \dots + (b - r_{n-1}) [f(p_{n-1}') - f(p_{n-1}'')]$$

$$p_k' \leq p_k'' \Rightarrow f(p_k') \geq f(p_k'') \text{ if } f \text{ is decreasing}$$

resp.

$$a \leq d = \max \{ f(x) \} - \min \{ f(x) \} < \epsilon$$

$$M - m < d [r_1 - a + r_2 - r_1 + \dots + b - r_{n-1}]$$

$$M - m < d(b - a)$$

$\epsilon < \delta$ etc.

$$\int_a^b f(x) dx = \lim \left\{ (x_1 - a) f(x_1) + (x_2 - x_1) f(x_2) + \dots + (b - x_{n-1}) f(x_{n-1}) \right\}$$

$$= \lim \left\{ (x_1 - a) f(x_1) + (x_2 - x_1) f(x_2) + \dots \right\}$$

Wallis $\int_0^1 x^k dx = \frac{1}{k+1}$ $[k \in \mathbb{N}]$

$a < x_1 < x_2 < \dots < x_n < b$ - geom. Progr. q

$$\sqrt[n]{\frac{b}{a}} = q$$

$$x_1 = aq$$

$$x_2 = aq^2$$

$$\dots$$

$$x_n = aq^n$$

$$\Sigma = \sum_{v=0}^{n-1} f(a\delta^v) a \delta^v (\delta-1)$$

z.B. $f(x) = x^k$.

$$\Sigma = a^{k+1} (\delta-1) \sum_0^{n-1} \delta^{(k+1)v}$$

$$\begin{aligned} \sum_0^{n-1} \delta^{(k+1)v} &= 1 + \delta^{k+1} + \delta^{2(k+1)} + \dots + \delta^{(n-1)(k+1)} \\ &= \frac{\delta^{n(k+1)} - 1}{\delta^{k+1} - 1} \end{aligned}$$

$$\Sigma = a^{k+1} \left[\left(\frac{\delta}{a} \right)^{k+1} - 1 \right] \frac{\delta-1}{\delta^{k+1} - 1} = (\delta^{k+1} - a^{k+1}) \frac{\delta-1}{\delta^{k+1} - 1}$$

$$\lim \frac{\delta-1}{\delta^{k+1} - 1} = \lim \frac{1}{1 + \delta + \delta^2 + \dots + \delta^k} = \frac{1}{k+1}$$

$$\Sigma = \left[\delta^{k+1} - a^{k+1} \right] \frac{1}{k+1}$$

Druck- \int

$$\sum_{v=0}^{n-1} (x_{v+1} - x_v) \left\{ \sum_{m=0}^{n-1} f(x_v, y_m) (y_{m+1} - y_m) \right\}$$

$$= \sum_{\mu=0}^{n-1} (y_{\mu+1} - y_\mu) \left\{ \sum_{v=0}^{n-1} f(x_v, y_\mu) (x_{v+1} - x_v) \right\}$$

$$\int_a^b dx \left\{ \int_g^h f(x, y) dy \right\} = \int_g^h dy \left\{ \int_a^b f(x, y) dx \right\}$$

$$\begin{aligned} \iint x^{h-1} y^{g-1} dx dy &= \int_a^h dy \int_a^h x^{h-1} dx = \int_a^h \frac{h y^g - a^g}{y} dy \\ &= \int_a^h dx \int_a^h x^{h-1} y^{g-1} dy = \int_a^h \frac{x^{h-1} - x^{g-1}}{h y^g} dx \end{aligned}$$

$$\begin{cases} a=0 \\ b=1 \end{cases} \int_0^1 \frac{x^{h-1} y^{g-1}}{h y^g} dx = \int_a^h \frac{dy}{y} = \log \frac{h}{g}$$

$$h > 0$$

$$g > 0$$

$$\int e^{-ax} dx = \frac{e^{-ax}}{-a} + C$$

$$\int e^{-(a+bi)x} dx = \frac{e^{-(a+bi)x}}{-(a+bi)} + C$$

$$\int_0^{\infty} \dots = \frac{1}{a+bi} \quad a > 0$$

$$e^{-(a+bi)x} = e^{-ax} (\cos bx - i \sin bx)$$

$$\int_0^{\infty} e^{-ax} \cos bx dx = \frac{a}{a^2+b^2} \quad a > 0$$

$$\int_0^{\infty} e^{-ax} \sin bx dx = \frac{b}{a^2+b^2} \quad a > 0$$

$$\frac{1}{y} = \int_0^{\infty} e^{-yx} dx$$

$$\int_0^h \frac{\sin \beta y}{y} dy = \int_0^h \sin \beta y dy \int_0^{\infty} e^{-y^x} dx$$

$$= \int_0^{\infty} dx \int_0^h e^{-y^x} \sin \beta y dy$$

$$\begin{cases} g=0 \\ h=\infty \end{cases}$$

$$\int_0^{\infty} \frac{\sin \beta y}{y} dy = \begin{cases} \frac{\pi}{2} & \beta > 0 \\ 0 & \beta = 0 \\ -\frac{\pi}{2} & \beta < 0 \end{cases}$$

$$\int_0^{\infty} \frac{\sin y}{y} \cos \mu y dy = \frac{1}{2} \int_0^{\infty} \frac{\sin (1+\mu) y}{y} dy + \frac{1}{2} \int_0^{\infty} \frac{\sin (1-\mu) y}{y} dy$$

$$\int_0^{\infty} \frac{\sin y}{y} \cos \mu y dy = \begin{cases} 0 & \mu > 0 \\ \frac{\pi}{4} & \mu = 1 \\ \frac{\pi}{2} & 1 > \mu > -1 \\ \frac{\pi}{4} & \mu = -1 \\ 0 & -1 > \mu \end{cases}$$

$$\int_0^{\infty} \frac{y \sin ky}{y^2 + b^2} dy = \int_0^{\infty} \cos bx dx \int_0^{\infty} e^{-xy} \sin ky dy$$

$$= \int_0^{\infty} \frac{k \cos bx}{k^2 + x^2} dx$$

$$y = \frac{b}{k} x$$

$$\underbrace{\int_0^{\infty} \frac{k \cos bx}{k^2 + x^2} dx}_{= ku} = \underbrace{\int_0^{\infty} \frac{x \sin bx}{k^2 + x^2} dx}_{= -\frac{du}{db}} \quad \frac{b}{k} > 0$$

$$k h = -b g u + b y c$$

$$u = c e^{-k h}$$

$$h=0 \quad u=c = \int_0^{\infty} \frac{dx}{k+x} = \frac{1}{k} \arctan \frac{x}{k} + c$$

$$c = \pm \frac{\pi}{2k}$$

$$\int_0^{\infty} \frac{x \sin bx}{k+x} dx = \begin{cases} \frac{\pi}{2} e^{-bk} & b > 0, k > 0 \\ -\frac{\pi}{2} e^{-bk} & b < 0, k < 0 \end{cases}$$

∴ er søges $\int_0^{\infty} \frac{e^{-bx} \sin kx}{k+x} dx$:

$$= \begin{cases} \frac{\pi}{2} e^{-b\sqrt{k^2}} & b > 0 \\ -\frac{\pi}{2} e^{-b\sqrt{k^2}} & b < 0 \end{cases}$$

∴ ∴:

$$\int_0^{\infty} \frac{e^{-bx} \sin kx}{k+x} dx = \begin{cases} \frac{\pi}{2\sqrt{k^2}} e^{-b\sqrt{k^2}} & b > 0 \\ \frac{\pi}{2\sqrt{k^2}} e^{b\sqrt{k^2}} & b < 0 \end{cases}$$

~~It~~

$k > 0$

$$\int_0^{\infty} \frac{\sin bx + k \cos bx}{k+x} dx = \begin{cases} \pi e^{-bk} & b > 0 \\ \frac{\pi}{2} & b = 0 \\ 0 & b < 0 \end{cases}$$

$$\int_0^{\infty} e^{-ax} dx = \frac{1}{a}$$

$$\int_0^{\infty} e^{-ax} x^{2n} dx = \frac{1}{a} \sqrt{\frac{\pi}{a}} \frac{(2n)!}{n!} \frac{1}{(4a)^n}$$

$$\begin{aligned} \int_0^{\infty} e^{-ax} \cos px dx &= \int_0^{\infty} da e^{-ax} \sum_{n=0}^{\infty} \frac{(-p^2 x^2)^n}{(2n)!} \\ &= \sum_0^{\infty} \frac{(-p^2)^n}{(2n)!} \frac{1}{a} \sqrt{\frac{\pi}{a}} \frac{(2n)!}{n!} \frac{1}{(4a)^n} \\ &= \frac{1}{a} \sqrt{\frac{\pi}{a}} e^{-\frac{p^2}{4a}} \end{aligned}$$

Unendliche Reihen.

$$\frac{1}{1} + \frac{1}{2} + \frac{1}{3} + \dots$$

$$= \int_0^1 \frac{1-x^n}{1-x} dx$$

$$\int_0^1 \frac{dx}{1-x} = -\log [1 - (1-k)] = -\log k$$

$$\lim_{k \rightarrow 0} \dots = \infty$$

$$\lim \left[\underbrace{1 + \frac{1}{2} - 1 + \frac{1}{3} + \frac{1}{4} - \frac{1}{2} + \dots}_{n} \right] = \lim \sum_{m=1}^{2n} \frac{1}{m}$$

$$= \lim \frac{1}{n} \left[\frac{1}{1+\frac{1}{n}} + \frac{1}{1+\frac{2}{n}} + \dots \right] = \int_1^2 \frac{dx}{x} = \log 2$$

$$\lim \left[\underbrace{1 + \frac{1}{2} + \dots + \frac{1}{k} - 1 + \frac{1}{k+1} + \frac{1}{k+2} + \dots + \frac{1}{2k} - \frac{1}{2} + \dots}_{k+1} \right] = \lim \sum_{m=1}^{kn} \frac{1}{m}$$

$$= \lim \frac{1}{n} \left\{ \frac{1}{1+\frac{1}{n}} + \dots + \frac{1}{1+\frac{(k-1)m}{n}} \right\} = \int_1^k \frac{dx}{x} = \log k$$

$$f(x) = a_1 \sin x + a_2 \sin 2x + \dots + a_{n-1} \sin(n-1)x$$

$$= 0 \quad \left. \begin{array}{l} x=0 \\ x=n \end{array} \right\}$$

$$0 < x < n$$

$(n-1)$ $\frac{1}{n}$ $\frac{2}{n}$ $\frac{3}{n}$ \dots $\frac{n-1}{n}$ $\frac{1}{n}$ $\frac{2}{n}$ $\frac{3}{n}$ \dots $\frac{n-1}{n}$ $\frac{1}{n}$ $\frac{2}{n}$ $\frac{3}{n}$ \dots $\frac{n-1}{n}$

$$f\left(\frac{x}{n}\right) = a_1 \sin \frac{x}{n} + a_2 \sin \frac{2x}{n} + \dots + a_{n-1} \sin \frac{(n-1)x}{n}$$

$$f\left(\frac{x}{n}\right) = a_1 \sin \frac{x}{n} + a_2 \sin \frac{2x}{n} + \dots + a_{n-1} \sin \frac{(n-1)x}{n} \quad \left| \begin{array}{l} 2 \sin \frac{x}{n} \\ 2 \sin \frac{2x}{n} \\ \dots \\ 2 \sin \frac{(n-1)x}{n} \end{array} \right.$$

$$f\left(\frac{2x}{n}\right) = a_1 \sin \frac{2x}{n} + a_2 \sin \frac{4x}{n} + \dots + a_{n-1} \sin \frac{2(n-1)x}{n}$$

$$f\left(\frac{(n-1)x}{n}\right) = a_1 \sin \frac{(n-1)x}{n} + \dots + a_{n-1} \sin \frac{(n-1)(n-1)x}{n} \quad \left| \begin{array}{l} 2 \sin \frac{(n-1)x}{n} \\ \dots \\ 2 \sin \frac{(n-1)(n-1)x}{n} \end{array} \right.$$

f ist Lagrange f ist n Gr. am i

$$e_k \sin kx: A_k = 2 \left(\sin k \frac{x}{n} \sin m \frac{x}{n} + \dots + \sin k \frac{(n-1)x}{n} \sin m \frac{(n-1)x}{n} \right)$$

$$= \cos(k-m) \frac{x}{n} + \cos(k-m) \frac{2x}{n} + \dots + \cos(k-m) \frac{(n-1)x}{n} - \left[\cos(k+m) \frac{x}{n} + \dots + \cos(k+m) \frac{(n-1)x}{n} \right]$$

$$s = \cos \theta + \cos 2\theta + \dots + \cos(n-1)\theta \quad \left| \begin{array}{l} 2 \cos \theta \\ 2 \cos 2\theta \\ \dots \\ 2 \cos(n-2)\theta \\ \cos(n-1)\theta \end{array} \right.$$

$$2s \cos \theta = 1 + \cos \theta + \cos 2\theta + \dots + \cos(n-2)\theta \quad \left. \begin{array}{l} \cos(n-2)\theta \\ \cos(n-1)\theta \end{array} \right\} = 1 - \cos(n-1)\theta + s$$

$$s = -\frac{1}{2} + \frac{\sin(2n-1)\frac{\theta}{2}}{2 \sin \frac{\theta}{2}}$$

$$k \geq m \quad A_k = 0$$

$$A_m = 1 + 1 + \dots + 1 -$$

$$\left[\cos 2m \frac{\pi}{n} + \cos 2m \frac{2\pi}{n} + \dots + \cos 2m \frac{(n-1)\pi}{n} \right] \Bigg\} = n$$

$$n a_m = 2 f\left(\frac{\pi}{n}\right) \sin \frac{m\pi}{n} + 2 f\left(\frac{2\pi}{n}\right) \sin \frac{2m\pi}{n} + \dots - 2 f\left(\frac{(n-1)\pi}{n}\right) \sin \frac{(n-1)m\pi}{n}$$

$$m = 1, 2, 3, \dots, n-1$$

sr $f \infty$ is

$$a_m = \frac{2}{n} \frac{\pi}{n} \left\{ f\left(\frac{\pi}{n}\right) \sin \frac{m\pi}{n} + \dots + f\left(\frac{(n-1)\pi}{n}\right) \sin \frac{(n-1)m\pi}{n} \right\}$$

$$\lim_{n \rightarrow \infty} \frac{\pi}{n}$$

$$a_m = \frac{2}{\pi} \int_0^{\pi} f(x) \sin mx \, dx$$

203.

I. $f(x) = x$

$$\frac{1}{2} x = \sin x - \frac{\sin 2x}{2} + \frac{\sin 3x}{3} - \frac{\sin 4x}{4} + \dots$$

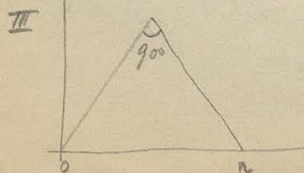
$$+\pi > x > -\pi$$

II. $f(x) = 1$

$$\frac{\pi}{4} = \frac{\sin x}{1} + \frac{\sin 3x}{3} + \frac{\sin 5x}{5} + \dots \quad x > 0$$

$$0 = \dots \quad x = 0$$

$$-\frac{\pi}{4} = \dots \quad x < 0$$



$$\int_0^{\pi} f(x) \sin mx \, dx = \int_0^{\pi/2} f(x) \sin mx \, dx + \int_{\pi/2}^{\pi} f(x) \sin mx \, dx$$

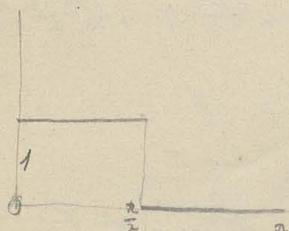
$$= \int_0^{\pi/2} x \sin mx \, dx + \int_{\pi/2}^{\pi} (x - \pi) \sin mx \, dx$$

$$f(x) = \frac{4}{\pi} \left\{ \frac{\sin x}{1^2} - \frac{\sin 3x}{3^2} + \frac{\sin 5x}{5^2} + \dots \right\}$$

$$\pi > x > 0$$

$$x = \frac{\pi}{2} \quad \frac{\pi^2}{8} = \frac{1}{1^2} + \frac{1}{3^2} + \frac{1}{5^2} + \dots$$

IV



$$\frac{\pi}{2} = \sin x + \sin 2x + \frac{\sin 3x}{3} + \frac{\sin 5x}{5} + \frac{\sin 6x}{7} + \dots$$

$$\frac{\pi}{2} > x > 0$$

$$0 = \sin x + \dots$$

$$\pi > x > \frac{\pi}{2}$$

$$\left(\frac{\pi}{4}\right) = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots$$

$$x = \frac{\pi}{2}$$

Cos-Reihen

$$f(x) = \frac{1}{2} f(x) \sin x + \frac{1}{2} f(x) \cos x$$

$$2 f(x) \sin x = a_0 \sin x + a_1 \sin 2x + a_2 \sin 3x + \dots$$

$$a_m = \frac{2}{\pi} \int_0^{\pi} f(x) \sin x \sin mx \, dx$$

$$b_m = \frac{2}{\pi} \int_0^{\pi} f(x) \cos x \, dx$$

$$a_m = b_{m-1} - b_{m+1}$$

$$2 f(x) \sin x = (b_0 - b_1) \sin x + (b_1 - b_2) \sin 2x + \dots$$

$$f(x) = \frac{1}{2} b_0 + b_1 \cos x + \dots$$

$$b_m = \frac{2}{\pi} \int_0^{\pi} f(x) \sin mx dx$$

98

$x \rightarrow$ Dann $\int_{-x}^x \cos \dots dx$

" " \cos " " " \sin " " "

" " $\sin + \cos$ " " \sin " " "

$$f(x) = \frac{f(x) + f(-x)}{2} + \frac{f(x) - f(-x)}{2}$$

$$\frac{f(x) + f(-x)}{2} = \frac{b_0}{2} + b_1 \cos x + \dots$$

$$b_m = \frac{2}{\pi} \int_0^{\pi} \frac{f(x) + f(-x)}{2} \cos mx dx$$

$$\frac{f(x) - f(-x)}{2} = a_1 \sin x + a_2 \sin 2x + \dots$$

$$a_m = \frac{2}{\pi} \int_0^{\pi} \frac{f(x) - f(-x)}{2} \sin mx dx$$

$$f(x) = \frac{b_0}{2} + b_1 \cos x + b_2 \cos 2x + \dots$$

$$+ a_1 \sin x + a_2 \sin 2x + \dots$$

$$b_m = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos mx dx$$

 $\pi > x > -\pi$

$$a_m = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin mx dx$$

Convergenz - Beweis [n. Dirichlet (Celle II, Seite I)]

$$S_{2n+1} = \frac{1}{2} b_0 + b_1 \cos x + b_2 \cos 2x + \dots + b_n \cos nx$$

$$+ a_1 \sin x + a_2 \sin 2x + \dots + a_n \sin nx$$

$$= \frac{i}{\pi} \int_{-n}^n f(x) \frac{\sin(2n+1) \frac{x-\alpha}{2}}{2 \sin \frac{x-\alpha}{2}} d\alpha$$

$$\frac{1}{2} + \cos p + \cos 4p + \dots + \cos 2np = \frac{\sin(2n+1)p}{2 \sin p}$$

$$\int \dots dp$$

$$\frac{n}{2} = \int_0^{\frac{n}{2}} \frac{\sin(2n+1)p}{\sin p} dp$$

$$= \int_0^{\frac{n}{2}} + \int_{\frac{n}{2}}^{\frac{3n}{2}} + \dots$$

$$+ \int_{\frac{3n}{2}}^{\frac{5n}{2}} + \dots + (-1)^n p_n$$

$$= p_0 - p_1 + p_2 - \dots$$

$$p_{2v-1} = (-1)^{v-1} \int_{\frac{(2v-1)n}{2}}^{\frac{2vn}{2}} \frac{\sin kp}{\sin p} dp$$

$$\frac{(-1)^{v-1}}{\sin \frac{(2v-1)n}{2}} \int_{\frac{(2v-1)n}{2}}^{\frac{2vn}{2}} \sin kp dp > p_{v-1} > \frac{(-1)^{v-1}}{\sin \frac{vn}{2}} \int_{\frac{(v-1)n}{2}}^{\frac{vn}{2}} \sin kp dp$$

$$\frac{2}{h \sin \frac{(v-1)n}{2}} > p_{v-1} > \frac{2}{h \sin \frac{vn}{2}} > p_v$$

$$\frac{n}{2} < p_0 - p_1 + p_2 - \dots + p_{2m}$$

$$\frac{n}{2} > p_0 - p_1 + p_2 - \dots - p_{2m-1} \quad 2m < n$$

$$T = \int_0^b f(p) \frac{\sin hp}{\sin p} dp = \int_0^{\frac{n}{2h}} + \int_{\frac{n}{2h}}^{\dots} \dots \int_{\frac{m\pi}{h}}^b$$

für m abnehmend

$$= r_0 - r_1 + r_2 - \dots + (-1)^m r_m$$

$$r_m = (-1)^m \int_{\frac{m\pi}{h}}^b f(p) \frac{\sin hp}{\sin p} dp$$

$$p_{v-1} f\left(\frac{v-1}{h} \pi\right) \geq r_{v-1} \geq p_{v-1} f\left(\frac{v\pi}{h}\right) > p_v f\left(\frac{v\pi}{h}\right) \geq r_v$$

$$T > p_0 f\left(\frac{\pi}{2h}\right) - p_1 f\left(\frac{\pi}{2h}\right) + p_2 f\left(\frac{3\pi}{2h}\right) - \dots - p_{2m-1} f\left(\frac{(2m-1)\pi}{2h}\right)$$

$$> f\left(\frac{2m\pi}{h}\right) [p_0 - p_1 + p_2 - \dots - p_{2m-1}]$$

$$T < p_0 f(0) - f\left(\frac{2m\pi}{h}\right) [p_0 - p_2 + p_4 - \dots - p_{2m}]$$

$$T > f\left(\frac{2m\pi}{h}\right) \left(\frac{n}{2} - p_{2m}\right)$$

$$T < p_0 \left\{ f(0) - f\left(\frac{2m\pi}{h}\right) \right\} + \left(\frac{n}{2} + p_{2m}\right) f\left(\frac{2m\pi}{h}\right)$$

$$n = \infty$$

$$h = \infty$$

$$m = \infty$$

$$\lim \frac{n}{h} = 0$$

$$\lim p_{2m} = 0$$

$$\lim T = \frac{n}{2} f(0)$$

$$\lim_{L \rightarrow \infty} S_{Lm} = \frac{1}{2} [\varphi(x+0) + \varphi(x-0)] \quad n > x > -n$$

$$= \frac{1}{2} [\varphi(-n+0) + \varphi(n-0)] \quad x = \pm n$$

$$\varphi\left(\frac{cx}{n}\right) = \frac{1}{2} b_0 + b_1 \cos 2x + b_2 \cos 4x + \dots$$

$$+ a_1 \sin 2x + a_2 \sin 4x + \dots$$

$$\left\{ \begin{aligned} b_m &= \frac{1}{\pi} \int_{-n}^n \varphi\left(\frac{cx}{n}\right) \cos m x dx \\ a_m &= \dots \end{aligned} \right.$$

$$n > x > -n$$

$$\varphi(x) = \frac{1}{2} b_0 + b_1 \cos \frac{n x}{c} + \dots$$

$$+ a_1 \sin \frac{n x}{c} + \dots$$

$$c > x > -c$$

$$\left\{ \begin{aligned} b_m &= \frac{1}{c} \int_{-c}^c \varphi(x) \cos \frac{m n x}{c} dx \\ a_m &= \dots \end{aligned} \right.$$

$$a_m = \dots$$

$$\varphi(x) = \frac{1}{c} \left\{ -\frac{1}{2} \int_{-c}^c \varphi(x) dx + \sum_{n=1}^{\infty} \int_{-c}^c dx \varphi(x) \cos \frac{m n x}{c} (1-x) \right\}$$

$$\frac{m n}{c} = \alpha \quad c = \infty \quad dx = \frac{n}{c}$$

$$\varphi(x) = \frac{1}{\pi} \int_0^{\infty} d\alpha \left\{ \int_{-\infty}^{\infty} dx \varphi(x) \cos \alpha (1-x) \right\}$$

Reithes dem zur Zurückführung einer D. Gl. mit aus. Gl.

$$1). a_n \frac{d^n y}{dx^n} + a_{n-1} \frac{d^{n-1} y}{dx^{n-1}} + \dots + a_1 y' + a_0 y = X$$

$$\int X \cdot z \, dx$$

$$\int a_n z y \, dx = \int \dots$$

$$\int a_{n-1} z \frac{dy}{dx} \, dx = a_{n-1} z y - \int y \frac{d(a_{n-1} z)}{dx} \, dx$$

$$\frac{d^n (a_0 z)}{dx^n} + \frac{d^{n-1} (a_1 z)}{dx^{n-1}} + \dots + (-1)^n a_n z = 0$$

$$\int X z \, dx = y (a_{n-1} z + \dots)$$

$$+ \frac{dy}{dx} (\dots)$$

$$+ \dots$$

$$+ \frac{d^n y}{dx^n} a_0 z$$

2). $y = u \cdot v$

$$X = u \left\{ a_0 \frac{d^n v}{dx^n} + a_1 \frac{d^{n-1} v}{dx^{n-1}} + \dots + a_{n-1} v' + a_n v \right\} = 0$$

$$+ \frac{du}{dx} \left\{ a_0 \frac{d^{n-1} v}{dx^{n-1}} + \dots + a_{n-1} v' + a_n v \right\}$$

$$+ \dots$$

$$+ \frac{d^n u}{dx^n} a_0 v$$

$$b_0 \frac{d^{n-1} u}{dx^{n-1}} + b_1 \frac{d^{n-2} u}{dx^{n-2}} + \dots + b_{n-1} u' = X$$

3). Variation der Constanten.

Bewegung der Wärme in festen Körpern

$$\frac{\partial u}{\partial t} = a^2 \frac{\partial^2 u}{\partial x^2}$$

$$t=0 \quad u=0$$

$$x=0 \quad u=\varphi t$$

$$u = \varphi(0) \quad \theta \geq t > 0$$

$$u = \varphi(\theta) \quad 2\theta \geq t > \theta$$

$$u = \varphi(\overline{m-1}\theta) \quad m\theta \geq t > \overline{m-1}\theta \quad u = u_1 + u_2 + \dots + u_m$$

$$\chi(x, t) = \frac{2}{\sqrt{\pi}} \int_0^{\infty} e^{-\mu^2} d\mu$$

$$\frac{x}{2a\sqrt{t}}$$

$$u_m = \varphi(\overline{m-1}\theta) \left\{ \chi(x, t - \overline{m-1}\theta) - \chi(x, t - m\theta) \right\}$$

$$u = \sum_{m=1}^{\infty} \varphi(\overline{m-1}\theta) \frac{\chi(x, t - \overline{m-1}\theta) - \chi(x, t - m\theta)}{\theta} \theta$$

$$= \int_0^t \varphi(\lambda) \frac{\partial \chi(x, t-\lambda)}{\partial t} d\lambda$$

$$u = \frac{x}{2a\sqrt{\pi}} \int_0^t \varphi(\lambda) e^{-\frac{x^2}{4a^2(t-\lambda)}} (t-\lambda)^{-\frac{3}{2}} d\lambda$$

$$\left\{ \begin{array}{l} \frac{\partial u}{\partial t} = a \frac{\partial u}{\partial x} \\ t=0 \quad u=0 \\ x=0 \quad u=0 \\ x=c \quad u=\psi(t) \end{array} \right.$$

$$u = \psi(t) \left\{ \frac{x}{c} + \frac{a}{n} \sum_1^{\infty} \frac{(-1)^n}{n} \sin \frac{n\pi x}{c} \right\} + \\ + \frac{2a\pi n}{c^2} \sum_1^{\infty} (-1)^{n-1} n \sin \frac{n\pi x}{c} \int_0^t \psi(\lambda) e^{-a^2 \left(\frac{n\pi}{c}\right)^2 (t-\lambda)} d\lambda$$

Für Kugel mit Berücks. der Ausstrahlung (Kugelfläche)

Schwingungen elastischer Körper

Bewegung der Flüssigkeiten.

Riemann: Schwere,

Elektrizität und Magnetismus.

Hg. v. Hattendorf

Hannover: C. Rimppler 76.

$$X = k_{\mu} \sum_{i=1}^n \frac{m_i (a_i - x)}{r_i^3}$$

$$Y = k_{\mu} \sum \frac{m_i (b_i - y)}{r_i^3}$$

$$Z = \dots$$

Lagrange hat δ auf e δ $\frac{\partial}{\partial x}$ δ auf f δ $\frac{\partial}{\partial x}$:

$$X = \sum m \frac{\partial(\frac{1}{r})}{\partial x} = \frac{\partial V}{\partial x}$$

$$Y = \sum \frac{m}{r}$$

oder δ $\frac{\partial}{\partial x}$
oder δ $\frac{\partial}{\partial x}$

$$Y =$$

$$Z =$$

$$X = \iiint \frac{a-x}{r^3} \rho \, da \, db \, dc$$

$$Y =$$

$$Z =$$

$$V = \iiint \frac{\rho \, da \, db \, dc}{r}$$

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0 \quad \text{oder } \underline{\Delta V = 0}$$

20. Ansatz einer Kugelschale

$$s^2 = x^2 + y^2 + z^2$$

$$V = F(s) \quad \frac{\partial V}{\partial x} = \frac{dV}{ds} \cdot \frac{\partial s}{\partial x}$$

$$\frac{\partial^2 V}{\partial x^2} = \frac{d^2 V}{ds^2} \left(\frac{\partial s}{\partial x} \right)^2 + \frac{dV}{ds} \frac{\partial^2 s}{\partial x^2}$$

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = \frac{d^2 V}{ds^2} + \frac{2}{s} \frac{dV}{ds} = 0$$

$$2s \frac{d^2 V}{ds^2} + 2 \frac{dV}{ds} = 0$$

$$V = \frac{\alpha}{s} + \beta$$

I. Ansatz $s < p$

$$\text{für } s < p: \quad V_0 = 4\pi \int_0^p \rho s ds$$

$$\alpha = 0 \quad \text{ob } s_0 = \infty$$

$$V = 2\pi \rho [p^2 - s^2]$$

II. $s > p$

$$\text{für } s = \infty: \quad V = 0$$

$$\rho = 0$$

$$\frac{1}{s-g} > \frac{1}{2} > \frac{1}{s+g}$$

$$\left| \rho \, d\epsilon \, dt \, d\epsilon \int \right.$$

103

$$\frac{M}{s-g} > V > \frac{M}{s+g}$$

$$\frac{1}{1-\frac{g}{s}} > \frac{\alpha}{11} > \frac{1}{1+\frac{g}{s}}$$

$$\left| \begin{array}{l} s \rightarrow \infty \\ \frac{\alpha}{M} = 1 \end{array} \right.$$

$$V = \frac{M}{s}$$

A. einer homog. Kugel

$$s < g \quad V_1 = \frac{4}{3} \pi \rho s^2 + 2\pi \rho (g^2 - s^2) \\ = 2\pi \rho g^2 - \frac{2}{3} \pi \rho s^2$$

$$s > g \quad V_2 = \frac{4\pi \rho g^3}{3s}$$

$$\left. \frac{\partial V_1}{\partial x} \right|_{s=g} = \left. \frac{\partial V_2}{\partial x} \right|_{s=g}$$

$$\left. \frac{\partial V_1}{\partial x} \right|_{s=g} \neq \left. \frac{\partial V_2}{\partial x} \right|_{s=g}$$

discontinuität
+ Defekt.

$$\frac{\partial V}{\partial x} + \dots = -4\pi \rho$$

Allgem.: V und ihre Deriv. für ein. inneren Punkt

$$V = \iiint \rho \frac{dx dy dz}{r} \quad \text{mit } L = \infty$$

$$a = x + r \sin \vartheta \cos \varphi$$

$$b = y + r \sin \vartheta \sin \varphi$$

$$c = z + r \cos \vartheta$$

$$V = \int_0^{2\pi} d\varphi \int_0^{\pi} \sin \vartheta d\vartheta \int_0^R \rho r dr$$

ebenso x, y, z el. C. Transf.:

$$X = \int_0^{2\pi} \cos \varphi d\varphi \int_0^{\pi} \sin^2 \vartheta d\vartheta \int_0^R \rho r dr$$

$$Y = \int_0^{2\pi} \sin \varphi d\varphi \int_0^{\pi} \sin^2 \vartheta d\vartheta \int_0^R \rho r dr$$

$$Z = \int_0^{2\pi} d\varphi \int_0^{\pi} \cos \vartheta \sin \vartheta d\vartheta \int_0^R \rho r dr$$

$$\therefore x = \frac{\partial V}{\partial a} \text{ etc. } \underline{\text{allgem. } r^2}$$

W-EM $r = 0$ / $\frac{\partial V}{\partial r}$ etc. ρ C. ρ f. ρ f.

$$\frac{\partial V}{\partial r} = \iiint \rho \sin \vartheta d\varphi d\vartheta dr \left\{ \frac{-1 + 3 \sin^2 \vartheta \cos^2 \varphi}{r} \right\}$$

etc. $r=0$ f. ρ f. ρ f. ρ f. ρ f. ρ f.

W-EM $r = 0$ / $\frac{\partial V}{\partial r}$ etc. ρ C. ρ f. ρ f.

$$\sigma = \int_{\Sigma} \rho \frac{dr}{r} \quad \rho = \rho(r, \theta, \phi), \quad \lim_{\epsilon \rightarrow 0}$$

101

er - Transformation:

$$\frac{\partial V}{\partial x} = \iiint \rho \frac{\partial(-\frac{1}{r})}{\partial a} da db dc$$

$$\int \rho \frac{\partial(-\frac{1}{r})}{\partial a} da = -\frac{1}{r} + \int \frac{1}{r} \frac{\partial \rho}{\partial a} da \quad \left. \vphantom{\int} \right\} \omega = \frac{1}{r} \sin \theta \sin \phi$$

$$\frac{\partial V}{\partial x} = \int \frac{\rho}{r} \cos \alpha db dc + \iiint \frac{1}{r} \frac{\partial \rho}{\partial a} da db dc$$

Voraussetz. $\rho = \rho(r, \theta, \phi)$

$\omega = \rho$ Transf. ω / R $\rho = \rho(r, \theta, \phi)$

$\omega < \text{Wert } 0 \text{ } - \pi \text{ Rad. } \epsilon \text{ } \sin \theta \sin \phi$

$$\rho \text{ Transf. } \int \rho = \int \rho \frac{d\omega}{r}$$

$$< \frac{\rho_1}{\epsilon} \int d\omega = < 4\pi \rho_1 \epsilon \quad \left. \vphantom{\int} \right\} \lim_{\epsilon \rightarrow 0} = 0$$

$$\rho \text{ Transf. } \iiint = 0, \quad \lim_{\epsilon \rightarrow 0} [r, R, \omega]$$

$$\rho \text{ Transf. } \frac{\partial V}{\partial x} = \iiint \rho \frac{\partial(-\frac{1}{r})}{\partial a} da db dc \quad \left. \vphantom{\iiint} \right\} \omega = \rho \text{ Transf.}$$

d. Diff:

$$\frac{\partial v}{\partial x^2} = \int \frac{\partial(\frac{1}{r})}{\partial x} \rho \cos \alpha dt + \iiint \frac{\partial(\frac{1}{r})}{\partial x} \frac{\partial \rho}{\partial t} dx dy dz$$

$$\frac{\partial v}{\partial y} =$$

$$\frac{\partial v}{\partial z} =$$

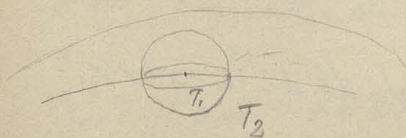
v = diff. d. $\int \rho dx dy dz$ [1, 2, 3, ...]
 für v = retard. unipol. til. line. eff. st.
 u. eff. st. ret. unipol. til. line. eff. st.

Die r. Vekt. v. s. d. s.

v = $\frac{\partial v}{\partial x}$ $\frac{\partial v}{\partial y}$ $\frac{\partial v}{\partial z}$ $\frac{\partial v}{\partial t}$ $\frac{\partial v}{\partial x}$ $\frac{\partial v}{\partial y}$ $\frac{\partial v}{\partial z}$ $\frac{\partial v}{\partial t}$

v $\frac{\partial v}{\partial x}$ = $\frac{\partial v}{\partial x}$ $\frac{\partial v}{\partial y}$ $\frac{\partial v}{\partial z}$ $\frac{\partial v}{\partial t}$

v $\frac{\partial v}{\partial x}$ $\frac{\partial v}{\partial y}$ $\frac{\partial v}{\partial z}$ $\frac{\partial v}{\partial t}$



$$-m \sqrt{R} \epsilon \gamma^0$$

$$v = v_1 + v_2$$

$$|T_2 \epsilon \gamma^0 - m| = \frac{\partial v_2}{\partial x} \rho \epsilon \gamma^0 \text{ u. } \lim d(\frac{\partial v_2}{\partial x}) = 0$$

$$\frac{\partial V_1}{\partial x} = \iiint \frac{a-x}{r^3} \rho \, da \, db \, dc \quad f = \sigma \sqrt{a^2 + b^2 + c^2}$$

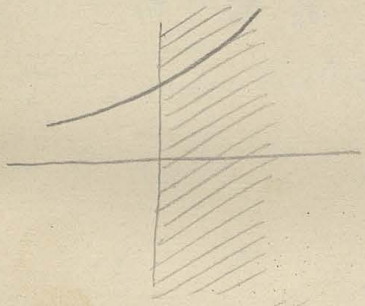
quando o volume $V_1 < \delta$ o ρ \approx ρ_0 \forall ρ \in V_1 \Rightarrow $\rho_0 < \delta$

$$\lim d\left(\frac{\partial V_1}{\partial x}\right) = 0$$

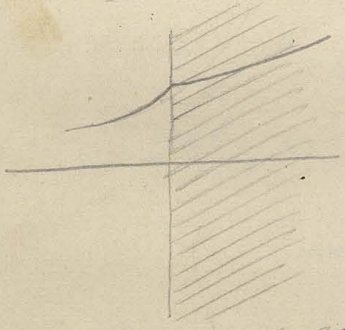
$$\lim d\left(\frac{\partial V}{\partial x}\right) = \lim d\left(\frac{\partial V_1}{\partial x}\right) + \lim d\left(\frac{\partial V_2}{\partial x}\right) = 0$$

\therefore $\frac{\partial V}{\partial x}$ \approx δ deriv. su ρ ρ_0

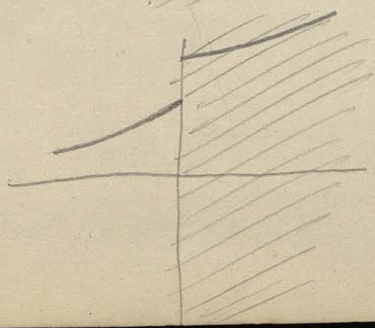
ρ :



$\frac{\partial V}{\partial x}$:



$\frac{\partial V}{\partial x}$:



$$\int \frac{1}{r^2} \cos(\alpha r) d\sigma; \text{ Satz von Gauss}$$

Es gilt $\vec{r} = x\vec{e}_1 + y\vec{e}_2 + z\vec{e}_3$; P komp. e -f. M $r = r(\alpha, \beta, \gamma)$ $\alpha, \beta, \gamma \in [0, 2\pi]$

$$\text{Normale } \vec{N} = r: N = \frac{1}{r^2} \cos(\alpha r) \quad [\alpha r = 1]$$

$$r = (x-a)^2 + (y-b)^2 + (z-c)^2$$

$$\frac{\partial x}{\partial r} = \cos(\alpha r) \quad \frac{\partial y}{\partial r} = \cos(\beta r), \quad \frac{\partial z}{\partial r} = \cos(\gamma r)$$

$$\cos(\alpha r) = \frac{\partial r}{\partial x} \frac{\partial x}{\partial r} + \frac{\partial r}{\partial y} \frac{\partial y}{\partial r} + \frac{\partial r}{\partial z} \frac{\partial z}{\partial r} = \frac{\partial r}{\partial r}$$

$$\int N d\sigma = \int \frac{1}{r^2} \cos(\alpha r) d\sigma$$

$$= \underbrace{4\pi}_{\text{Ueberfl. d. Kgl.}}$$

$$= 0 \text{ Ueberfl. d. Kgl.}$$

$$= 2\pi \cdot \underbrace{r^2}_{\text{Ueberfl. d. Kgl.}} \cdot \frac{1}{r^2} = 2\pi$$

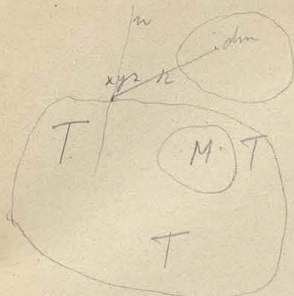
$$\text{U. R. } \rho = \rho_0 \cdot r^2 \quad \text{U. R. } \rho = \rho_0 \cdot r^2$$

$$N = \int \frac{1}{r^2} \cos(\alpha r) d\sigma = \int \frac{1}{r^2} \cos(\alpha r) d\sigma$$

$$\int N d\sigma = \int N d\sigma \int \frac{1}{r^2} \cos(\alpha r) d\sigma =$$

$$\int d\sigma \left\{ \int \frac{1}{r^2} \cos(\alpha r) d\sigma \right\} = 4\pi M$$

M = ...



M = ...

C M x - p - ...

$$\int N d\sigma = 2\pi M$$

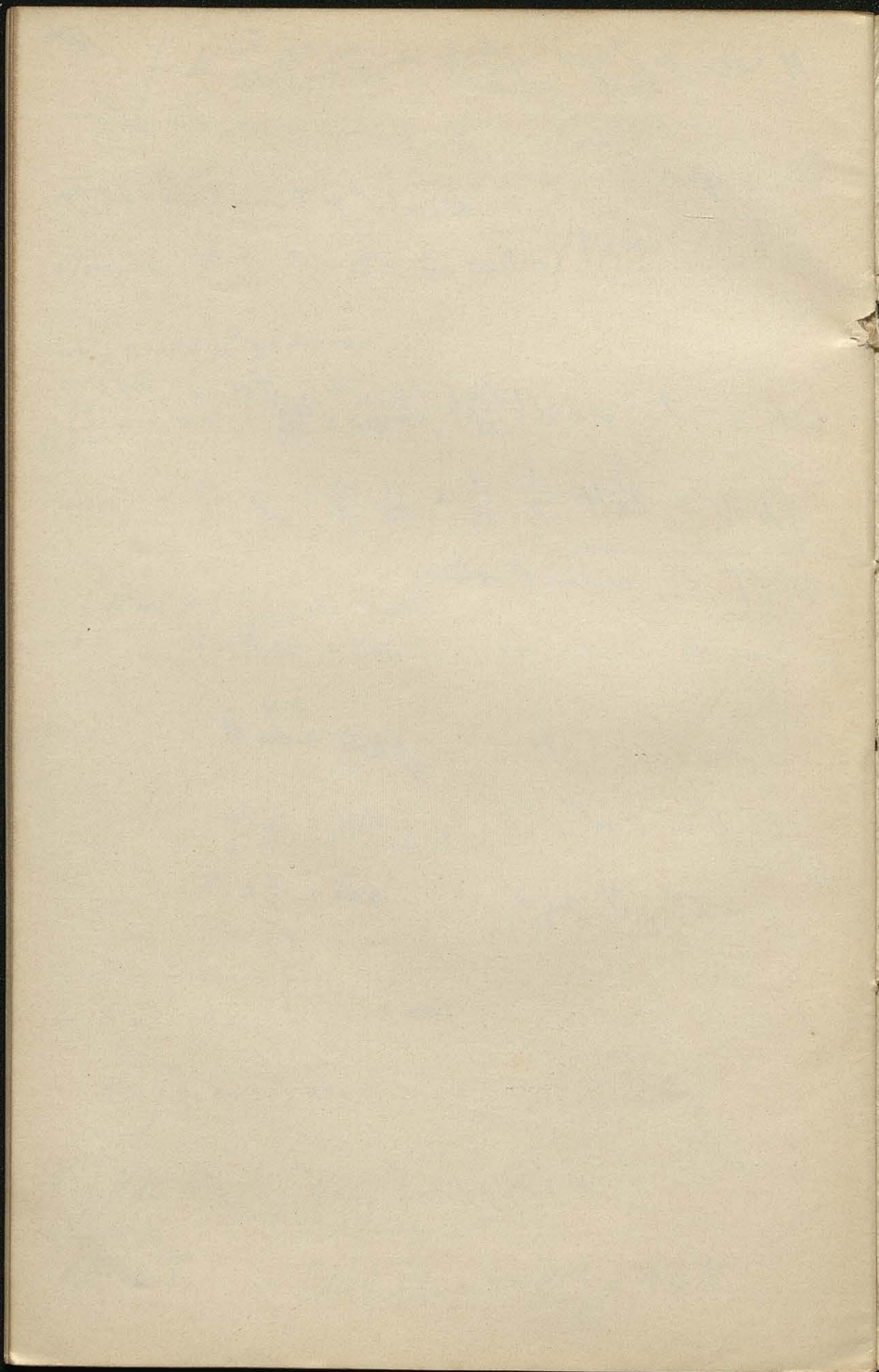
2D. T ~ rechte. Parallelep.

ω 2m M, \leftarrow ... $\int N d\sigma = 4\pi M$

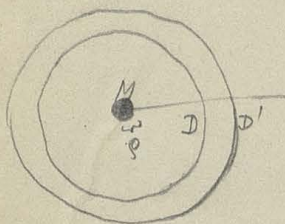
$\sim 2D, M \int N d\sigma = 2\pi M$

ω M ... $\int N d\sigma = \pi M$

ω ... $\int N d\sigma = \frac{1}{2}\pi M$



Verdehnung einer gasförmigen Materie um einen Massenmittelpunkt herum.



$$D = f(r) \quad D' = f(r + \Delta r)$$

$$D : D' = P : P'$$

$$P = \frac{(r + \Delta r)^2 4\pi \cdot P' + \frac{M m k}{r^2}}{4\pi r^2}$$

$$m = 4\pi r^2 dr \cdot \frac{D'}{g}$$

$$P = \frac{(r + \Delta r)^2 P' + M r^2 \Delta r \frac{D' k}{g}}{r^2} \quad \left| \quad P' = P - M \Delta r \frac{D k}{g} \right.$$

$$\frac{D' - D}{\Delta r} : D = \frac{P' - P}{\Delta r} : P$$

$$\frac{D' - D}{\Delta r} = \frac{P' - P}{\Delta r} \frac{D}{P}$$

$$D : D_0 = P : P_0$$

$$\frac{D}{P} = \frac{D_0}{P_0}$$

$$\frac{dD}{dr} = \frac{dP}{dr} \frac{D_0}{P_0}$$

$$\frac{P - P'}{\Delta r} = \frac{P - M \Delta r \frac{D' k}{g} - P'}{\Delta r} = - \frac{M D' k}{g}$$

$$\frac{dD}{dr} = - \frac{D_0}{P_0} \frac{M D k}{g}$$

$$\frac{dP}{dr} = - \frac{M D k}{g} + \frac{D_0}{P_0} \frac{M k}{g} = a$$

$$\frac{dD}{D} = -a dr$$

$$\int D = -a r + b$$

$$\int D_0 = -a r_0 + b$$

$$\int \frac{D}{D_0} = -a(r - r_0)$$

$$D = D_0 e^{-a(r-r_0)} = D_0 e^{\frac{-ar}{l} - \frac{-ar_0}{l}} = c e^{\frac{-ar}{l}}$$

$$c = \frac{D_0}{e^{\frac{-ar_0}{l}}} = D_0 e^{\frac{ar_0}{l}}$$

$$\left. \begin{array}{l} \text{für } M=0: a=0 \\ D=D_0 \end{array} \right\}$$

$$P = \frac{P_0}{D_0} D = P_0 e^{\frac{-ar}{l}}$$

$$G = \int_0^R 4\pi r^2 n D dr$$

$$= 4\pi n \int_0^R e^{-\frac{ar}{l}} r^2 dr$$

$$\int e^{-\frac{ax}{l}} x^2 dx = -\frac{l}{a} e^{-\frac{ax}{l}} x^2 + \frac{2}{a} \int e^{-\frac{ax}{l}} x dx$$

$$= -\frac{l}{a} e^{-\frac{ax}{l}} x^2 + \frac{2}{a} \left\{ -\frac{l}{a} e^{-\frac{ax}{l}} x + \frac{1}{a} \int e^{-\frac{ax}{l}} dx \right\}$$

$$= -\frac{l}{a} e^{-\frac{ax}{l}} x^2 + \frac{2}{a} \left\{ -\frac{l}{a} e^{-\frac{ax}{l}} x - \frac{1}{a^2} e^{-\frac{ax}{l}} \right\} = -l e^{-\frac{ax}{l}} \left[\frac{x^2}{a} + \frac{2x}{a^2} + \frac{2}{a^3} \right]$$

$$\begin{aligned} \beta &= -4cr e^{-ar} \left[\frac{r^2}{a} + \frac{2r}{a^2} + \frac{2}{a^3} \right] \Big|_0^R \\ &= 4cr e^{-ar} \left[\frac{r^2}{a} + \frac{2r}{a^2} + \frac{2}{a^3} \right] \Big|_0^R \end{aligned}$$

in eine Reihe entwickelt:

$$\begin{aligned} \beta &= 4cr \left[1 - ar + \frac{a^2 r^2}{1 \cdot 2} - \frac{a^3 r^3}{1 \cdot 2 \cdot 3} + \frac{a^4 r^4}{1 \cdot 2 \cdot 3 \cdot 4} - \frac{a^5 r^5}{5!} \right] \left[\frac{r^2}{a} + \frac{2r}{a^2} + \frac{2}{a^3} \right] \\ &= 4cr \left[\frac{r^2}{a} - r^3 + \frac{ar^4}{2} - \frac{a^2 r^5}{3!} + \frac{a^3 r^6}{4!} - \frac{a^4 r^7}{5!} + \dots \right. \\ &\quad \left. + \frac{2r}{a^2} - \frac{2r^2}{a} + \frac{2r^3}{2} - \frac{2ar^4}{3!} + \frac{2a^2 r^5}{4!} - \frac{2a^3 r^6}{5!} + \dots \right. \\ &\quad \left. + \frac{2}{a^3} - \frac{2r}{a^2} + \frac{2r^2}{2a} - \frac{2r^3}{3!} + \frac{2ar^4}{4!} - \frac{2a^2 r^5}{5!} + \dots \right] \\ &= 4cr \left[\frac{2}{a^3} - \frac{2r^3}{3!} + ar^4 \left(\frac{1}{2} - \frac{2}{3!} + \frac{2}{4!} \right) - a^2 r^5 \left(\frac{1}{3!} - \frac{2}{4!} + \frac{2}{5!} \right) + \dots \right] \end{aligned}$$

Beim Substituieren der Grenzwerte fällt das erste Glied heraus, also auch:

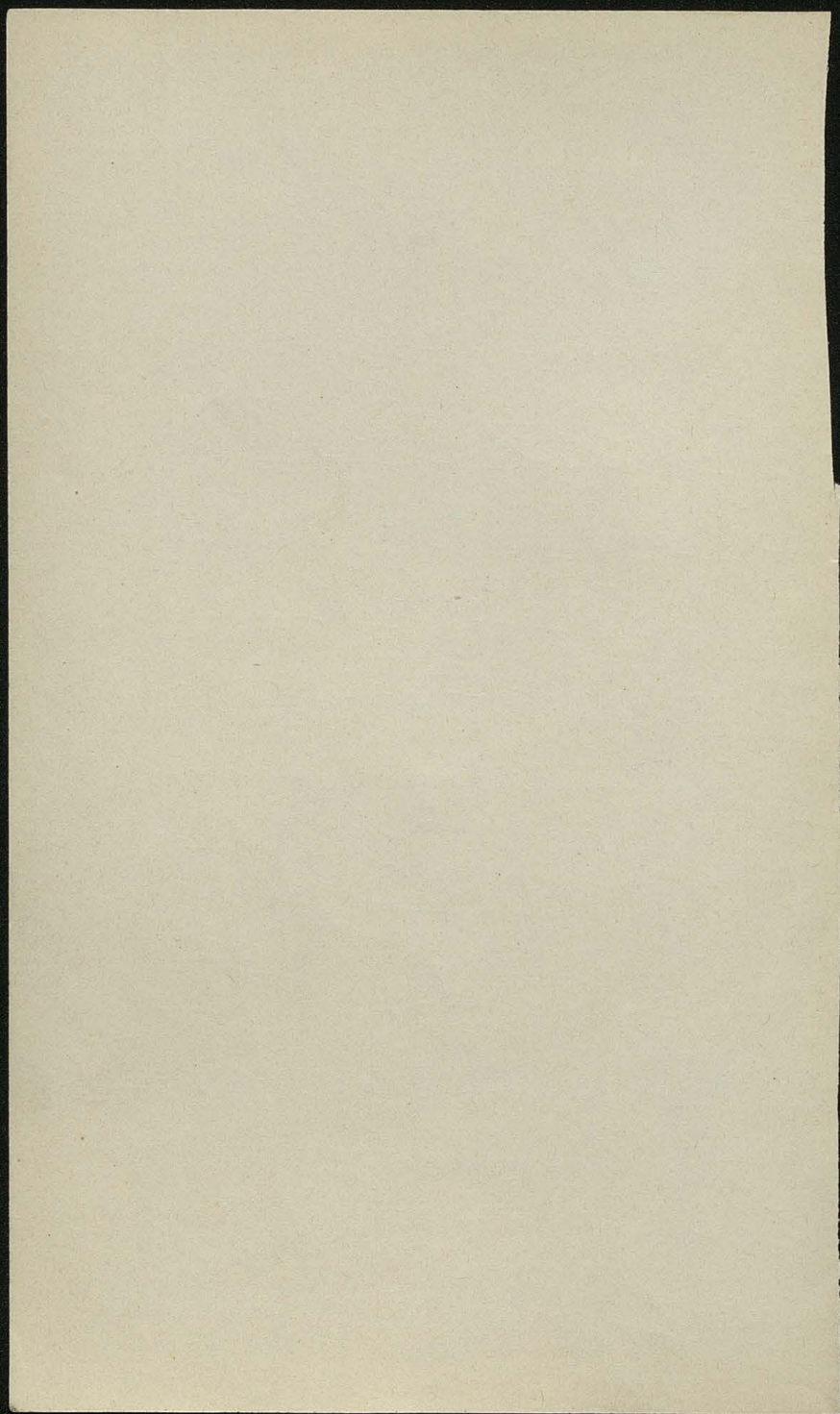
$$\begin{aligned} \beta &= 4cr \left[-\frac{r^3}{3} + \frac{ar^4}{4} - \frac{a^2 r^5}{10} + \frac{a^3 r^6}{36} - \dots \right] \Big|_0^R \\ &= 4cr \left[\frac{r^3}{3} - \frac{ar^4}{4} + \frac{a^2 r^5}{10} - \frac{a^3 r^6}{36} - \dots \right] \Big|_0^R \end{aligned}$$

In der Tat gilt dies für $M=0$, $a=0$: $\beta = \frac{4cr}{3} [R^3 - \rho^3]$

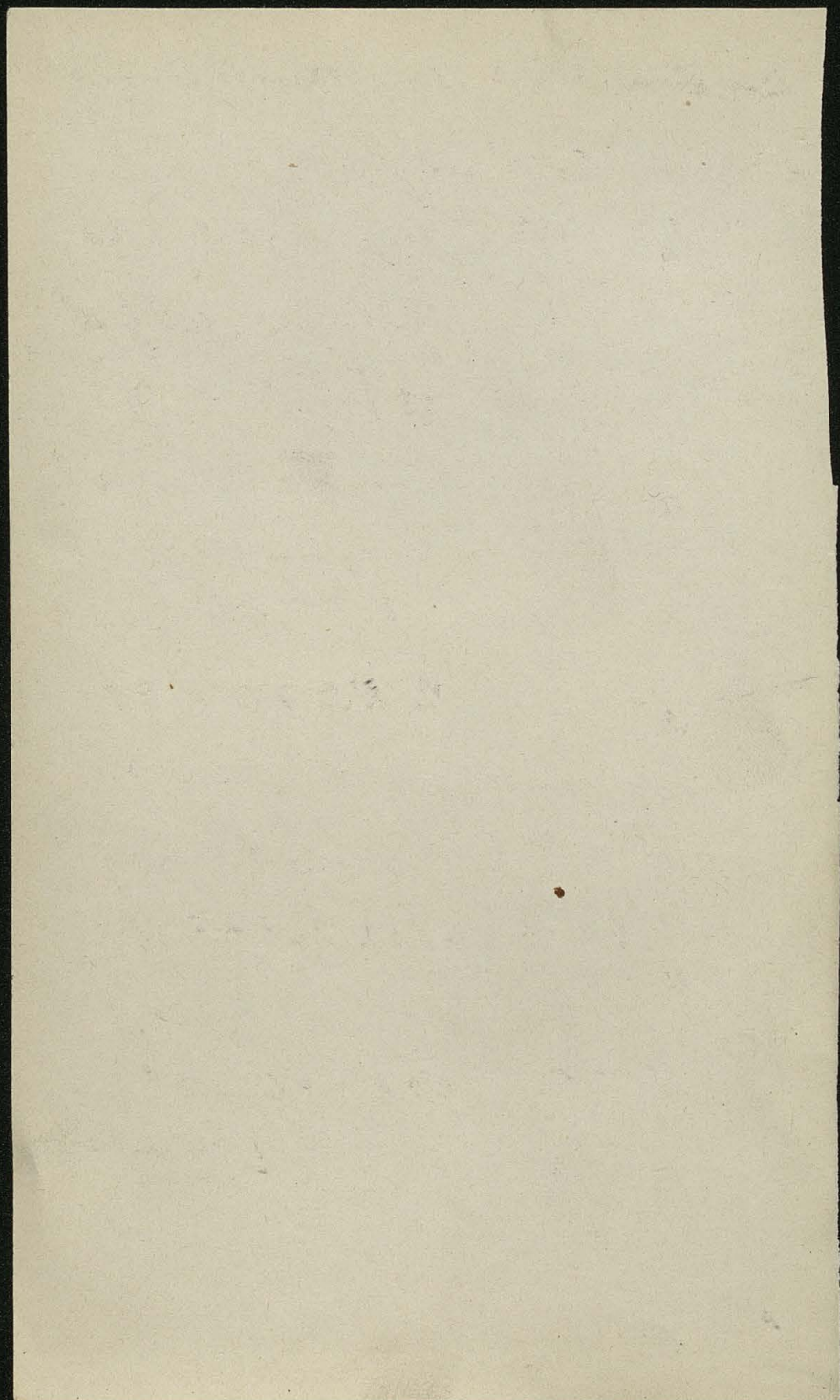
Der Überschuss des Gesamtgewichtes über jenes,

wobei $M=0$, ist:

$$\beta - \beta_0 = G = 4cr \left[\dots \right]$$

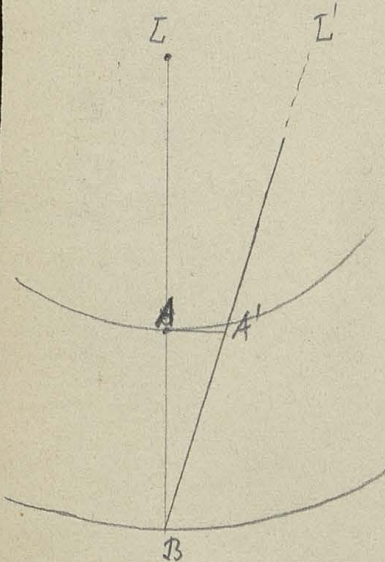


110



Aberration $\approx \frac{v}{c} \sin \theta$ in theorie $\approx \frac{v}{c} \sin \theta$ in re

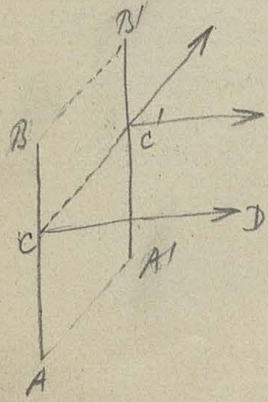
in $\approx \frac{v}{c} \sin \theta$ in re



theorie $\approx \frac{v}{c} \sin \theta$ in re
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in re $\approx \frac{v}{c} \sin \theta$ in re $\approx \frac{v}{c} \sin \theta$ in re



$\approx \frac{v}{c} \sin \theta$ in re
 $\approx \frac{v}{c} \sin \theta$ in re

[Quelle: ...]

Gleichgewichtverteilung von Elektrizität auf
 einem System von 2 einander berührenden Kugeln

2 Kugeln mit Radien a und b , Abstand c
 geladene Kugel mit Ladung Q , andere ungeladet

oder

$$\frac{Q}{a} = \frac{q}{b} = \frac{Q+q}{a+b}$$

$$Q+q = \frac{QaC}{a+b} = 10^9 \text{ Q.V.}$$

$$= \frac{C}{a+b}$$

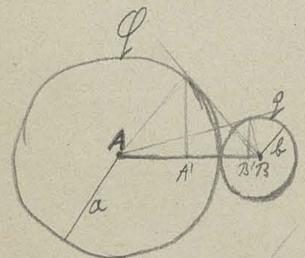
oder

$$Q = \frac{aC}{a+b}$$

$$q = \frac{bC}{a+b}$$

$$Q:q = a:b$$

$a < b$:



oder Q und q induc. geladen
 geladene Kugel Q auf A und B
 $Q+q = Q = 10^9 \text{ Q.V.}$

oder Q und q induc. geladen

$$Q \text{ auf } B = \frac{Qb}{a+b} + \frac{qa}{(a+b)(a+b - \frac{a^2}{a+b})}$$

$$Q \text{ auf } A = \frac{Qa}{a+b}$$

oder Q und q induc. geladen

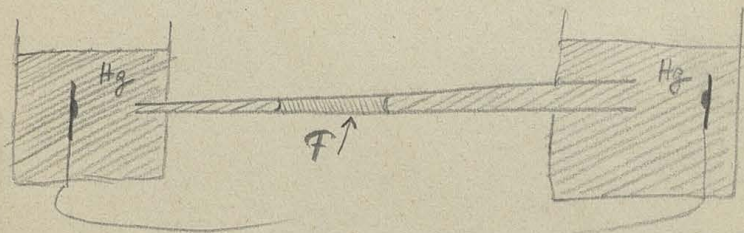
$$Q+q = \frac{Q}{a} + \frac{Q}{a+b} + \frac{q}{b} + \frac{q}{a+b}$$

$$Qab + Qb^2 + qab = qa^2 + qb^2$$

~~$$P: q = a^2: b^2$$~~

$$C = P + q + \frac{Pb + qa}{a+b} = \frac{P(a+b) + q(2a+b)}{a+b}$$

Apparat zur Calibration von Röhren auf elektrischem Wege (Widerstandsbestimmung)

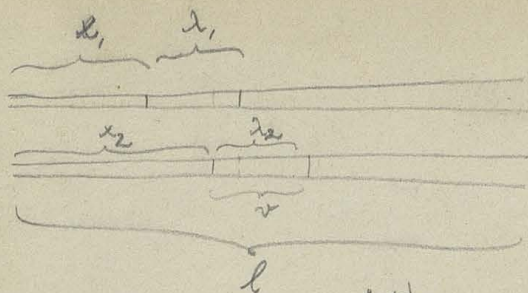


In die Röhre wird ein Stück einer schlecht leitenden Flüssigkeit eingeblasen F ; sonst ist die Röhre ganz mit (20%) Quecksilber Hg angefüllt. Falls nun der Querschnitt ungleich ist, so ändert sich auch die Länge der schlecht leitenden Flüssigkeitssäule um aus beiden Gründen der Widerstand.

$$W = \rho \cdot \frac{l}{S}$$

$$\Sigma W = W = \rho \int \frac{dl}{S}$$

ρ ist dann das Ohm'sche Gesetz noch genau gültig: wenn der Querschnitt verschieden ist



$$W_1 = \alpha \int_0^{x_1} \frac{dx}{q} + \beta \int_{x_1}^{x_1 + \lambda_1} \frac{dx}{q} + \alpha \int_{x_1 + \lambda_1}^l \frac{dx}{q}$$

$$W_2 = \alpha \int_0^{x_2} \frac{dx}{q} + \beta \int_{x_2}^{x_2 + \lambda_2} \frac{dx}{q} + \alpha \int_{x_2 + \lambda_2}^l \frac{dx}{q}$$

$$x_2 = x_1 + \xi$$

$$W_1 - W_2 = -\alpha \int_{x_1}^{x_2} \frac{dx}{q} + \beta \left[\int_{x_1}^{x_2} \frac{dx}{q} - \int_{x_1 + \lambda_1}^{x_2 + \lambda_2} \frac{dx}{q} \right] + \alpha \int_{x_1 + \lambda_1}^{x_2 + \lambda_2} \frac{dx}{q}$$

$$= (\beta - \alpha) \left[\int_{x_1}^{x_2} \frac{dx}{q} - \int_{x_1 + \lambda_1}^{x_2 + \lambda_2} \frac{dx}{q} \right]$$

$$v = \text{const.} = \int_{x_1}^{x_1 + \lambda_1} q dx = \int_{x_2}^{x_2 + \lambda_2} q dx = \lambda_2 \langle q \rangle_m$$

$$\cancel{W_1 = W_2 = (\beta - \alpha)}$$

Wenn der Widerstand α gegen den spec
Widerstand β zu vernachlässigen ist, so
kann man setzen:

$$W_1 = \beta \int_{x_1}^{x_1 + \lambda_1} \frac{dx}{r} = \frac{\beta \lambda_1}{(r_1)}$$

$$(r) = \sqrt{d}$$

$$W_2 = \beta \int_{x_2}^{x_2 + \lambda_2} \frac{dx}{r} = \frac{\beta \lambda_2}{(r_2)}$$

$$\frac{W_1}{W_2} = \frac{\lambda_1 (r_2)}{\lambda_2 (r_1)}$$

$$v = \lambda_1 (r_1) = \lambda_2 (r_2)$$

$$= \left(\frac{\lambda_1^2}{\lambda_2} \right) = \left[\frac{(r_2)}{(r_1)} \right]^2$$

$$\frac{\lambda_1}{\lambda_2} = \frac{(r_2)}{(r_1)}$$

$$\frac{(r_2)}{(r_1)} = \sqrt{\frac{W_1}{W_2}}$$

Eine genauere Formel erhält man wenn man
auch den Widerstand W_0 = ohne den Tropfen

der schlecht bestehenden Flüssigkeit misst; dann
~~1 - mit Formelabgang von λ gegen λ :~~

$$W_1 = \frac{\rho \lambda_1}{(\rho_1)} + W_0$$

$$\left(\frac{\rho_2}{\rho_1}\right) = \sqrt{\frac{W_1 - W_0}{W_2 - W_0}}$$

$$W_2 = \frac{\rho \lambda_2}{(\rho_2)} + W_0$$

Beweis:

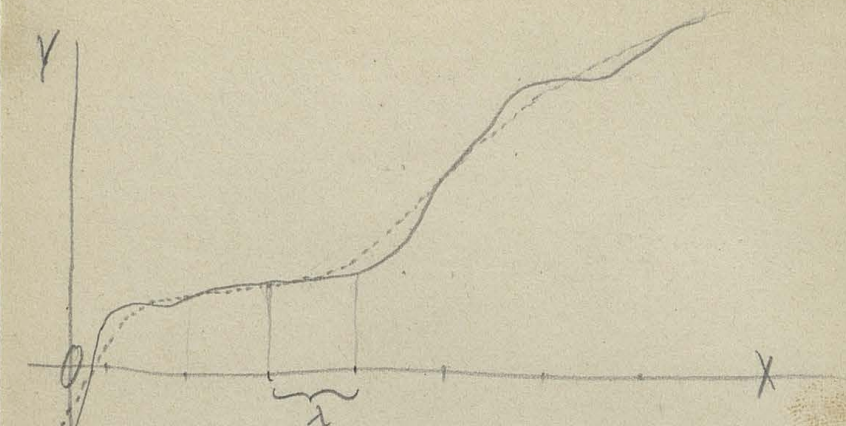
~~Nachgewiesen: wenn α und β bekannt.~~

$$W_1 = \alpha \int_0^{\lambda} \frac{dx}{\rho} - \alpha \int_{x_1}^{x_1 + \lambda_1} \frac{dx}{\rho} + \beta \int_{x_1}^{x_1 + \lambda_1} \frac{dx}{\rho}$$

$$= W_0 + (\beta - \alpha) \int_{x_1}^{x_1 + \lambda_1} \frac{dx}{\rho} = W_0 + (\beta - \alpha) \frac{\lambda_1}{(\rho_1)}$$

$$W_2 = W_0 + (\beta - \alpha) \frac{\lambda_2}{(\rho_2)}$$

$$\frac{W_1 - W_0}{W_2 - W_0} = \frac{\lambda_1}{\lambda_2} \frac{(\rho_2)}{(\rho_1)} = \left[\frac{(\rho_2)}{(\rho_1)}\right]^2$$



$f(x)$

$$y = f(x)$$

$$y = \varphi(x)$$

$$\eta = \frac{1}{\lambda} \int_{x - \frac{\lambda}{2}}^{x + \frac{\lambda}{2}} f(x) dx$$

Sehr oft bei physikal. Körper
wird eine $\varphi(x)$ bestimmt,
während $f(x)$ gesucht wird;

gine stellt den arithm. Mittelwert in einem
festen Intervalle dar. z.B. Spectrallinien

[1]. Übereinanderlagern der Spaltbilder

2). Bei Untersuchung mit Bolometer, Deckedess
Calibrieren von Röhren, Temperatur-Messung
mittels thermom. etc.]

Wie findet man $f(x)$?

$$\varphi(x) = \frac{1}{\lambda} \int_{x-\frac{\lambda}{2}}^{x+\frac{\lambda}{2}} f(x) dx = \frac{1}{\lambda} \int_{x-\frac{\lambda}{2}}^{x+\frac{\lambda}{2}} f(z) dz$$

$$\frac{d\varphi}{dx} = \frac{1}{\lambda} \left[f\left(x+\frac{\lambda}{2}\right) - f\left(x-\frac{\lambda}{2}\right) \right]$$

$$= \frac{1}{\lambda} \left[f(x) + \frac{\lambda}{2} f'(x) + \frac{\lambda^2}{4} \frac{1}{2} f''(x) + \frac{\lambda^3}{8} \frac{1}{2 \cdot 3} f'''(x) + \dots - f(x) + \dots + \dots \right]$$

$$= f'(x) + \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} f''(x) + \frac{\lambda^4}{16} \frac{1}{5!} f^{IV}(x) + \dots$$

$$\lambda \frac{d\varphi}{dx} = \lambda f'(x) + \frac{\lambda^3}{4} \frac{1}{2 \cdot 3}$$

$$\varphi(x) = \int \frac{d\varphi}{dx} dx = f(x) + \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} f''(x) + \frac{\lambda^4}{16} \frac{1}{5!} f^{IV}(x) + \dots$$

I Näherungs-methode:

$$f_0(x) = \varphi(x) - \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} \cdot \varphi''(x) - \frac{\lambda^4}{16} \frac{1}{5!} \varphi^{IV}(x) - \frac{\lambda^6}{64 \cdot 7!} \varphi^{VI}(x) - \dots$$

$$f_2(x) = \varphi(x) - \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} f_0''(x) - \frac{\lambda^4}{16} \frac{1}{5!} f_0^{IV}(x) - \dots$$

$$f_3(x) = \varphi(x) - \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} f_2''(x) - \dots \quad \text{etc.}$$

oder:

$$f_2(x) = \varphi(x) - \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} \left[\varphi''(x) - \frac{\lambda^2}{4} \frac{1}{2 \cdot 3} \varphi^{IV}(x) - \frac{\lambda^4}{16} \frac{1}{5!} \varphi^{VI}(x) - \dots \right]$$

$$\begin{aligned}
 & -\frac{\lambda^4}{16} \cdot \frac{1}{5!} \varphi^{IV}(x) - \frac{\lambda^2}{4} \cdot \frac{1}{2 \cdot 3} \varphi^{VI}(x) - \frac{\lambda^4}{16} \cdot \frac{1}{5!} \varphi^{VIII}(x) \\
 & - \frac{\lambda^6}{64} \cdot \frac{1}{7!} \left[\varphi^{VI}(x) - \frac{\lambda^2}{4} \cdot \frac{1}{2 \cdot 3} \varphi^{VIII}(x) \right]
 \end{aligned}$$

$$\begin{aligned}
 & = \varphi(x) - \frac{\lambda^2}{4} \cdot \frac{1}{2 \cdot 3} \varphi''(x) + \left[\frac{\lambda^4}{16} \left(\frac{1}{2 \cdot 3} \right)^2 - \frac{\lambda^4}{16} \cdot \frac{1}{5!} \right] \varphi^{IV}(x) \\
 & + \frac{\lambda^6}{64} \left[\frac{1}{2 \cdot 3} \cdot \frac{1}{5!} + \frac{1}{5!} \cdot \frac{1}{2 \cdot 3} - \frac{1}{7!} \right] \varphi^{VI}(x)
 \end{aligned}$$

etc.

mit Beschränkung auf 4te Potenzen

$$\left[\text{Coefficient } 6^{\text{ter}} = \frac{1}{322560} ! \right]$$

$$f_2(x) = \varphi(x) - \frac{\lambda^2}{24} \varphi''(x) + \frac{7\lambda^4}{5760} \varphi^{IV}(x) \leftarrow$$

$$f_3(x) = \varphi(x) - \frac{\lambda^2}{4} \cdot \frac{1}{2 \cdot 3} \left[\varphi''(x) - \frac{\lambda^2}{24} \varphi^{IV}(x) + \frac{7}{5}
 \right.$$

$$\left. + \frac{\lambda^4}{16} \cdot \frac{1}{5!} \varphi^{IV}(x) \right]$$

$$= \varphi(x) - \frac{\lambda^2}{24} \varphi''(x) + \frac{7\lambda^4}{5760} \varphi^{IV}(x) =$$

$$d) = \text{problem} = \varphi(x) - \frac{\lambda^2}{24} \varphi''(x) + \frac{7\lambda^4}{5760} \varphi^{(4)}(x)$$

$$\varphi\left(x + \frac{\lambda}{2}\right) + \varphi\left(x - \frac{\lambda}{2}\right) = \frac{1}{\lambda} \int_x^{x+\lambda} f(x) dx + \frac{1}{\lambda} \int_{x-\lambda}^x f(x) dx = \frac{1}{\lambda} \int_{x-\lambda}^{x+\lambda} f(x) dx$$

$$\frac{d}{dx} \left[\varphi\left(x + \frac{\lambda}{2}\right) + \varphi\left(x - \frac{\lambda}{2}\right) \right] = \frac{1}{\lambda} \left[f(x+\lambda) - f(x-\lambda) \right]$$

Coordinaten-Verschiebung

[e Index $\nu_2 \neq \nu_1$ (gr!)]

$$\varphi(x) = \frac{1}{\lambda} \int_x^{x+\lambda} f(x) dx$$

$$\frac{d\varphi(x)}{dx} = \frac{1}{\lambda} \left[f(x+\lambda) - f(x) \right]$$

$$\varphi(x+\lambda) = \frac{1}{\lambda} \left[f(x+2\lambda) - f(x+\lambda) \right]$$

$$\frac{d}{dx} [\varphi(x) + \varphi(x+\lambda)] = \frac{1}{\lambda} \left[f(x+2\lambda) - f(x) \right]$$

$$= 2 \frac{d\varphi_2(x)}{dx}$$

$$\varphi_2(x) = \frac{1}{2\lambda} \int_x^{x+2\lambda} f(x) dx$$

$$\varphi_2(x) = \frac{1}{2} [\varphi_1(x) + \varphi_1(x+\lambda)] + C_1$$

$$\varphi_1(x) = \frac{1}{2} [\varphi_{\frac{1}{2}}(x) + \varphi_{\frac{1}{2}}(x+\lambda)] + C_2$$

$$\varphi_{\frac{1}{2}}(x+\lambda) = \frac{1}{2} [\varphi_{\frac{1}{2}}(x+\lambda) + \varphi_{\frac{1}{2}}(x+2\lambda)] + C_2$$

$$\varphi_{\frac{1}{2}}(x) = \frac{1}{2} \left[\frac{1}{2} \varphi_{\frac{1}{2}}(x) + \varphi_{\frac{1}{2}}(x+\lambda) + \frac{1}{2} \varphi_{\frac{1}{2}}(x+2\lambda) \right] + C_1 + C_2$$

$$= \frac{1}{4} \varphi_{\frac{1}{2}}(x) + \frac{1}{2} \varphi_{\frac{1}{2}}(x+\lambda) + \frac{1}{4} \varphi_{\frac{1}{2}}(x+2\lambda) + C_1 + C_2$$

etc. etc.

$$f(x+\lambda) = \lambda \frac{d\varphi(x)}{dx} + f(x)$$

$$f(\lambda) = \lambda \frac{d\varphi(0)}{dx} + f(0)$$

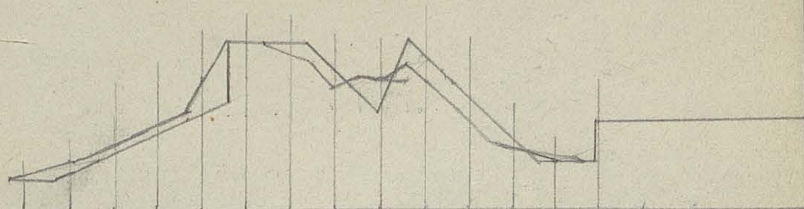
$$f(2\lambda) = \lambda \frac{d\varphi(2\lambda)}{dx} + f(2\lambda)$$

$$= \lambda \left[\frac{d\varphi(2\lambda)}{dx} + \frac{d\varphi(0)}{dx} \right] + f(0)$$

$$f(3\lambda) = \lambda \left[\frac{d\varphi(2\lambda)}{dx} + \frac{d\varphi(\lambda)}{dx} + \frac{d\varphi(0)}{dx} \right] + f(0)$$

etc. wenn also ein Werth $f(0)$ bekannt ist
so können daraus alle anderen gefunden werden.

MF



Bewegung eines Punktes welcher von einem fixen Centrum ausgeh. perp. der 3 Pot. der Entfernung angesetzt wird.

$$\frac{d^2x}{dt^2} = -k \frac{x}{r^4} \quad \left| \quad \frac{dx}{dt} \right.$$

$$\frac{d^2y}{dt^2} = -k \frac{y}{r^4} \quad \left| \quad \frac{dy}{dt} \right.$$

$$\frac{d}{dt} \left[\left(\frac{dx}{dt} \right)^2 + \left(\frac{dy}{dt} \right)^2 \right] = -\frac{k}{r^3} \frac{dr}{dt}$$

$$= \frac{d}{dt} \left[\left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\phi}{dt} \right)^2 \right] = \frac{k}{2} \frac{d}{dt} \left(\frac{1}{r^2} \right)$$

$$\left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\phi}{dt} \right)^2 = \frac{k}{2r^2} + A$$

$$A = \left(\frac{dr}{dt} \right)_0^2 + r_0^2 \left(\frac{d\phi}{dt} \right)_0^2 - \frac{k}{2r_0^2}$$

$$r^2 \frac{d\phi}{dt} = 2c$$

$$\left(\frac{dr}{dt} \right)^2 + \frac{4c^2}{r^2} = \frac{k}{2r^2} + A$$

$$\left(\frac{dr}{dt}\right)^2 = \frac{k - \rho c^2}{2r^2} + A$$

118

$$\frac{dr}{\frac{k - \rho c^2}{2r^2} + A} = dt$$

$$\frac{k - \rho c^2}{2r^2} + A$$

$$\frac{r dr \sqrt{2}}{\sqrt{k - \rho c^2 + 2Ar^2}} = dt$$

$$t = \frac{1}{A\sqrt{2}} \sqrt{k - \rho c^2 + 2Ar^2} + \beta$$

$$4A^2(t - \beta)^2 = k - \rho c^2 + 2Ar^2$$

$$r = \sqrt{\frac{2A^2(t - \beta)^2 + \rho c^2 - k}{2A}}$$

$$= \frac{1}{2} \sqrt{2A(t - \beta)^2 + \frac{\rho c^2 - k}{A}}$$

$$\frac{4c^2}{r^4} \left(\frac{dr}{dt}\right)^2 = \frac{k - \rho c^2}{2r^2} + A$$

$$\rho c^2 \left(\frac{dr}{dt}\right)^2 = \frac{k - \rho c^2}{\rho c^2} r^2 + \frac{2Ar^4}{4c^2}$$

$$\frac{dr}{\sqrt{(k - \rho c^2)r^2 + 2Ar^4}} = \frac{dt}{2c\sqrt{2}}$$

$$2c = r^2 \frac{df}{dr}$$

$$= r^2 \frac{df}{dr} \frac{dr}{dt}$$

$$= \frac{r^2 2c\sqrt{2}}{\sqrt{(k-\rho c^2)r^2 + 2Ar^4}}$$

$$\frac{\sqrt{k-\rho c^2} + \sqrt{k-\rho c^2 + 2Ar^2}}{r\sqrt{2}}$$

$$\frac{\varphi}{2c\sqrt{2}} = \int \frac{dr}{r\sqrt{(k-\rho c^2) + 2Ar^2}} + C$$

I. $k - \rho c^2 > 0$

$$\frac{\varphi}{2c\sqrt{2}} = -\frac{1}{\sqrt{k-\rho c^2}} \log \left[\frac{\sqrt{k-\rho c^2} + \sqrt{k-\rho c^2 + 2Ar^2}}{r} \right]$$

$$(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}$$

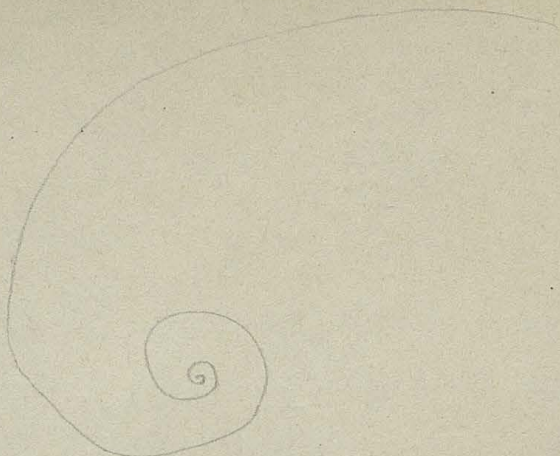
$$= \frac{\sqrt{k-\rho c^2} + \sqrt{k-\rho c^2 + 2Ar^2}}{r}$$

$$k - \rho c^2 + 2Ar^2 = k - \rho c^2 + 2r \sqrt{k-\rho c^2} e^{(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}} + r^2 e^{2(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}}$$

$$r_1 = 0$$

$$(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}$$

$$r_2 = \frac{2\sqrt{k-\rho c^2} e^{(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}}}{2A - e^{2(C-\varphi) \frac{\sqrt{k-\rho c^2}}{2c\sqrt{2}}}}$$



II. $k - \rho c^2 < 0$

$$\frac{\varphi}{\rho c \sqrt{2}} = -\frac{1}{\sqrt{\rho c^2 - k}} \arcsin \frac{1}{2} \sqrt{\frac{\rho c^2 - k}{2A}} + \text{Const.}$$

$$\frac{1}{2} \sqrt{\frac{\rho c^2 - k}{2A}} = \sin \frac{(C - \varphi) \sqrt{\rho c^2 - k}}{2c \sqrt{2}}$$

$$r = \sqrt{\frac{\rho c^2 - k}{2A}} \frac{1}{\frac{\sin \frac{(C - \varphi) \sqrt{\rho c^2 - k}}{2c \sqrt{2}}}{2c \sqrt{2}}}$$

$t = \infty \quad r = \infty$

II. $C - \varphi_1 = 0$

$$\varphi_1 = C$$

$$C - \varphi_2 = \pi$$

$$\varphi_2 = C - \pi$$

$$I. \quad 2A - e^{(C-\varphi) \frac{\sqrt{k-\rho c^2}}{c\sqrt{2}}} = 0$$

$$(C-\varphi) \frac{\sqrt{k-\rho c^2}}{c\sqrt{2}} = \log 2A$$

$$\varphi = C - \frac{c\sqrt{2}}{\sqrt{k-\rho c^2}} \log(2A)$$

$$t=0 \quad r = \sqrt{A^2 \Omega^2 + \frac{\rho c^2 - k}{2A}}$$

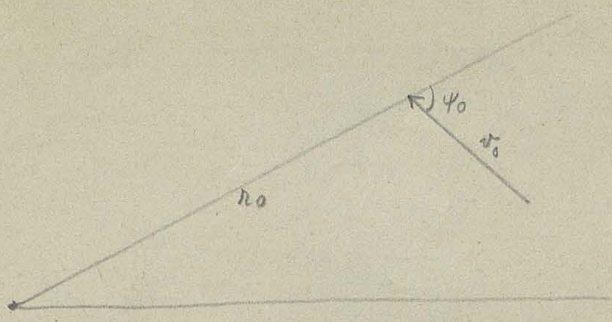
$$I. \quad \varphi = -\frac{2c\sqrt{2}}{\sqrt{k-\rho c^2}} \log \left[\frac{\sqrt{k-\rho c^2} + \sqrt{2A\Omega}}{\sqrt{A^2 \Omega^2 + \frac{\rho c^2 - k}{2A}}} \right] + C$$

II.

$$\varphi = -\frac{2c\sqrt{2}}{\sqrt{\rho c^2 - k}} \arcsin \left(\frac{\sqrt{\rho c^2 - k}}{2A} \frac{1}{\sqrt{A^2 \Omega^2 + \frac{\rho c^2 - k}{2A}}} \right) + C$$

$$= -\frac{2c\sqrt{2}}{\sqrt{\rho c^2 - k}} \arcsin \left(\sqrt{\frac{\rho c^2 - k}{2A^2 \Omega^2 + \rho c^2 - k}} \right) + C$$

$$= -\frac{2c\sqrt{2}}{\sqrt{\rho c^2 - k}} \arcsin \left[\sqrt{\frac{1}{\frac{2A^2 \Omega^2}{\rho c^2 - k} + 1}} \right] + C$$



$\therefore t=0$

$$\underline{\varphi=0, \quad r_0, \quad v_0, \quad \varphi_0}$$

$$\left(\frac{dr}{dt}\right)_0 = -v_0 \cos \varphi_0$$

$$\left(\frac{d\varphi}{dt}\right)_0 = \frac{1}{r_0} v_0 \sin \varphi_0$$

$$A = (-v_0 \cos \varphi)^2 + v_0^2 \sin^2 \varphi - \frac{k}{2r_0^2}$$

$$\underline{A = v_0^2 - \frac{k}{2r_0^2}}$$

$$2C = r_0^2 \left(\frac{d\varphi}{dt}\right)_0$$

$$\underline{C = \frac{r_0 v_0 \sin \varphi_0}{2}}$$

$$A > 0 \quad - \frac{\sqrt{k - 2C^2 + 2Ar_0^2}}{Ar_0^2} = \frac{B}{M}$$

$$A = - \frac{\sqrt{k - \rho c v + 2 v_0^2 n_0^2 - k}}{\left(v_0^2 - \frac{k}{2 n_0^2}\right) \sqrt{2}}$$

$$B = - \frac{\sqrt{v_0^2 n_0^2 - 4c^2}}{v_0^2 - \frac{k}{2 n_0^2}}$$

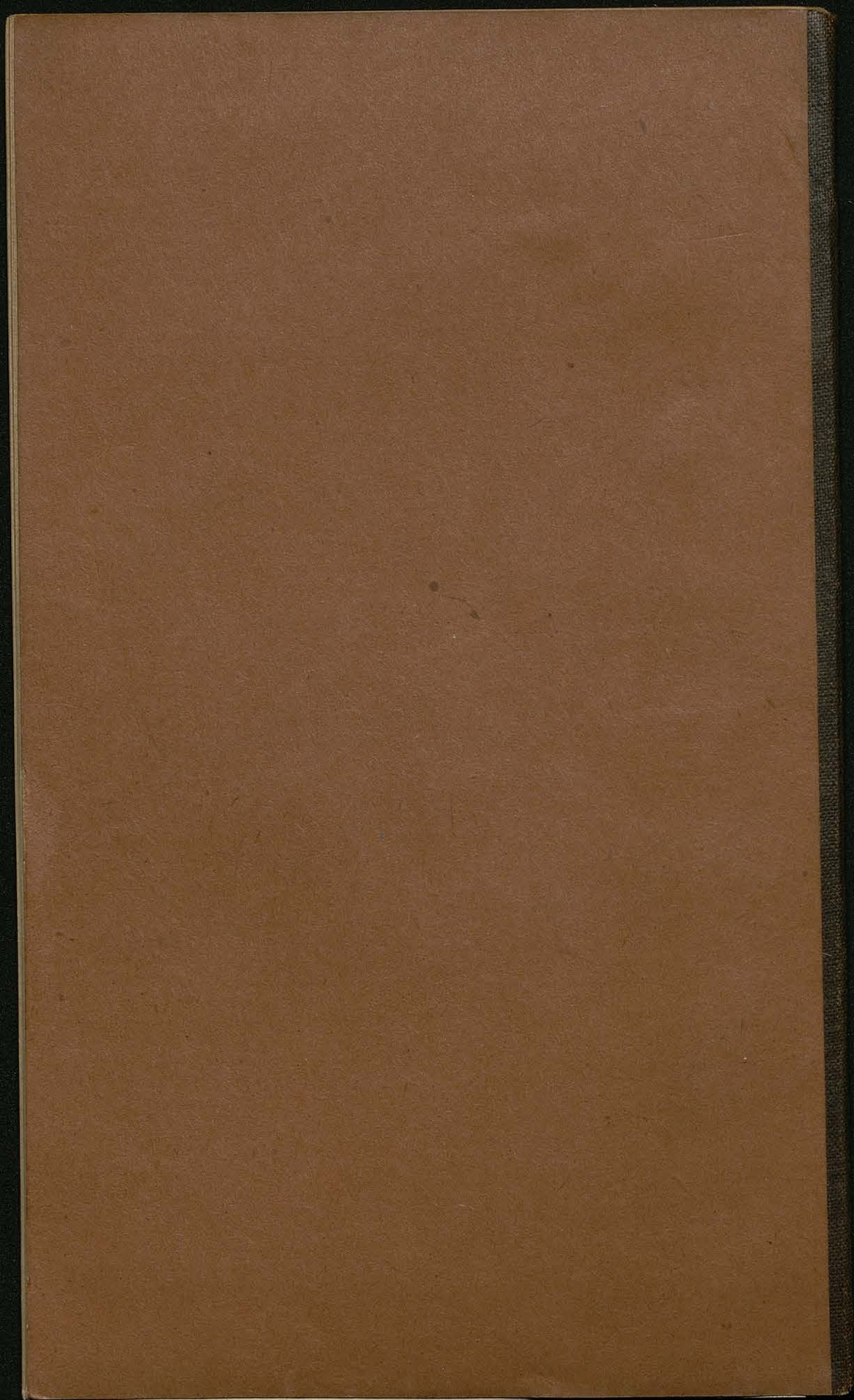
$$C = \frac{2c\sqrt{2}}{\rho c v - k} \arcsin \left[\frac{1}{2 \frac{4 n_0^4 (v_0^2 n_0^2 - 4c^2)}{(2 n_0^2 v_0^2 - k)^2} + 1} \right]$$

~~$$= \frac{2c\sqrt{2}}{\rho c v - k} \arcsin \left[\frac{1}{2 v_0^2 n_0^2 \cos^2 \varphi_0} \right]$$~~

$$= \frac{2c\sqrt{2}}{\sqrt{\rho c v - k}} \arcsin \left[\frac{\sqrt{\rho c v - k}}{\sqrt{2 v_0^2 n_0^2 - k}} \right] + \dots$$

BJ

121



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122



0078

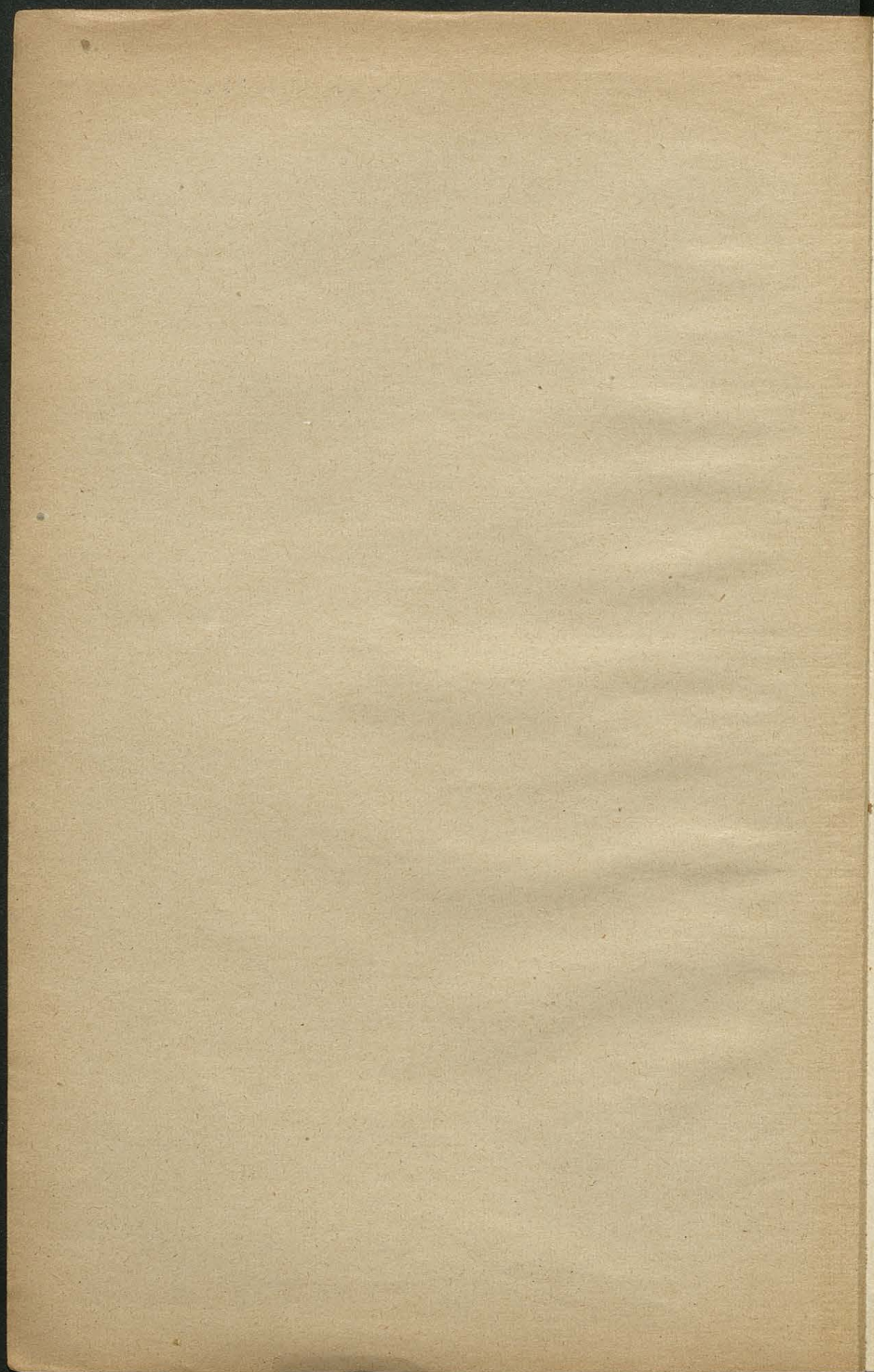
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$\int \frac{\partial u}{\partial t} = k \frac{\partial}{\partial r} \left[2\pi r \frac{\partial u}{\partial r} \right]$

$= k \frac{\partial}{\partial r} \left[2\pi r \frac{\partial u}{\partial r} \right] dx dr$

$r \frac{\partial u}{\partial t} = m \frac{\partial}{\partial r} \left(r \frac{\partial u}{\partial r} \right) = m \left[\frac{\partial u}{\partial r} + r \frac{\partial^2 u}{\partial r^2} \right]$

$u = r^\alpha t^\beta$

~~$\beta r^{\alpha+1} t^\beta = m \alpha r^{\alpha-1} t^\beta$~~

$u = r^\alpha e^{t^\beta}$

~~$\beta r^{\alpha+1} e^{t^\beta} = m \alpha r^{\alpha-1} e^{t^\beta}$~~

$u = e^{\alpha r} t^\beta$

~~$\beta r e^{\alpha r} t^\beta = m \left[\alpha e^{\alpha r} t^\beta + r \alpha^2 e^{\alpha r} t^\beta \right]$~~



$$\frac{\partial v}{\partial r^2} + \frac{\partial v}{\partial \theta^2} = 0 \quad ?$$

~~$$r^2 = c$$

$$v + \mu \frac{\partial v}{\partial r} = 0$$

$$\frac{\partial v}{\partial r} + \frac{\mu}{r^2} v = 0$$~~

$$r = \frac{\frac{c^2}{\mu}}{1 + \sqrt{1 + \frac{2c^2 C}{\mu^2}} \cos \varphi}$$

$$\mu = k(m + m_1)$$

$$c = \sqrt{\mu k(m + m_1)}$$

$$C = \frac{k(m + m_1)(e^2 - 1)}{\mu}$$

$$\frac{v^2}{2} = \frac{\mu}{r} + C$$

$$C = \frac{v^2}{2} \Big|_{r=\infty}$$

~~$$\sqrt{1 + \frac{2c^2 C}{\mu^2}} \cos \varphi = -1$$

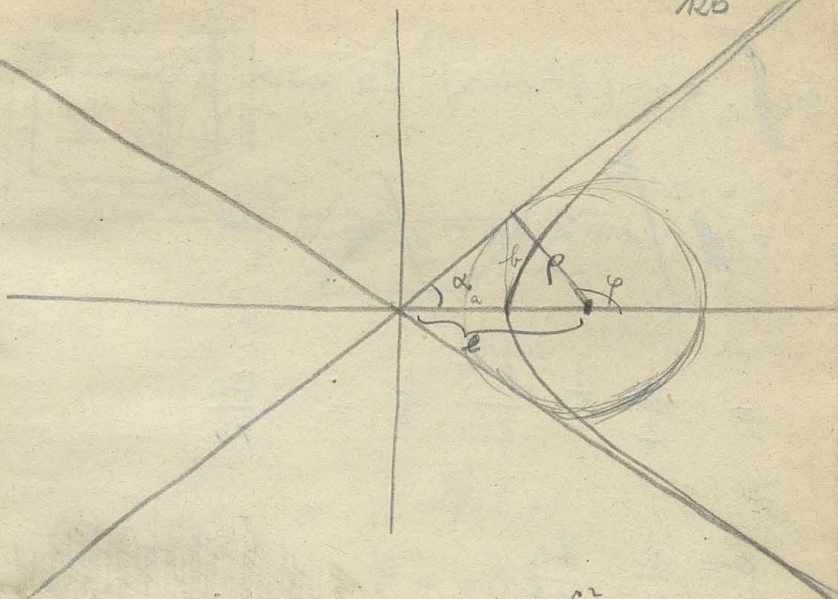
$$\mu + 2c^2 C \cos^2 \varphi = \mu^2$$~~

$$\cos^2 \varphi \sqrt{1 + \frac{2c^2 C}{\mu^2}} = -1$$

$$\cos^2 \alpha (\mu + 2c^2 C) = \mu^2$$

$$c = \sqrt{\frac{\mu^2(1 - \cos^2 \alpha)}{2C}}$$

$$= \frac{\mu \sin \alpha}{\sqrt{2C}} = \frac{\mu \sin \alpha}{v}$$



$$\rho = l \sin \alpha$$

$$\sqrt{\quad} = \frac{l}{a}$$

$$r = \frac{\frac{\mu \sin^2 \alpha}{v^2}}{1 + \cos \varphi \sqrt{1 + \sin^2 \alpha}}$$

$$\varphi = \pi$$

$$r_1 = \frac{\mu \sin^2 \alpha}{v^2} \frac{1}{1 + \mu \sqrt{\quad}}$$

$$\frac{r_1 - r_2}{2} = a$$

$$r = \frac{\mu \sin^2 \alpha}{v^2} \frac{1}{1 - \sqrt{\quad}}$$

$$= \frac{\mu \cancel{\quad}}{v^2} \frac{\sqrt{1 + \sin^2 \alpha}}{\cancel{\quad}}$$

$$l = a \sqrt{\quad} = \frac{\mu}{v^2} (1 + \sin^2 \alpha)$$

$$\rho = \frac{\mu \sin^2 \alpha}{v^2} (1 + \sin^2 \alpha)$$

$$\frac{1}{2\pi} \int_0^{\frac{\pi}{2}} \sin \alpha (1 + \sin^2 \alpha) 2\pi \sin \alpha d\alpha$$

$$= \int_0^{\frac{\pi}{2}} \sin^2 \alpha d\alpha + \int_0^{\frac{\pi}{2}} \sin^4 \alpha d\alpha$$

$$= \frac{\pi}{4} + \frac{1.3}{2.4} \frac{\pi}{2} = \frac{7\pi}{16}$$

$$p = \frac{\mu}{v^2} \frac{7\pi}{16}$$

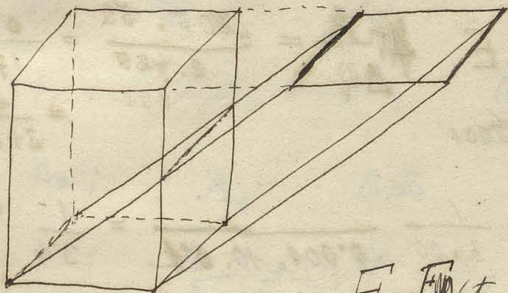
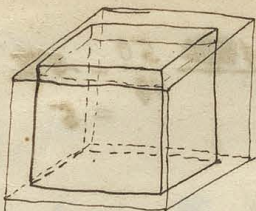


$$n = \frac{10^{-11} \cdot 6 \cdot 10^{-8}}{400^2} = 10^{-15} = 6 \cdot 10^{-23}$$

$$t_1, \mu_1, c_1 \quad \int_{r_1}^{r_2} \frac{\rho dr}{\rho E} = \frac{r_2^2 - r_1^2}{2\rho E} = \frac{(r_2 - r_1)(r_2 + r_1)}{2\rho E}$$

$$d = \frac{\mu}{\rho E}$$

$$\Delta d = \frac{\Delta \mu}{\rho E}$$



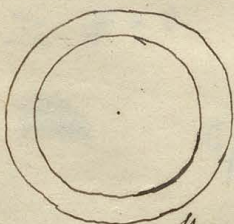
$$v, E = \Phi(t, r)$$

$$\theta = \varphi(r)$$

$$m = f(t, r)$$

$$l = \Phi(\theta)$$

$$\rho = \mu \left(1 + \frac{r}{E}\right)$$



$$\int_{r_1}^{r_2} \int_{\theta_1}^{\theta_2} 4\pi r^2 \left(1 + \frac{r}{E}\right)^2 \rho r \frac{dr}{E} - \int_{r_1}^{r_2} l \rho r dr +$$

$$+ \int_{t_1}^{t_2} \int_{r_1}^{r_2} c_{r_2} dt + \int_{t_1}^{t_2} \int_{r_1}^{r_2} c'_{r_2} dt$$

$$l = 1 \text{ m}$$

$$q = 1 \text{ mm}^2$$

$$P = 50 \text{ kg}$$

$$E = 10000$$

$$\lambda = \frac{lP}{E q} = \frac{1 \cdot 50}{10 \cdot 1} = 5$$

$$\frac{dE}{dt} = 0.0004 \cdot E$$

~~0.0001~~ 0.0001

$$\frac{\Delta E}{E} = \frac{0.0005 \cdot 50}{2.425} = \frac{0.05}{170}$$

$$= \frac{1}{3400} \text{ Cal}$$

$$\Delta t = \frac{\Delta Q}{q l c s} = \frac{1}{3400 \cdot 0.001 \cdot 10 \cdot 0.1} = \frac{1}{3.4} = 0.3^{\circ}$$

$$\frac{\Delta E}{E} = \frac{0.0004}{2} = 0.00012$$

$$\frac{\Delta \lambda}{\lambda} = 0.0006$$

$$\alpha = 0.00002$$

$$\frac{\Delta \lambda}{\lambda} = 0.02 \cdot 0.3 = 0.006 = 0.1 \%$$

$$l = 1 \text{ m}$$

$$q = 1 \text{ mm}^2$$

$$P = 1000 \text{ mm}$$

$$E = 0.1$$

$$P = 0.1 \text{ kg}$$

$$R = \frac{\lambda E q}{l}$$

$$A = \frac{\lambda P}{2}$$

$$\Delta Q = \frac{0.05}{425} = \frac{1}{8500} \text{ Cal}$$

Kreisprozess:

$E, \theta, p, A, \varphi, \alpha, c, s$

127

1). θ_0, p_0

2). θ_0, p_1 ; ~~A_2~~ $\Delta\varphi_2$

3). θ_1, p_1 ; $\Delta_E A_3 + \Delta_\alpha A_3$ $\varphi_3 = c_{(p)}(\theta_1 - \theta_0) \Delta$

4). θ_1, p_0 ; A_4 $\Delta\varphi_4$

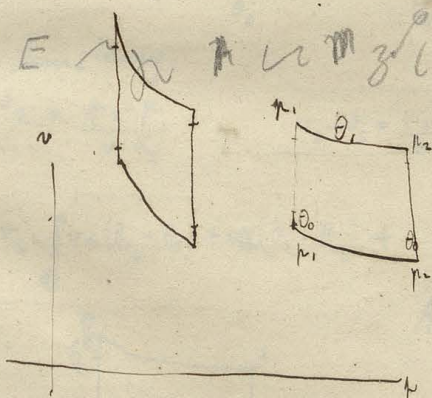
5). θ_0, p_0 ; ~~$\Delta_E A_5 + \Delta_\alpha A_5$~~ φ_5

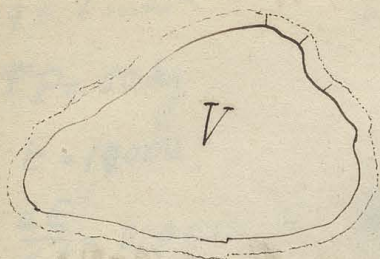
$$\sum \frac{\varphi}{T} = 0$$

~~$$c_f(\theta_1 - \theta_0) + h_{s_1} + c_e(\theta_1 - \theta_0)$$~~

cf

$\forall E$ ~~A~~ ~~φ~~ ~~α~~ ~~c~~ ~~s~~





$$t_1, \rho_1$$

$$t_1, \rho_2$$

$$t_2, \rho_2$$

$$t_2, \rho_1$$

$$1). t_1, \rho_1$$

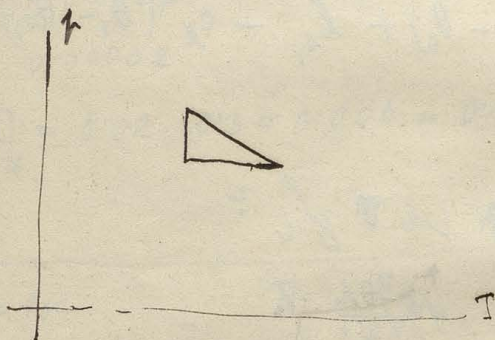
$$2). t_1, \rho_2$$

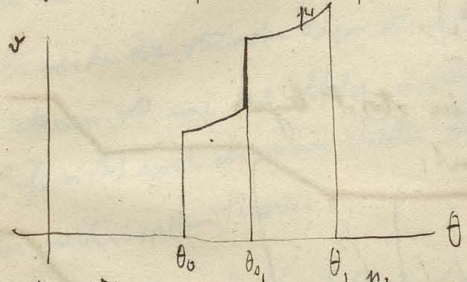
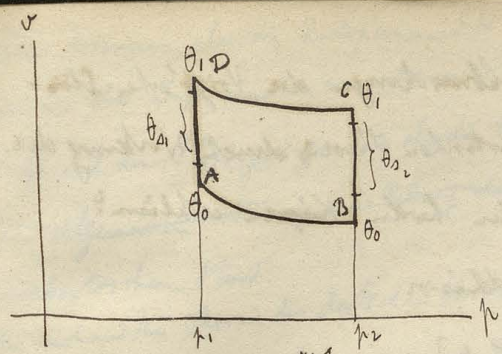
$$A = \frac{(\rho_2^2 - \rho_1^2)}{2KV}$$

$$Q =$$



$$\int \frac{dq}{T} = 0$$





$v_A - v_D = \int_{p_1}^{p_2} \frac{dp}{vK}$

$v_C - v_B = \int_{\theta_0}^{\theta_1} v \alpha d\theta + \int_{\theta_2}^{\theta_1} v \beta d\theta + \left(\frac{1}{s_{I\uparrow}} - \frac{1}{s_{I\downarrow}} \right) v$

Annahme:

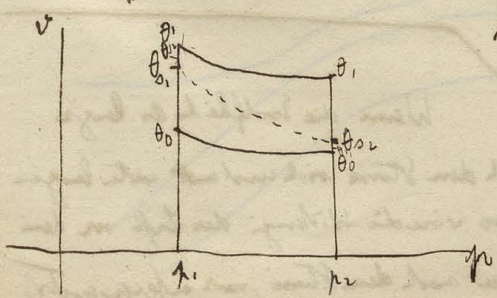
$v_A - v_B = \frac{p_2 - p_1}{v K_I}$

$v_D - v_C = \frac{p_2 - p_1}{v K_{II}}$

$\theta_{S1} > \theta_{S2}$

$v_D - v_A = \int_{\theta_0}^{\theta_{S1}} v \alpha d\theta + \int_{\theta_{S1}}^{\theta_1} v \alpha d\theta + \dots$

$A = \frac{(p_1 - p_2)(v_D - v_A + v_C - v_B)}{2}$

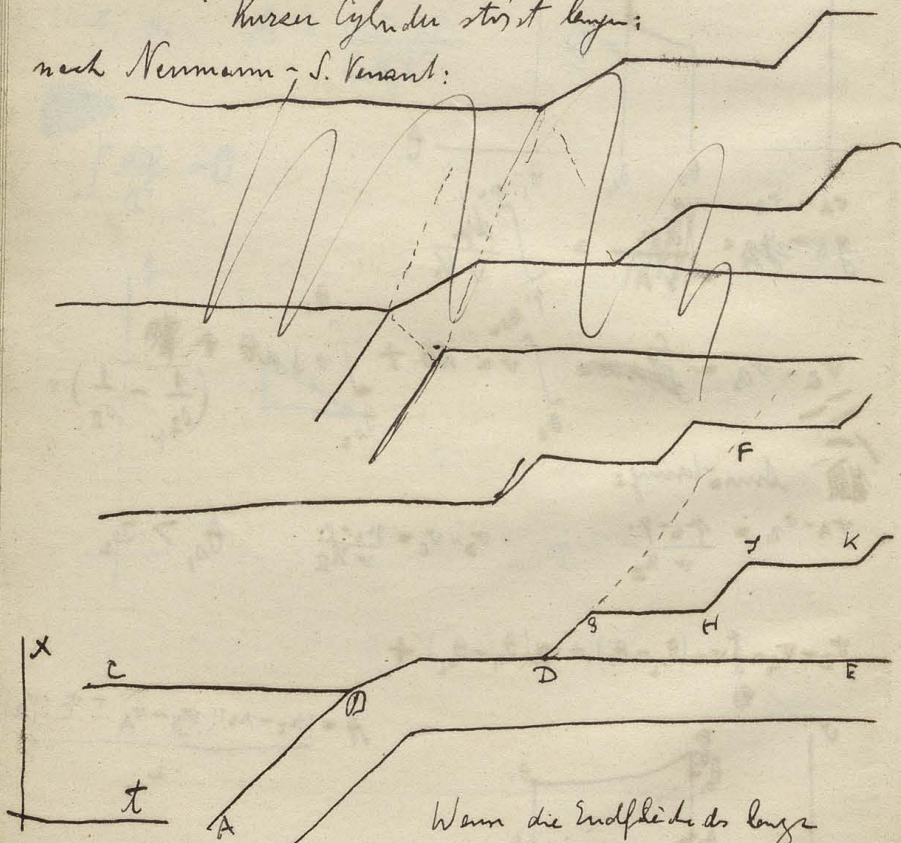


2. Ob sich nicht die Abweichungen der Voigt'schen Strom-Experimente von der elastischen Theorie durch Wirkung der Luftschichte zwischen den beiden Körpern erklären?

Stefan Schermbare Adhäsion

$$t = \frac{32 \mu R^3}{49} \left[\frac{1}{a^2} - \frac{1}{x^2} \right]$$

Kurven Cyklen der Stöße liegen:
nach Neumann - S. Venant:



Wenn die Endfläche des Körpers nach dem Stosse sich unstant weiter bewegen würde ~~an~~ [D S F] so würde die Wirkung der Luft von dem Stosse annähernd gleich genau nach dem Stosse, nur entgegengesetzt

Nun kommen aber die Ruhepausen S_H , J_K etc., vor, welche
 welche die Luft einströmen kann; daher wird die verurtheilte
 Kraft nach der Stärke geringer sein als die trennende vor der Stärke.
~~In Allgemeinen Wirkung einer Pflanz, daher Annäherung an 129~~

und elastischen Stoff.

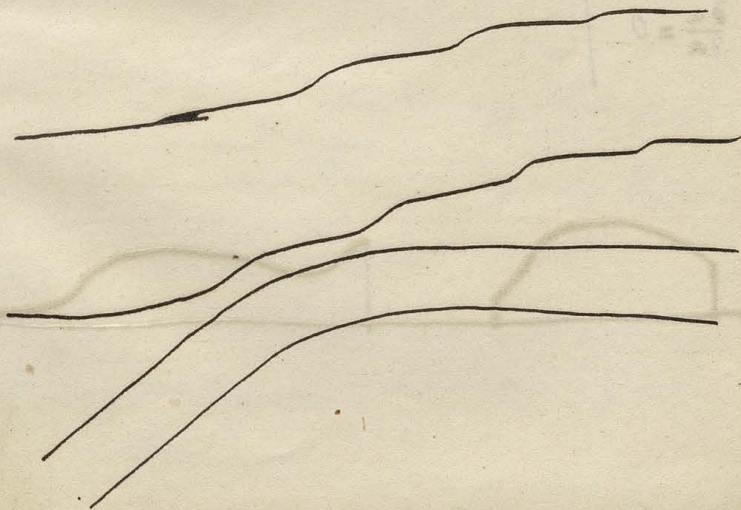
Die mechanische Theorie des elastischen Stoffes würde stimmen wenn
 W. wüsste die stehenden Körper ein sehr bewegliches elastisches Körper-
 Netzwerk W. eine Spiralfeder eingelegt wäre. Die sogenannte
 Luft wirkt auch als so ein Pulver daher Annäherung an letzter
 der mechanischen Theorie.

Die ~~Wirkung~~ Kraft welche in Momenten t , um das ~~Netz~~

Endflächen u v zusammen sind, und die ~~W~~

relativ höchste Geschwindigkeit $\frac{d\alpha}{dt}$ haben ist:

(bei Cylindern mit Radius R):
$$q = \frac{3\pi n}{2\alpha^3} R^4 \frac{d\alpha}{dt}$$



$$t = -\infty :$$

$$\frac{\partial u}{\partial t} = c_1$$

$$\xi < x < -(\xi + a_1)$$

$$\frac{\partial u}{\partial t} = c_2$$

$$x - \xi < x < x - \xi + a_2$$

$$\xi : x - \xi = c_1 t$$

$$x - \xi = \frac{c_2}{c_1} \xi$$

$$x = \frac{c_2 + c_1}{c_1} \xi$$

$$\frac{\partial u}{\partial t} = u f'(x + ut) - u \varphi'(x - ut)$$

$$\frac{c_1}{u} =$$

$$\left\{ \begin{array}{l} f'(x) = \frac{c_1}{2u} \\ f'(x) = \frac{c_2}{2u} \end{array} \right.$$

$$\varphi'(x) = -\frac{c_1}{2u}$$

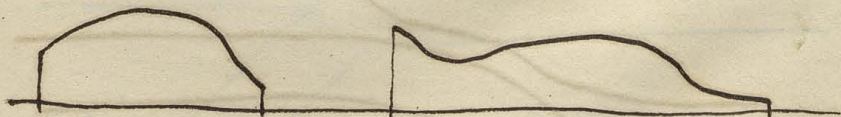
$$\varphi'(x) = -\frac{c_2}{2u}$$

$$\xi = \infty$$

$$-\xi < x < -(\xi + a_1)$$

$$\frac{c_2}{c_1} \xi < x < \frac{c_2}{c_1} \xi + a_2$$

$$\frac{\partial u}{\partial x} = 0$$



$$X_x = \left(2K \frac{1+3L}{1+2L} \right) \frac{\partial u_0}{\partial x} = A \frac{\partial u_0}{\partial x}$$

130

$$A \left(\frac{\partial u_0}{\partial x} \right)_1 = A \left(\frac{\partial u_0}{\partial x} \right)_2 = \underbrace{\frac{3\pi \mu R^4}{2} \frac{1}{\alpha^3} \frac{d\alpha}{dt}}_B \frac{d\alpha}{dt}$$

$$-\frac{8}{9} \alpha^2 - (\xi + \eta)$$

$$x_1 = -\xi$$

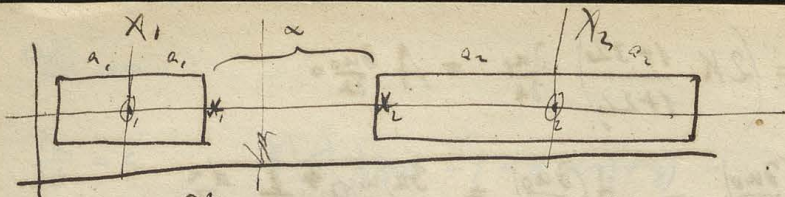
$$x_2 = \alpha - \xi - \eta$$

$$\frac{\partial u_0}{\partial x} = 0$$

$$x_1 = -(\xi + \eta)$$

$$x_2 = \alpha - \xi + \eta$$

u



$$m_1 \frac{d^2 x_1}{dt^2} = + \frac{B}{\alpha^3} \frac{dx}{dt} \quad m_1 \frac{dx_1}{dt} = - \frac{B}{2\alpha^2} + C_1, m_1$$

$$m_2 \frac{d^2 x_2}{dt^2} = - \frac{B}{\alpha^3} \frac{dx}{dt} \quad m_2 \frac{dx_2}{dt} = + \frac{B}{2\alpha^2} + C_2, m_2$$

~~$\alpha =$~~

$$\alpha = X_2 - X_1 - (a_1 + a_2) - (u_1 + u_2) \Big|_{x=l_1, u=l_2}$$

$$\frac{\partial^2 u}{\partial t^2} = \omega^2 \frac{\partial^2 u}{\partial x^2}$$

$$\left\{ \begin{array}{l} x = -l_1 \\ x = +l_1 \end{array} \right. \left. \begin{array}{l} \frac{\partial u}{\partial x} = 0 \\ A \frac{\partial u}{\partial x} = \frac{B}{\alpha^3} \frac{dx}{dt} = - \frac{B}{2} \frac{d(\frac{1}{\alpha^2})}{dt} \end{array} \right.$$

$$t = -T: \quad u = g(x) \quad x < +l_1$$

$$\frac{\partial u}{\partial t} = h(x)$$

$$u = f(x + \omega t) + \varphi(x - \omega t)$$

$$f(x + \omega T) + \varphi(x - \omega T) = g(x)$$

$$\omega f'(x + \omega T) - \omega \varphi'(x - \omega T) = h(x)$$

~~$f(x + \omega T) = \varphi(x - \omega T)$~~
 ~~$f'(x + \omega T) = \varphi'(x - \omega T)$~~
 ~~$f(x + \omega T) = \varphi(x - \omega T)$~~
 ~~$f'(x + \omega T) = \varphi'(x - \omega T)$~~

$$\frac{\partial u}{\partial t} = \omega f'(x + \omega t) - \omega \varphi'(x - \omega t)$$

$$\cancel{m_1 \frac{dX_1}{dt}} + m_2 \frac{dX_2}{dt} = \cancel{m_1 c_1 + m_2 c_2}$$

$$\frac{d\alpha}{dt} = \frac{dX_2}{dt} - \frac{dX_1}{dt} - \left(\frac{du_1}{dt} + \frac{du_2}{dt} \right)$$

$$= +\frac{\beta}{2\alpha^2 m_1} + \frac{\beta}{2\alpha^2 m_2} - c_1 + c_2 - \left(\frac{du_1}{dt} + \frac{du_2}{dt} \right)$$

$$= -\frac{\beta}{2\alpha^2} \frac{1}{m_1}$$

$$= + \frac{\beta (m_2 + m_1)}{2 m_1 m_2} \frac{1}{\alpha^2} \dots$$

$$\cancel{\frac{d(\alpha + u_1 + u_2)}{dt} \neq 0} \quad \frac{d(\alpha + u_1 + u_2)}{dt} = + \frac{\beta (m_2 + m_1)}{2 m_1 m_2} \frac{1}{\alpha^2} + c_2 - c_1$$

~~A~~

$$f'(x + \omega T) = \frac{1}{2} (L(x) + p'(x))$$

$$\varphi'(x - \omega T) = \frac{1}{2} (-L(x) + p'(x))$$

$$f(x + \omega T) = \frac{1}{2} [g(x) + \int L(x) dx] \quad -d_1 < x < +d_2$$

$$\varphi(x - \omega T) = \frac{1}{2} [p(x) - \int L(x) dx]$$

$$f(-d_1 + \omega t) + \varphi'(-d_1 - \omega t) = 0$$

$$A_1 \left\{ f(l_1 + \omega t) + \varphi'(l_1 - \omega t) \right\} = -\frac{\beta}{2} \frac{d\left(\frac{1}{\alpha^2}\right)}{dt}$$

$$u_1 = a_1 + \frac{b_1}{t} + \frac{c_1}{t^2} + \frac{d_1}{t^3} + \frac{e_1}{t^4} + \frac{f_1}{t^5} + \dots$$

$$u_2 = a_2 + \frac{b_2}{t} + \frac{c_2}{t^2} + \frac{d_2}{t^3} + \dots$$

$$a, b, \dots = f(x)$$

$$\frac{\partial u_1}{\partial t} = - \left[\frac{b_1}{t^2} + \frac{2b_2}{t^3} + \frac{3c_3}{t^4} + \dots \right]$$

$$\frac{\partial^2 u_1}{\partial t^2} = + \frac{2b_1}{t^3} + 2 \cdot 3 \frac{c_2}{t^4} + 3 \cdot 4 \frac{d_3}{t^5} + \dots$$

$$\cancel{2 \frac{b_1}{t^3} + 2 \cdot 3 \frac{c_2}{t^4} + 3 \cdot 4 \frac{d_3}{t^5} + \dots} = \frac{\partial^2 u_1}{\partial t^2}$$

$$2 \frac{b_1}{t^3} + 2 \cdot 3 \frac{c_2}{t^4} + 3 \cdot 4 \frac{d_3}{t^5} + \dots =$$

$$\frac{\partial^2 a_1}{\partial x^2} + \frac{1}{t} \frac{\partial^2 b_1}{\partial x^2} + \frac{1}{t^2} \frac{\partial^2 c_1}{\partial x^2} + \dots$$

$$\frac{\partial^2 a_1}{\partial x^2} = 0 \quad 2b_1 = \frac{\partial^2 d_1}{\partial x^2}$$

$$\frac{\partial^2 b_1}{\partial x^2} = 0 \quad 2 \cdot 3 \cdot c_1 = \frac{\partial^2 c_1}{\partial x^2}$$

$$\frac{\partial^2 c_1}{\partial x^2} = 0 \quad 3 \cdot 4 \cdot d_1 = \frac{\partial^2 f_1}{\partial x^2}$$

$$a_1 = M_1 x + N_1$$

$$b_1 = M_2 x + N_2$$

$$c_1 = M_3 x + N_3$$

$$d_1 = M_2 \frac{x^3}{3} + N_2 x^2 + O_2 x + P_2$$

$$e_1 = M_3 \frac{x^4}{4} + 3N_3 x^2 + O_3 x + P_3$$

$$2M_2x + 2N_2 = \frac{d^2d_1}{dx^2}$$

$$M_2x^2 + 2N_2x = \frac{d^4d_1}{dx^4} + O_2$$

$$M_2 \frac{x^3}{3} + N_2x^2 + O_2x + P_2 = d_1$$

$$\frac{d^2f_1}{dx^2} = 3.4 \cdot [P_2 + O_2x + N_2x^2 + M_2 \frac{x^3}{3}]$$

$$\frac{df_1}{dx} = 3.4 \cdot [P_2x + O_2 \frac{x^2}{2} + N_2 \frac{x^3}{3}] + \frac{M_2x^4}{4} + Q_4$$

$$f_1 = 3.4 [P_2 \frac{x^2}{2} + O_2 \frac{x^3}{2.3} + N_2 \frac{x^4}{3.4}] + M_2 \frac{x^5}{5} + Q_4x + R_4$$

~~6.4~~

$$f_1 = M_2 \frac{x^5}{5} + N_2x^4 + 2 \cdot O_2x^3 + 6P_2x^2 + Q_4x + R_4$$

$$\left\{ \begin{aligned} -\frac{B}{2} \frac{m_1+m_2}{m_1m_2} \frac{1}{x^2} &= \frac{d(\alpha+u_1+u_2)}{dt} + c_1 - c_2 \\ A_1 \frac{\partial u_1}{\partial x} &= -\frac{B}{2} \frac{d}{dt} \left(\frac{1}{x^2} \right) \end{aligned} \right\} \parallel x=a_1$$

$$A_1 \frac{m_1+m_2}{m_1m_2} \frac{\partial u_1}{\partial x} = \frac{d^2}{dt^2} (\alpha+u_1+u_2) \parallel x=a_1$$

~~$$= \frac{d^2}{dt^2} [x_2 - x_1 - (\alpha + u_1 + u_2)]$$~~

$$2A_1 \left\{ \frac{1}{t} \frac{\partial c_1}{\partial x} + \frac{1}{t^2} \frac{\partial c_1}{\partial x} + \frac{1}{t^3} \frac{\partial c_1}{\partial x} + \frac{1}{t^4} \frac{\partial c_1}{\partial x} + \frac{1}{t^5} \frac{\partial c_1}{\partial x} + \dots \right\}$$

$$= -B \frac{d(\frac{1}{x^2})}{dt}$$

$$\frac{1}{2} \frac{d^2 x}{dt^2} = -\frac{g}{2} x$$

$$M \frac{d^2 x}{dt^2} + kx = 0$$

$$\frac{d^2 x}{dt^2} + \frac{k}{M} x = 0$$

$$\frac{d^2 x}{dt^2} + \omega^2 x = 0$$

$$x = A \cos(\omega t + \phi) + B \sin(\omega t + \phi)$$

$$x = A \cos(\omega t + \phi)$$

$$\left. \begin{aligned} \frac{dx}{dt} &= -A\omega \sin(\omega t + \phi) \\ \frac{d^2 x}{dt^2} &= -A\omega^2 \cos(\omega t + \phi) \end{aligned} \right\}$$

$$A \cos(\omega t + \phi) = x$$

$$\frac{dx}{dt} = -A\omega \sin(\omega t + \phi)$$

$$x = A \cos(\omega t + \phi)$$

$$\frac{dx}{dt} = -A\omega \sin(\omega t + \phi)$$

damit einer der beiden Platte eini des prupt, müsste ein 133

Moment sein, wo $\frac{dX_1}{dt}$ oder $\frac{dX_2}{dt} = 0$ also:

$$-\frac{\beta}{2\alpha^2} + c_1 m_1 = 0 \quad \text{oder} \quad +\frac{\beta}{2\alpha^2} + c_2 m_2 = 0$$

~~verlangt ein $c_2 < 0$~~

verlangt ein $c_2 < 0$

$$\frac{\beta}{2\alpha^2} = m_1 c_1$$

dann ist:

$$m_2 \frac{dX_2}{dt} = m_2 c_2 + m_1 c_1$$

$$\alpha = \sqrt{\frac{\beta}{2m_1 c_1}}$$

~~dx~~
A

Wenn die Länge des stromdurchlässigen Zylinders im Verhältnis zu seinem Querschnitt sehr kurz ist, so dass diese als Scheiben betrachtet werden können, wird die elektrische Verschiebung sehr klein werden; wenn man annimmt:

$$\frac{dx}{dt} = \frac{dx}{dt}; \text{ so wird sein:}$$

$$m, \frac{dx}{dt} = + \frac{B}{2x^2} - m, c,$$

$$\frac{dx}{dt} - \frac{B}{2m, x^2} + m, c = 0$$

$$\frac{dx}{dt} = \frac{B}{2m, x^2} - m, c$$

$$\int \frac{x^2 dx}{\frac{B}{2m,} - m, c, x^2} = \int dt$$

$$\int \frac{\frac{2m,}{B} x^2 dx}{1 - \frac{2m, c,}{B} x^2} = t + C$$

$$x \sqrt{\frac{2m, c,}{B}} = y$$
$$x = \sqrt{\frac{B}{2m, c,}} y$$

$$= \frac{2m, B}{B 2m, c,} \sqrt{\frac{B}{2m, c,}}$$

$$\int \frac{x^2 dx}{1-x^2} = \int \frac{1}{2} \left(\frac{-x}{1+x} + \frac{x}{1-x} \right) dx$$

$$\frac{1}{1+x} = y$$

$$J_1 = \int \frac{y-1}{y} dy = y - \ln y$$
$$= (1+x) - \ln(1+x)$$

$$1-x = y$$

$$J_2 = \int \frac{1-y}{y} dy = -\ln y + y$$
$$= (1-x) - \ln(1-x)$$

$$J = \frac{1}{2} \left(-2x + \ln \frac{1+x}{1-x} \right)$$
$$= -x + \ln \sqrt{\frac{1+x}{1-x}}$$

$$t+C = \sqrt{\frac{B}{2m, c_1^3}} \left(\cancel{\dots} - x \sqrt{\frac{2m, c_1}{B}} + \log \left| \frac{1+x \sqrt{\frac{2m, c_1}{B}}}{1-x \sqrt{\frac{2m, c_1}{B}}} \right| \right)$$

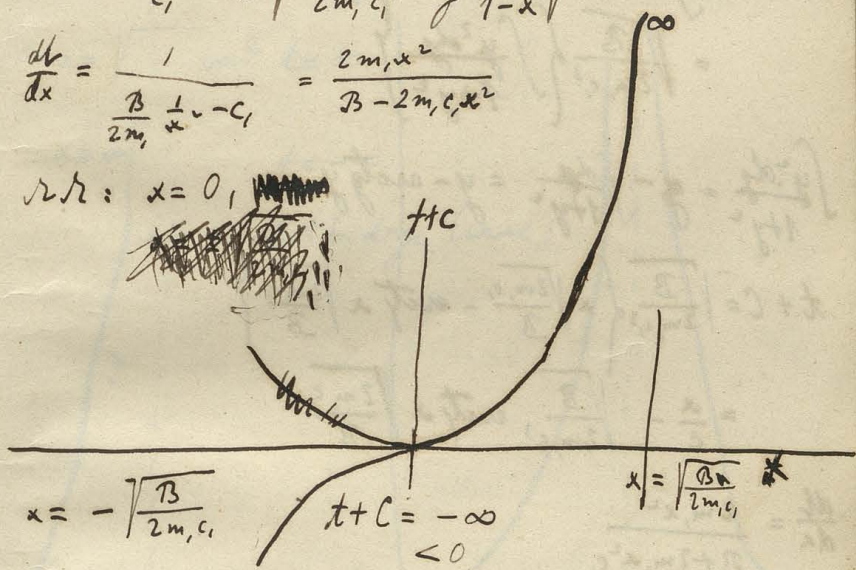
134

$$= -\frac{x}{c_1} + \frac{1}{2} \sqrt{\frac{B}{2m, c_1^3}} \log \left| \frac{1+x \sqrt{\frac{2m, c_1}{B}}}{1-x \sqrt{\frac{2m, c_1}{B}}} \right|$$

$$= -\frac{x}{c_1} + \frac{1}{2} \sqrt{\frac{B}{2m, c_1^3}} \log \left| \frac{1+x \sqrt{\frac{2m, c_1}{B}}}{1-x \sqrt{\frac{2m, c_1}{B}}} \right|$$

$$\frac{dt}{dx} = \frac{1}{\frac{B}{2m, x} - c_1} = \frac{2m, x^2}{B - 2m, c_1 x^2}$$

z.z.: $x=0$, ~~...~~



$$x = -\sqrt{\frac{B}{2m, c_1}}$$

$$t+C = -\infty < 0$$

$$x = \sqrt{\frac{B}{2m, c_1}}$$

$$x = \dots$$

~~...~~

was wenn $x > \dots$?!

$$x = 0$$

$$t+C = 0$$

$$x = + \dots$$

dies wenn Geschwindigkeit der

~~Schiff~~ ~~wenn~~ ~~Reibung~~ ~~der~~ ~~Wand~~

Scherbe gegen die feste Wand

~~von~~ ~~der~~ ~~Wand~~ ~~weg~~
 also asymptotische Annäherung an den Punkt $x = \sqrt{\dots}$
 für $t \rightarrow \infty$

$$m_1 \frac{dx}{dt} = \frac{B}{2x^2} + m_1 c_1$$

$$x \sqrt{\frac{2m_1 c_1}{B}} = y$$

$$\frac{dx}{\frac{B}{2m_1 x^2} + c_1} = dt$$

$$t + C = \int \frac{1}{\frac{B}{2m_1 x^2} + c_1} dx$$

$$= \sqrt{\frac{B}{2m_1 c_1^3}} \left\{ \int \frac{y^2 dy}{1+y^2} \right\}$$

$$\int \frac{y^2 dy}{1+y^2} = y - \int \frac{dy}{1+y^2} = y - \arctan y$$

$$t + C = \sqrt{\frac{B}{2m_1 c_1^3}} \left\{ x \sqrt{\frac{2m_1 c_1}{B}} - \arctan x \sqrt{\frac{2m_1 c_1}{B}} \right\}$$

$$= \frac{x}{c} - \sqrt{\frac{B}{2m_1 c_1^3}} \arctan x \sqrt{\frac{2m_1 c_1}{B}}$$

$$\frac{dt}{dx} = \frac{2m_1 x^2}{B + 2m_1 x^2 c_1}$$

$$M.H.: x=0 \quad t+C=0$$

ist frühere Annahme: c_1

eine Lösung ist auch:

$$t+C = -\frac{x}{c_1} + \frac{1}{2} \sqrt{\frac{B}{2m_1 c_1^3}} \operatorname{Log} \left\{ \frac{x \sqrt{\frac{2m_1 c_1}{B}} + 1}{x \sqrt{-1}} \right\}$$

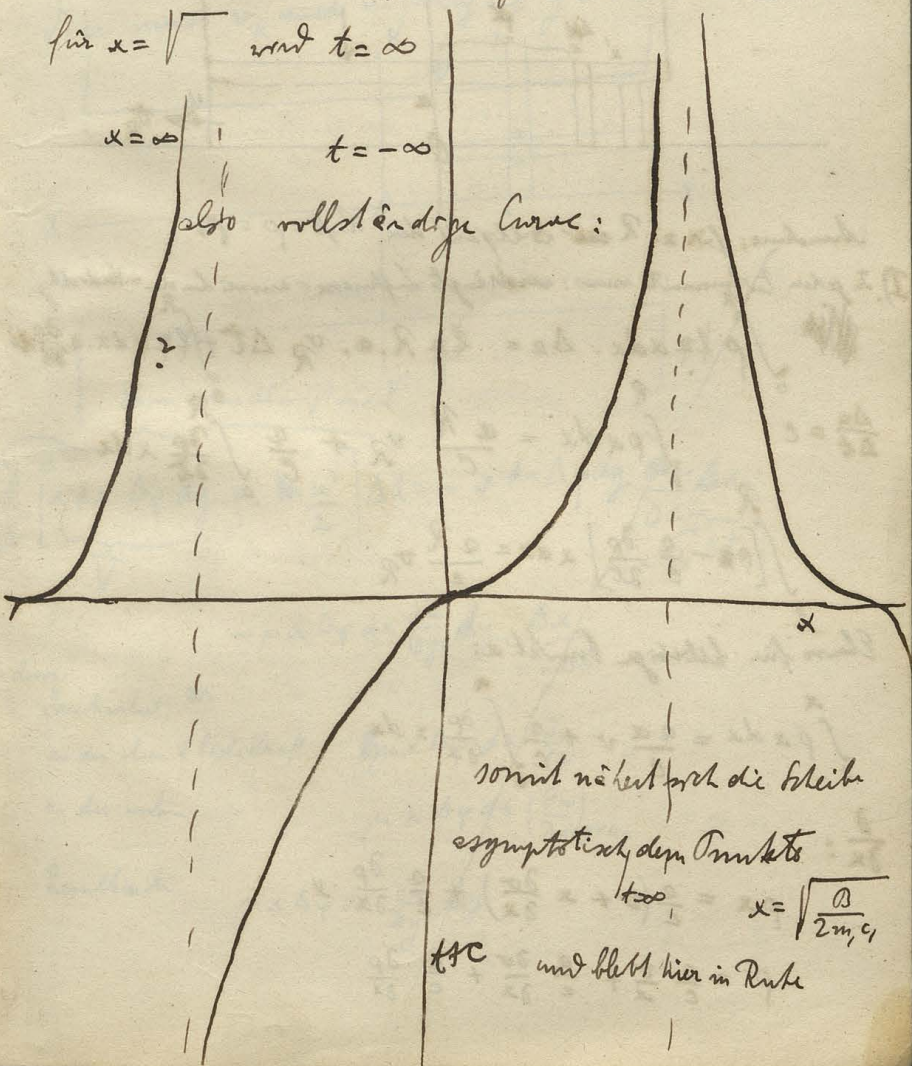
dies gilt für $x > \sqrt{\frac{B}{2m_1 c_1}}$ also gewöhnlicher Fall

für $x = \sqrt{\frac{B}{2m_1 c_1}}$ und $t = \infty$

$x = \infty$

$t = -\infty$

also vollständige Curve:



somit nähert sich die Gleichung asymptotisch dem Punkte

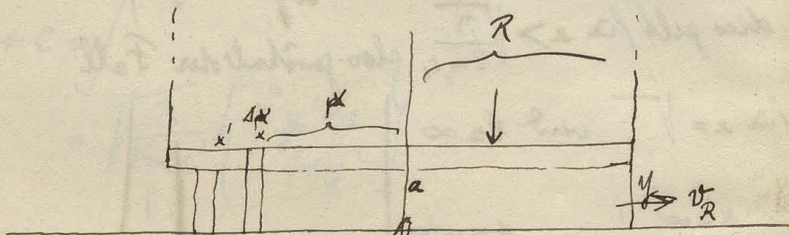
$$t \rightarrow \infty \quad x = \sqrt{\frac{B}{2m_1 c_1}}$$

und bleibt hier in Ruhe

Vorübung:

(Körperliche)

Eine Platte bewegt sich mit der constanten Geschwindigkeit c senkrecht gegen eine Wand μ (in Luft), welches sind die durch Compression der Luft entstehenden Kräfte?



Annahme: für $x=R$ sei Dichtigkeit der Luft $\rho = \rho_0$

D. In jedem Zeitmomente muss: verdrängte Luftmasse = anströmende + Verdichtung

$$\int_0^R \rho x dx \cdot \Delta a = \int_0^R \rho x dx \cdot a \cdot v_R \Delta t + \int_0^R \rho x dx \cdot a \frac{\partial \rho}{\partial t} \Delta t$$

$$\frac{\Delta a}{\Delta t} = c \quad \int_0^R \rho x dx = \frac{a R}{c} \cdot v_R + \frac{a}{c} \int_0^R \frac{\partial \rho}{\partial t} x dx$$

$$\int_0^R \left[\rho x - \frac{a}{c} \frac{\partial \rho}{\partial t} x \right] x dx = \frac{a R}{c} v_R$$

Ebenso für beliebigen Punkt x :

$$\int_0^x \rho x dx = \frac{a x}{c} v + \frac{a}{c} \int_0^x \frac{\partial \rho}{\partial t} x dx$$

$\frac{\partial}{\partial x}$:

$$\rho x = \frac{a}{c} \left(v + x \frac{\partial v}{\partial x} \right) + \frac{a}{c} \frac{\partial \rho}{\partial t} \cdot x$$

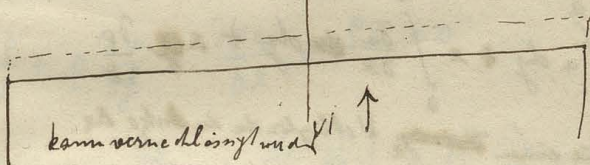
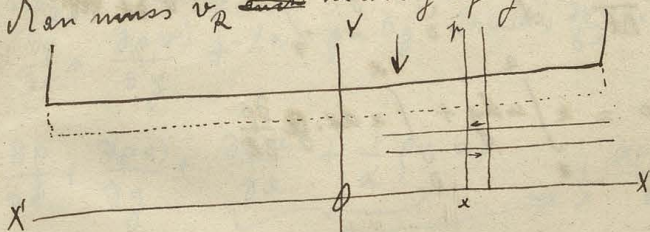
$$\rho = \frac{a}{c} \frac{v}{x} + \frac{a}{c} \frac{\partial v}{\partial x} + \frac{a}{c} \frac{\partial \rho}{\partial t}$$

$$\int \rho dx = \frac{a}{c} \int \frac{v}{x} dx + \frac{e}{c} (v + p) + \text{const}$$

II). Kräfte die auf ein Volumenelement wirken
Arbeit bei Verschiebung des Elementes von x nach x'

$$\Delta \varphi \cdot dx \cdot dy \cdot dz$$

Man muss v_R ~~und~~ ^{und überlegt v_x} in Bezug auf y als variabel ansehen



kann vernachlässigt werden v_i ↑

$$\frac{d}{dt} \left[\underbrace{x dx \Delta \varphi dy \cdot \rho \cdot \frac{u^2}{2}}_{V} \right] \Delta t = - x dx \Delta \varphi dy \frac{d\varphi}{dx} \Delta x$$

$[= u \Delta t]$

$$- \rho x \Delta \varphi dx \frac{\partial u}{\partial y} dy \cdot \frac{\Delta x}{u \Delta t}$$

dann:

Druckarbeit: ρx

an der oben Flächekraft $\rho x \Delta \varphi dx \frac{\partial u}{\partial y}$

an der unteren

$$\rho x \Delta \varphi dx \left(\frac{\partial u}{\partial y} \right)_{y+dy}$$

Resultante

$$\rho x \Delta \varphi dx \frac{\partial^2 u}{\partial y^2} dy$$

$$\frac{\partial p}{\partial x} = -\mu \frac{\partial^2 u}{\partial y^2} \quad \leftarrow \frac{\partial}{\partial y}$$

$$\frac{\partial p}{\partial y} = 0 \quad \left| \quad \frac{\partial}{\partial y} \left(\frac{\partial p}{\partial x} \right) = \frac{\partial}{\partial x} \left(\frac{\partial p}{\partial y} \right) = 0 = -\mu \frac{\partial^3 u}{\partial y^3}$$

II Vordränge Masse

$$\int_0^x 2\alpha x dx \cdot \rho_0 \frac{\Delta a}{\Delta t} = \int_0^a 2\alpha x \int_0^x u dy + \int_0^x 2\alpha x dx \cdot \rho_0 \frac{\partial p}{\partial t}$$

$$c \int_0^x x dx \cdot \rho = x \int_0^a u dy + \int_0^x x dx \cdot \rho \frac{\partial p}{\partial t}$$

$$\frac{\partial}{\partial x} : \quad c x \cdot \rho_0 = \int_0^a u dy + x \int_0^a \frac{\partial u}{\partial x} dy + x \rho \frac{\partial p}{\partial t}$$

auch direkt für einen ~~Winkel~~ Δx Willkürlich da die Dicke Δx :

$$c \cdot \rho \cdot 2\alpha x \Delta x \cdot \frac{\Delta a}{\Delta t} = 2\alpha x \int_0^a u dy - 2\alpha(x+\Delta x) \int_0^{a+\Delta a} u dy + 2\alpha x \Delta x \cdot \rho \frac{\partial p}{\partial t}$$

$$c \rho 2\alpha x \Delta x = -2\alpha \Delta x \int_0^a u dy - 2\alpha x \int_0^a \frac{\partial u}{\partial x} dy \Delta x + 2\alpha x \Delta x c \rho$$

wie oben:

$$c x \rho_0 = \int_0^a u dy + x \int_0^a \frac{\partial u}{\partial x} dy + x \rho \frac{\partial p}{\partial t}$$

III $\frac{\mu}{\rho_0} = \frac{\rho}{\rho_0} \quad \rho = A \mu$

$$\frac{\partial \Pi}{\partial y} = \frac{\partial \Pi}{\partial y} =$$

$$\frac{\partial \Pi}{\partial x} = c\rho + c x \frac{\partial \rho}{\partial x} = 2 \int_0^a \frac{\partial u}{\partial x} dy + x \int_0^a \frac{\partial^2 u}{\partial x^2} dy + a \frac{\partial \rho}{\partial t} + ax \frac{\partial(\rho u)}{\partial t \partial x}$$

Annahme für einen Kontrollvolumen von der Dicke Δy in der Höhe y :

$$2n x \Delta x \Delta y \frac{\partial \rho}{\partial t} + 2n x \Delta x \Delta y \frac{\partial(\rho v)}{\partial y} - 2n x \rho u \Delta y - [n(x+\Delta x) \frac{\partial(\rho u)}{\partial x} \Delta x] \Delta y = 0$$

$$\frac{\partial \rho}{\partial t} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho u)}{\partial x} + \frac{1}{x} \rho u = 0 \quad \rightarrow \quad \frac{1}{x} \frac{\partial(\rho u x)}{\partial x}$$

$$\frac{\partial \rho}{\partial t} + \rho \frac{\partial v}{\partial y} + \frac{\partial(\rho u)}{\partial x} + \frac{\rho u}{x} = 0$$

$$\frac{\partial^2}{\partial y^2} : \quad \rho \frac{\partial^3 v}{\partial y^3} + \frac{\partial(\rho \frac{\partial^2 u}{\partial y^2})}{\partial x} + \frac{\rho}{x} \frac{\partial^2 u}{\partial y^2} = 0$$

$$\frac{1}{x} \frac{\partial[\rho x \frac{\partial^2 u}{\partial y^2}]}{\partial x}$$

$$\rho \frac{\partial^3 v}{\partial y^3} = \frac{1}{x \mu} \frac{\partial[\rho x \frac{\partial \tau}{\partial x}]}{\partial x} \quad \# \quad = \frac{A}{2x \mu} \frac{\partial[x \frac{\partial(\mu \dot{\gamma})}{\partial x}]}{\partial x}$$

$$\mu \rho \frac{\partial^3 v}{\partial y^3} = \frac{\partial(\rho \frac{\partial \tau}{\partial x})}{\partial x} + \frac{1}{x} \rho \frac{\partial \tau}{\partial x}$$

$$= \frac{\partial \rho}{\partial x} \frac{\partial \tau}{\partial x} + \rho \frac{\partial^2 \tau}{\partial x^2} + \frac{\rho}{x} \frac{\partial \tau}{\partial x}$$

$$= A \left(\frac{\partial \tau}{\partial x}\right)^2 + A \mu \frac{\partial^2 \dot{\gamma}}{\partial x^2} + A \frac{\mu}{x} \frac{\partial \dot{\gamma}}{\partial x}$$

$$\frac{\partial^4 v}{\partial y^4} = 0 \quad v = c_0 + c_1 y + c_2 y^2 + c_3 y^3 + \text{~~...}~~$$

$$c = f(x, t)$$

$$\frac{\partial^3 v}{\partial y^3} = \text{const} = c_3$$

$$\frac{\partial^3 v}{\partial y^3} \rho = -\frac{1}{x} \cdot \frac{\partial(\rho x \frac{\partial^2 u}{\partial y^2})}{\partial x}$$

$$\frac{\partial^3 v}{\partial y^3} \rho x = -\frac{\partial}{\partial x} \left[\rho x \frac{\partial^2 u}{\partial y^2} \right] = \frac{1}{\mu} \frac{\partial}{\partial x} \left[\rho x \frac{\partial \mu}{\partial x} \right]$$

~~... $\mu \frac{\partial \rho x}{\partial x}$...~~

$$\mu \frac{\partial \rho x}{\partial x} = \frac{\partial}{\partial x} \left[\rho x \frac{\partial \mu}{\partial x} \right] = \frac{\partial(\rho x)}{\partial x} \frac{\partial \mu}{\partial x} + \rho x \frac{\partial^2 \mu}{\partial x^2}$$

$$\frac{\partial^2 \mu}{\partial x^2} + \frac{1}{\rho x} \frac{\partial(\rho x)}{\partial x} \frac{\partial \mu}{\partial x} = \mu \frac{\partial^3 v}{\partial y^3}$$

$$\frac{\partial^2 \mu}{\partial x^2} + \frac{\partial(\log(\rho x))}{\partial x} \frac{\partial \mu}{\partial x} = \mu \frac{\partial^3 v}{\partial y^3}$$

$$\frac{\partial^2 \mu}{\partial x^2} \partial x + \frac{\partial \log(\rho x)}{\partial x} \partial x = \frac{\mu \frac{\partial^3 v}{\partial y^3}}{\frac{\partial \mu}{\partial x}} \partial x$$

$$\frac{\partial(\log(\rho x))}{\partial x}$$

$$\frac{\partial \log(\rho x \frac{\partial \mu}{\partial x})}{\partial x} \cdot \frac{\partial \mu}{\partial x} = \mu \frac{\partial^3 v}{\partial y^3}$$

$$u = a_0 + a_1 y + a_2 y^2 + \dots$$

$$a = f(x, t)$$

138

$$p = f(x, t)$$

$$\frac{\partial p}{\partial x} = -2\mu a_2 \quad \text{fehlt!} \quad p = b_0 - 2\mu \int a_2 dx$$

$$p = b_0 - 2a_2 \mu x + \dots \quad b = f_c(t)$$

II auch in der Form:

$$\frac{\partial p}{\partial t} + p \frac{\partial v}{\partial y} + \frac{1}{x} \frac{\partial (p u x)}{\partial x} = 0 \quad p_0 = (b_0) - 2(a_2) \mu R + \dots$$

$$\frac{\partial p}{\partial x} = -2\mu \frac{\partial a_2}{\partial x} = -\mu \quad \text{für } x=R$$

$$p = p_0 + 2(a_2) \mu (R-x)$$

III in der Form:

$$\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial y} + \frac{1}{x} \frac{\partial (u x)}{\partial x} = 0$$

$u = 0$ für alle t und $x=0$

$$y = a = E - ct$$

$$0 = a_0 + a_1 (E - ct) + a_2 (E - ct)^2$$

$$= a_0 + a_1 E + a_2 E^2 - (a_1 + 2a_2) ct$$

$$- [a_1 + 2a_2 E] ct + a_2 c^2 t^2$$

für $y=0$ wird u ein rel. Maximum haben also

$$\frac{\partial u}{\partial y} \Big|_{y=0} = 0$$

und für $y=0$ $x=R$ ein absolutes; dann muss

$$\text{also auch } \frac{\partial u}{\partial x} \Big|_{x=R, y=0} = 0 \text{ sein also: } \frac{\partial a_0}{\partial x} \Big|_{x=R} = 0$$

aus Kugelsystem:

$$A \exp = \int_0^a u dy + x \int_0^a \frac{\partial u}{\partial x} dy + A \exp \frac{\partial \mu}{\partial t}$$

$$A \exp [\rho_0 + 2a_1 \mu (R-x)] = a_0 a + a_1 \frac{a^2}{2} + a_2 \frac{a^3}{3}$$

$$A \exp [b_0 + 2a_1 \mu \int_0^a dx] + x \left(\frac{\partial a_0}{\partial x} a + \frac{\partial a_1}{\partial x} \frac{a^2}{2} + \frac{\partial a_2}{\partial x} \frac{a^3}{3} \right)$$

$$+ A \exp \frac{\partial \mu}{\partial t} (R-x) \frac{\partial a_2}{\partial t}$$

$$+ A \exp \left[\frac{\partial b_0}{\partial t} - 2a_1 \int_0^a \frac{\partial a_2}{\partial t} dx \right]$$

$$\frac{\partial \mu}{\partial t} + \mu \frac{\partial v}{\partial y} + \frac{1}{x} \frac{\partial (\mu u x)}{\partial x} = 0 \quad (10)$$

$$\frac{\partial}{\partial y} \uparrow =$$

$$\frac{\partial}{\partial x} \left(\frac{\partial \mu}{\partial y} \right) + \frac{\partial \mu}{\partial y} \frac{\partial v}{\partial y} + \mu \frac{\partial^2 v}{\partial y^2} + \frac{1}{x} \frac{\partial}{\partial x} \left\{ x \frac{\partial (\mu u)}{\partial y} \right\} = 0$$

$$\frac{\partial^2 (\mu u)}{\partial x \partial y} + \frac{1}{x} \frac{\partial (\mu u)}{\partial y} = 0$$

$$\frac{\partial}{\partial x} \left\{ u \frac{\partial \mu}{\partial y} + \mu \frac{\partial u}{\partial y} \right\} + \frac{1}{x} \left\{ u \frac{\partial \mu}{\partial y} + \mu \frac{\partial u}{\partial y} \right\} = 0$$

$$\mu \frac{\partial^2 v}{\partial y^2} + \frac{\mu}{x} \frac{\partial u}{\partial y} + \frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial y} \right) = 0$$

$$\frac{\partial \mu}{\partial x} \frac{\partial u}{\partial y} + \mu \frac{\partial^2 u}{\partial x \partial y}$$

$$\frac{1}{x} \frac{\partial}{\partial x} \left(x \mu \frac{\partial u}{\partial y} \right)$$

$$\mu \frac{\partial^2 v}{\partial y^2} + \frac{1}{x} \frac{\partial}{\partial x} \left\{ \mu x \frac{\partial u}{\partial y} \right\} = 0$$

139

$$\mu x \frac{\partial^2 v}{\partial y^2} + \frac{\partial}{\partial x} \left\{ \mu x \frac{\partial u}{\partial y} \right\} = 0$$

$$\mu x \frac{\partial^2 v}{\partial y^2} + \mu \frac{\partial}{\partial x} \left\{ x \frac{\partial u}{\partial y} \right\} - \mu x \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial y^2} = 0$$

$$u + \mu x \frac{\partial^2 u}{\partial y \partial x} + \mu \frac{\partial u}{\partial y} + x \frac{\partial \mu}{\partial x} \frac{\partial u}{\partial y} = 0$$

$$\mu x \frac{\partial}{\partial y} \left\{ \frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \right\} + \frac{\partial u}{\partial y} \frac{\partial}{\partial x} (\mu x) = 0$$

$$\frac{\partial}{\partial y} \left\{ \frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \right\} + \frac{\partial u}{\partial y} \frac{\partial \log(\mu x)}{\partial x} = 0$$

$$\frac{\partial}{\partial y} \left\{ \frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \right\} + \frac{\partial}{\partial y} \left\{ u \frac{\partial \log(\mu x)}{\partial x} \right\} = 0$$

$$\frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} + u \frac{\partial \log \mu x}{\partial x} = \text{const}_y = f_c(x, t)$$

selbstverständlich! nach $(\mathbb{R}) = \frac{1}{\mu} \frac{\partial \mu}{\partial x}$

$$\frac{\partial^2 v}{\partial y^2} = 2c_2 + 2\beta \cdot y c_3$$

$$\mu x (2c_2 + 2\beta c_3 y) = - \frac{\partial}{\partial x} \left(\mu x \frac{\partial u}{\partial y} \right)$$

$$\mu x \frac{\partial u}{\partial y} \frac{\partial^2 v}{\partial y^2} + \frac{\partial u}{\partial y} \frac{\partial}{\partial x} \left\{ \mu x \frac{\partial u}{\partial y} \right\} = 0$$

$$\frac{\partial^2 v}{\partial y^2} + \frac{\partial u}{\partial y} \frac{\partial}{\partial x} \left\{ \log \left(\mu x \frac{\partial u}{\partial y} \right) \right\} = 0$$

$$a \int_{-\infty}^{+\infty} e^{-yx^2} dx = 1 = \frac{a}{\sqrt{y}} \int_{-\infty}^{+\infty} e^{-y^2 x^2} dx \sqrt{y} = a \sqrt{\frac{\pi}{y}} = 1$$

$$a = \sqrt{\frac{y}{\pi}}$$

$$\sqrt{\frac{y}{\pi}} \int_0^x e^{-y^2 x^2} dx = \text{erf} \left(\frac{y x}{\sqrt{\pi}} \right)$$

$$= \frac{1}{\sqrt{\pi}} \int_0^{\frac{y x}{\sqrt{\pi}}} e^{-y^2 x^2} dy$$

$$\frac{x}{l_{50}} = y$$

$$l_{50} \sqrt{\frac{y}{\pi}} \int_0^{l_{50} y} e^{-y^2 x^2} dy = \frac{1}{2}$$

$$dy = \frac{1}{l_{50}}$$

$$y = \frac{1}{2} \left(\frac{\pi}{y} \right)$$

$$2y^2 = 1 \Rightarrow y = 0.5$$

$$l_{50} = 2.0845!$$

$$l_{50} \sqrt{\frac{y}{\pi}} \int_0^{l_{50}} e^{-y^2 x^2} dy = \frac{1}{2}$$

$$y l_{50} = x$$

$$f(x) = e^{-y^2 x^2}$$

$$\sqrt{\frac{y}{\pi}} \int_0^{l_{50}} e^{-y^2 x^2} dx = \frac{1}{4} = \frac{1}{\sqrt{\pi}} \int_0^{\frac{y l_{50}}{2}} e^{-y^2 x^2} dy$$

$$y = \frac{l_{50}}{2}$$

$$y = \frac{1}{4} \left(\frac{\pi}{y} \right)$$

~~erf~~

$$l_{50} = \frac{0.65546}{\sqrt{y}}$$

$$y = \frac{0.65546^2}{l_{50}^2}$$

$$\int_0^a e^{-x^2} dx = \frac{1}{\sqrt{\pi}} \int_0^{\frac{y}{\sqrt{\pi}}} e^{-\frac{y^2}{\pi}} dy$$

$$\frac{\sum x^2}{\sum N} = \xi \cdot 140$$

~~$\frac{y}{\sqrt{\pi}} = x$~~

$$\frac{x}{\sqrt{\pi}} = y$$

$$x = \frac{\sqrt{\pi} y}{1}$$

$$\int_0^{\infty} e^{-\left(\frac{\sqrt{\pi} y}{1}\right)^2} d\left(\frac{\sqrt{\pi} y}{1}\right)$$

$$\xi = \frac{\sum nx}{\sum n} = \frac{\int_0^{\infty} x e^{-x^2} dx \cdot \sqrt{\pi}}{\int_0^{\infty} e^{-x^2} dx \cdot \sqrt{\pi}}$$

$$l_{50} = 2 \cdot \xi \cdot 0.845$$

$$x \sqrt{\pi} = y$$

$$= \frac{\int_0^{\infty} y e^{-y^2} dy}{\int_0^{\infty} e^{-y^2} dy}$$

$$= \frac{\frac{-1}{2\sqrt{\pi}} e^{-y^2}}{\frac{1}{2\sqrt{\pi}}} = \frac{1}{\sqrt{\pi}}$$

$$l_{50} = \frac{1.69}{\sqrt{\pi}}$$

$$\frac{1}{\sqrt{\pi}} = 0.56419$$

0.40770	0.22780	
0.80103	0.39709	
81815	0.8880	0.6742
0.97591	1	l_{50}

Prob: $\frac{1}{2} = \frac{1}{\sqrt{2}} \int_0^{0.955} e^{-y^2} dy$ ~~0.955~~ 0.674

~~$$y - \frac{y^3}{3} + \frac{y^5}{1.2.5} - \frac{y^7}{1.2.3.5.7} + \frac{y^9}{1.2.3.4.6}$$~~

~~0.65760 - by approx 15.6.11
 0.97280 - 2 3 0.09393 : 3
3
110
40~~

0.28800 - 2 10 0.01941 : 10

0.60320 - 3 6.7 0.00401 : 6.7

0.91840 - 4 6.4.9 0.00083 : 6.4.9
24.9
216

0.45457	0.03131
0.00194	0.00009
<hr/> 0.45651	<hr/> 0.03140
-0.03140	
<hr/> 0.42511	

$$\sqrt{n} = 1.77245$$

$$\frac{\sqrt{n}}{2} = 0.88622$$

$$0.166666$$

$$0.041666$$

$$0.0046295$$

1

0.1

$$0.0046$$

$$0.3333$$

$$0.0001$$

$$0.0238$$

$$1.047$$

$$\frac{35}{0.85}$$

$$0.679$$

$$\int_0^1 (1+x+\frac{x^2}{2} + \dots) dx = 2 - 1 = 1$$

Part 1: $\frac{\sqrt{n}}{4} = 0.3371$

$$\begin{array}{r} 0.22789 \\ - 0.24857 \\ \hline 0.97932 \\ 95864 \end{array}$$

$$90946$$

$$0.95350$$

$$l_{50} = \frac{0.95350}{\sqrt{e}}$$

$$e = \frac{0.90916}{l_{50}^2}$$

2%

~~1.75~~

~~4~~

~~7%~~

~~1.5~~

~~10.5~~

~~16%~~

~~1.7~~

~~16~~

~~25%~~

~~0.51~~

~~12.5~~

1 . $\frac{15}{8}$

15

1 . $\frac{13}{8}$

13

$\frac{5}{2}$. $\frac{11}{8}$

27.5

$\frac{9}{2}$. $\frac{9}{8}$

40.5

7 . $\frac{7}{8}$

49

9 . $\frac{5}{8}$

45

12 . $\frac{3}{8}$

36

13 . $\frac{1}{8}$

13

50 l_{50}

$$\frac{239.0}{7.9} : 8 = \frac{30}{5}$$

$$\frac{3}{5} l_{50} = \frac{6}{5} a_{50}$$

$$a_{50} = \frac{5}{6} m = 0.83$$

Prob: 0.47675

$$\frac{\sqrt{x}}{4} = \int e^{-y^2} dy$$

0.1084

0.67830 - 1

0.03490 - 1

0.002463

0.39150 - 2

0.005599

0.74810 - 7

142

0.47675

- 0.0361 -

0.00246

0.0001

0.47921

0.0362

- 0.0362

0.4430

= 0.4431

stimmt!

~~kg~~ ^{kg} A

n_1

N_1

$v_1 = \frac{n_1}{N_1}$

~~kg~~ a

$a n_1 = 1$

c_1

~~kg~~ ^{kg} B

n_2

N_2

$v_2 = \frac{n_2}{N_2}$

b

$b n_2 = 1$

c_2

~~kg~~ ^{kg} Molekulargew.

= J. e. Atome

= J. e. Moleküle

= J. e. St. u. Molek.

= Störungswert

= spez. W. = c_2 / ρ_2

0,510822

~~kg~~ $a c_1 = b c_2$ (Dulong Petit)

$\frac{c_1}{n_1} = \frac{c_2}{n_2}$

(u = ~~v~~ ^{Strom} ~~gesch.~~)

$$c_1 = \frac{J n_1 a \Delta(u_1^2)}{2} = \frac{J \Delta(u_1^2)}{2} \quad c_2 = \frac{J n_2 b \Delta(u_2^2)}{2} = \frac{J \Delta(u_2^2)}{2}$$

$$\frac{c_1}{n_1} = J \frac{\Delta(u_1^2)}{2} \quad \frac{c_2}{n_2} = J \frac{\Delta(u_2^2)}{2}$$

Wenn gleiche Temp. so $\frac{\Delta(u_1^2)}{2} = \frac{\Delta(u_2^2)}{2}$

Besteht Gleichheit der Temp. darin, dass die mittl. kinetische Energie der Atome oder der Moleküle gleich ist?

Glaube des letzteren.

Dann ist Beweis in Nulls-D. mit Winkelmann falsch!

wenn man ~~kennt~~ ~~kennt~~ ~~kennt~~

$v = \text{gleiche Moleküle}$

des Temp. Gleichheit:

$$\frac{A(v_1^2)}{2} = \frac{B(v_2^2)}{2}$$

$$A = \frac{n_1 a}{N_1} \quad B = \frac{n_2 b}{N_2}$$

$$\frac{\overset{=1}{n_1} a(v_1^2)}{2 N_1} = \frac{\overset{=1}{n_2} b(v_2^2)}{2 N_2}$$

$$\frac{\Delta(v_1^2)}{N_1} = \frac{\Delta(v_2^2)}{N_2}$$

$$c_1 = \int \frac{N_1 A(v_1^2)}{2} \quad c_2 = \int \frac{N_2 B(v_2^2)}{2}$$

von Dulong Petit:

$$\frac{c_1}{n_1} = \frac{c_2}{n_2}$$

$$\int \frac{N_1 A(v_1^2)}{2 n_1} = \int \frac{N_2 B(v_2^2)}{2 n_2} \quad \frac{\Delta(v_1^2)}{n_1} = \frac{\Delta(v_2^2)}{n_2}$$

~~$n_1 a = n_2 b$~~
 ~~$n_1 a v_1^2 = n_2 b v_2^2$~~



$$\frac{n_1}{N_1} = \frac{n_2}{N_2}$$

das würde bei genauer Gültigkeit

des Dulong-Petit oder Gesetzes folgen,

(aller Körper bei derselben Temperatur)

das die Moleküle aus gleichem Atom zusammengesetzt sind!

wenn man dabei die räumliche Abit = 0 voraussetzt!

$$\begin{aligned}
 a c_1 : b c_2 &= a \Delta(v_1) : b \Delta(v_2) \\
 &= a N_1 \frac{\Delta(v_1)}{N_1} : b N_2 \frac{\Delta(v_2)}{N_2} \\
 &= a N_1 : b N_2 \\
 &= \frac{N_1}{n_1} : \frac{N_2}{n_2} = v_1 : v_2
 \end{aligned}$$

Also geben die Verhältnisse der Stoffmengen
bei Voraussetzung eines inneren Arbeits = 0 zugleich
die Verhältnisse der Zusammensetzung der Moleküle
aus Atomen!

Das alles bei Voraussetzung von gleichartigen Atomen
wenn wir jetzt ^{vor} voraussetzen

A besteht aus v_1 Atomen mit Atomgew. a_1

• v_2 " " a_2

B μ_1 " b_1

μ_2 " b_2

so folgt ~~zusammensetzung~~ ~~Arbeits = 0~~

$$A = v_1 a_1 + v_2 a_2 \quad B = \mu_1 b_1 + \mu_2 b_2$$

Temp. Gleichheit:

144

$$A \Delta v_1^2 = B \Delta v_2^2$$

$$c_1 = \int \frac{A N_1 \Delta v_1^2}{2}$$

$$c_2 = \int \frac{B N_2 \Delta v_2^2}{2}$$

$$= \int \frac{\Delta v_1^2}{2}$$

$$= \int \frac{\Delta v_2^2}{2}$$

so folgt: $A c_1 = B c_2$ } = Überlegung allgemeinster
Jedes durch den Diloy 'Polst' in
und Normen 'hat' }

~~$$(q_1 v_1 + q_2 v_2) c_1 = (q_1 \mu_1 + q_2 \mu_2) c_2$$~~

~~$$\text{da besteht: } (q_1 v_1 + q_2 v_2) c_1 = q_1 c_1 (v_1 + v_2)$$~~

~~$$2B: (q_1 v_1 + q_2 v_2) c_3 = q_1 \mu_1 c_1$$~~

~~$$= q_2 \mu_2 c_2$$~~

~~$$c_3 = \frac{q_1 \mu_1 c_1}{q_1 v_1 + q_2 v_2} = \frac{q_2 \mu_2 c_2}{q_1 v_1 + q_2 v_2}$$~~

es würde also aus den Annahmen

1. Gleichheit der Temp. zweier Körper besteht darin, dass die kinetische Energie der Proben gleich ist
2. Innere Arbeit = 0 also spec. Wärme = Arbeit zur Vermehrung der kinetischen Energie der Proben

folgen, dass die Probenwärme aller Körper gleich ist;

Aus nicht der Fall, also falsch!

Annahme 1 ist richtig dann ist. Angenommen Seite 1 die
 die Annahme 2 zu corrigieren

$$c_1 = \int \frac{N_1 A \Delta(v_1^2)}{2} + \int \frac{N_1 A \Delta(v_1^2) P_1}{2} + \int \frac{N_1 A \Delta(v_1^2) P_1}{2}$$

$$c_2 = \int \frac{N_2 B \Delta(v_2^2)}{2} + \int \frac{N_2 B \Delta(v_2^2) P_2}{2} \quad P_1 = \dots \text{ pro 1 Molent}$$

Temperatur Gleichheit erfordert:

$$A \Delta(v_1^2) = B \Delta(v_2^2)$$

~~Übung Petri: a c_1 = b c_2~~

$$A \Delta(v_1^2) \left(\frac{\int N_1}{N_1} + \frac{\int N_1 P_1}{N_1} \right) = c_1$$

$$\frac{c_1}{N_1 (1 + \frac{P_1}{N_1})} = \frac{c_2}{N_2 (1 + \frac{P_2}{N_2})}$$

$$\frac{c_1}{N_1 (1 + \frac{P_1}{N_1})} = \frac{c_2}{N_2 (1 + \frac{P_2}{N_2})}$$

Übung Petri: a c_1 = b c_2

~~a N_1 (1 + \frac{P_1}{N_1}) = b N_2 (1 + \frac{P_2}{N_2})~~

$$a N_1 (1 + \frac{P_1}{N_1}) = b N_2 (1 + \frac{P_2}{N_2})$$

$$\frac{1 + \frac{P_1}{N_1}}{v_1} = \frac{1 + \frac{P_2}{N_2}}{v_2}$$

~~... unter Voraussetzung~~

~~$\frac{1+P_1}{V_1} = \frac{1+P_2}{V_2}$~~

also bei gleicher Ertragskraft des Darlehens
folgt nicht dass die innere Arbeit = 0 ist
(Winkelmann, Wöllner etc.) !

sondern, aus je mehr Atome ein Solenoid
zusammengesetzt ist, desto größer ist das Verhältnis
der inneren Arbeit eines Solenoids zu dem Turnverhältnis
des L K desselben.

Bei einem Solenoid wäre innere Arbeit
proportional der L K (selbstverst.)
 $V_1 = V_2 = 1$

wenn Index in der Bedeutung = v

$$\frac{1+P_1}{1} = \frac{1+P_2}{2} = \frac{1+P_3}{3} = \frac{1+P_4}{4} = \frac{1+P_5}{5}$$

Bei Körpern, deren Solenoid aus gleichviel Atomen
zusammengesetzt sind, ist das ^{Verhältnis der} inneren Arbeit eines
Solenoids zu dem Turnverhältnis des L K desselben gleich

(mit N multipliziert)
oder mit anderen Worten:

$$P_1 = \frac{2V_1}{A \cdot V_1^2}$$

$$\frac{V_1}{A \cdot V_1^2} = \frac{V_2}{B \cdot V_2^2}$$

oder weil Nenner gleich
 $V_1 = V_2$ das heißt:

bei genauer Betrachtung der DP Situation
 müsste bei Körper, deren Molekül aus gleichviel Atomen
 bestehen, die innere Arbeit pro Molekül bei gleicher Temperatur-
 erhöhung gleich sein, also gesamt innere Arbeit
 proportional der Anzahl Moleküle resp. umgekehrt prop.
 dem ~~der~~ Molekulargewicht.

bei verschiedenen atomigen Körpern:

je mehr atomig desto größer die innere Arbeit (wählt $k=1$)
 pro Molekül (aber nicht proportional)

(dies stimmt insofern als desto größer
 Moleküle wenn je mehr die Körper zusammengesetzt sind)

wenn P sehr groß ~~ist~~^{wäre}, so dass die erste Potenz zu
 vernachlässigen so würde folgen:

$$C_1 = J N_1 A_1 \frac{v_1^2}{2} P_1$$

$$\frac{C_1}{N_1 P_1} = \frac{C_2}{N_2 P_2} =$$

$$A \frac{v_1^2}{2} = B \frac{v_2^2}{2} \quad \text{also}$$

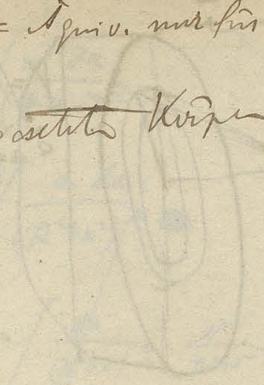
$$\frac{C_1}{N_1 P_1} = \frac{C_2}{N_2 P_2}$$

$$\frac{P_1}{V_1} = \frac{P_2}{V_2}$$

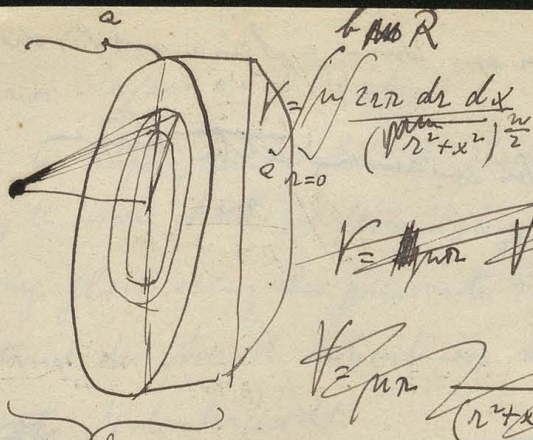
innere Arbeit proportional der Anzahl Atome pro Mole.
 also gesamt innere Arbeit proportional der Gesamtzahl
 der Atome also umgekehrt proportional dem
 Molekulargewicht

wil davon aber spec. Wärm. = $\int q_{in}$ nur für die
innere

bei ist dies bei zusammengeordneten Körper
Joules des Gerets ?



[Faint, illegible handwritten notes and mathematical scribbles covering the lower two-thirds of the page.]



$$V = \int_a^b \frac{2\pi r (r+x)}{1-\frac{n}{2}} dx =$$

$$= \frac{2\pi R}{2-n} \int_a^b \left(\frac{R^2+x^2}{(R^2+x^2)^{\frac{n}{2}}} - x \right) dx$$

$$\int \frac{R^2+x^2}{(R^2+x^2)^{\frac{n}{2}}} dx = (R^2+x^2)^{\frac{1-n}{2}} x - \int \frac{x(1-\frac{n}{2})2x dx}{(R^2+x^2)^{\frac{n}{2}}}$$

$$R^2 \int \frac{dx}{(R^2+x^2)^{\frac{n}{2}}} + \int \frac{x^2 dx}{(R^2+x^2)^{\frac{n}{2}}} = (R^2+x^2)^{\frac{2-n}{2}} x - (2-n) \int \frac{x^2 dx}{(R^2+x^2)^{\frac{n}{2}}}$$

$$I_1 = \frac{(R^2+x^2)^{\frac{2-n}{2}} x}{R^2} - \frac{(3-n)}{R^2} I_2$$

$$\int \frac{dx}{(R^2+x^2)^{\frac{n}{2}}} = \frac{x}{(R^2+x^2)^{\frac{n}{2}}} + n \int \frac{x^2 dx}{(R^2+x^2)^{\frac{n}{2}+1}}$$

$$-n \int \frac{dx}{(R^2+x^2)^{\frac{n}{2}+1}} = \int \frac{n dx}{(R^2+x^2)^{\frac{n}{2}+1}}$$

$$(1-n) \int_1 = \frac{x}{(R^2+x^2)^{\frac{n}{2}}} + n \int \frac{dx}{(R^2+x^2)^{\frac{n}{2}}} \frac{x^2 - R^2}{R^2+x^2}$$

$$= \frac{x}{(R^2+x^2)^{\frac{n}{2}}} - n R^2 \int \frac{dx}{(R^2+x^2)^{\frac{n}{2}+1}}$$

$$n = n-2$$

$$(1-n+2) \int_1 = \frac{x}{(R^2+x^2)^{\frac{n-2}{2}}} - (n-2) R^2 \int \frac{dx}{(R^2+x^2)^{\frac{n}{2}}}$$

$$\int = R^2 \int_1 + \int_2 = (R^2+x^2)^{\frac{2-n}{2}} x - (3-n) \int_2$$

$$(3-n) \int = \frac{x}{(R^2+x^2)^{\frac{n-2}{2}}} - (n-2) R^2 \int_1$$

$$\int = \frac{x}{(R^2+x^2)^{\frac{n-2}{2}}} - (n-2) (R^2+x^2)^{\frac{2-n}{2}} x - (n-2)(3-n) \int_2$$

$$(3-n) \int_2 = (R^2+x^2)^{\frac{2-n}{2}} x (n-2) + \dots$$

$$\int_2 = x (R^2+x^2)^{\frac{2-n}{2}}$$

$$\frac{d}{dx} = (R^2+x^2)^{\frac{2-n}{2}} + (2-n) x^2 (R^2+x^2)^{-\frac{n}{2}}$$

$$= (R^2+x^2)^{-\frac{n}{2}} [2x + (3-n)x^2]$$

Voraussetz. $n > 2$

$$v_1 \quad v_2$$

$$a v_1 v_1^2 = b v_2 v_2^2$$

$$a u_1^2 = b u_2^2$$

~~$$c = \int \frac{a v_1 v_1^2 + b v_2 v_2^2}{\Delta}$$

$$= \int \frac{a v_1 v_1^2 (1 + \lambda)}{\Delta}$$~~

$$\frac{u_1^2}{v_1 v_1^2} = \frac{u_2^2}{v_2 v_2^2} = \lambda = 1 \quad \text{c. d. h. } \lambda = 1$$

$$u_1^2 = \lambda v_1 v_1^2$$

$$c_1 = \int N_1 \Delta(a v_1 v_1^2) + \int N_2 v_2 \Delta(a u_1^2)$$

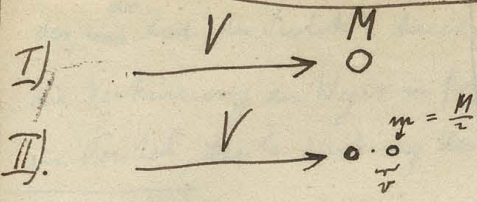
~~$$= \int N_1 \Delta(a v_1 v_1^2) [1 + \lambda v_1] = \int N_1 \Delta(a v_1 v_1^2) (1 + v_1)$$~~

$$c_1^x$$

$$\frac{c_1}{N_1 (1 + v_1)} = \frac{c_2}{N_2 (1 + v_2)} \quad \vee \quad P_1 = v_1$$

$$\frac{1 + v_1}{v_1} = \frac{1 + v_2}{v_2} \quad \text{c. d. h. } \frac{1}{v_1} = \frac{1}{v_2} \quad \vee \quad v_1 = v_2$$

en $\frac{x}{p}$ e, λ konstante so λ spec. Vol. von x^2



$$E_I = \frac{MV^2}{2}$$

$$E_{II} = \frac{m}{2} \left[(V+v)^2 + (V-v)^2 \right]$$

$$= \frac{M}{2} [V^2 + v^2] = E_I + 2 \cdot \frac{m v^2}{2}$$

$$E_n = \frac{m_n}{2} \sum (V+v_n)^2$$

~~$$\sum (V+v) = Vn$$~~

$$= \frac{m_n}{2} \left(\sum V^2 + \sum v_n^2 \right) \quad \sum v_n = 0 \text{ da } \sum v_n = 0$$

$$E_n = E_I + n \frac{m v^2}{2} = E_I \left[1 + \left(\frac{v}{V} \right)^2 \right]$$

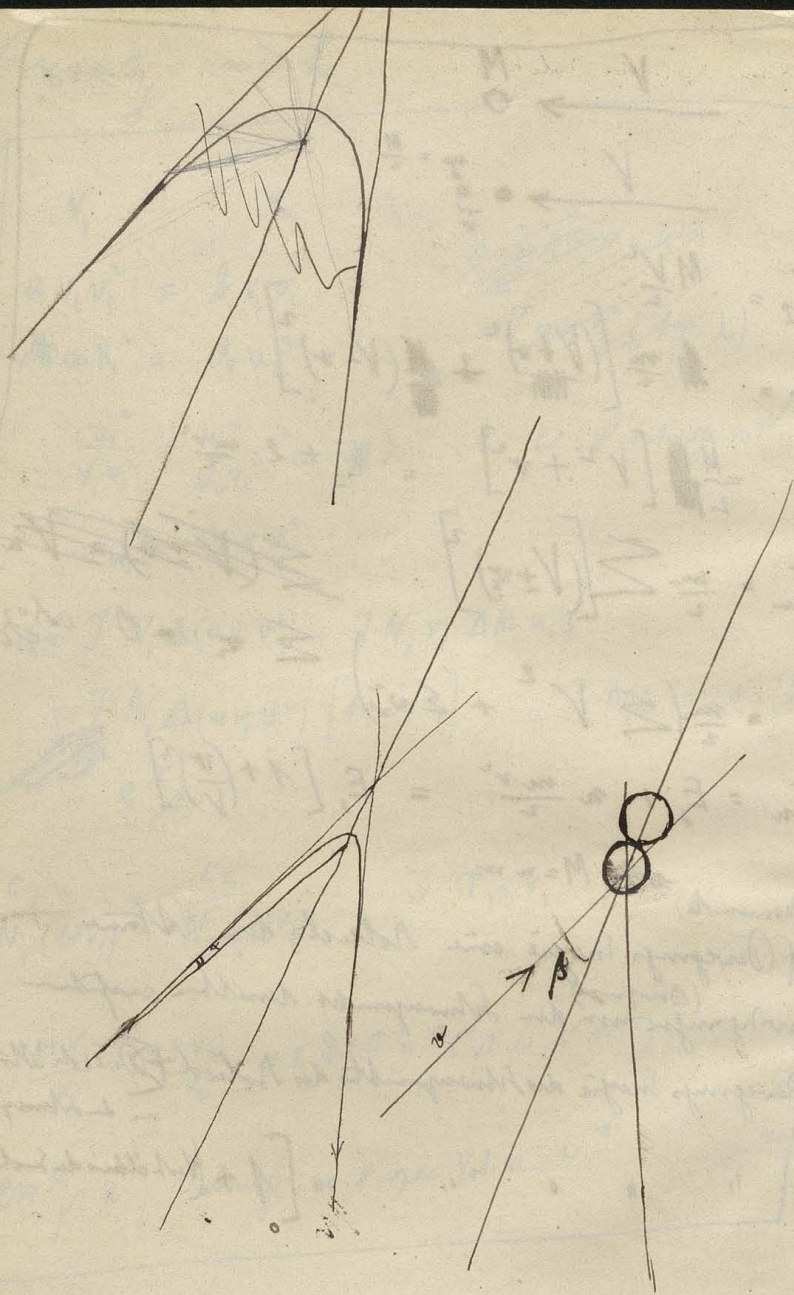
~~$$M = n m$$~~

(Gesamt)

die Bewegungsenergie eines Rollen, der Stone
 Schwingungen um den Schwerpunkt derselben erfordern
 (Bewegung)

= Bewegungsenergie des Schwerpunktes des Rollen + $\frac{1}{2} I \omega^2$ d. Stone
 um den Schwerpunkt

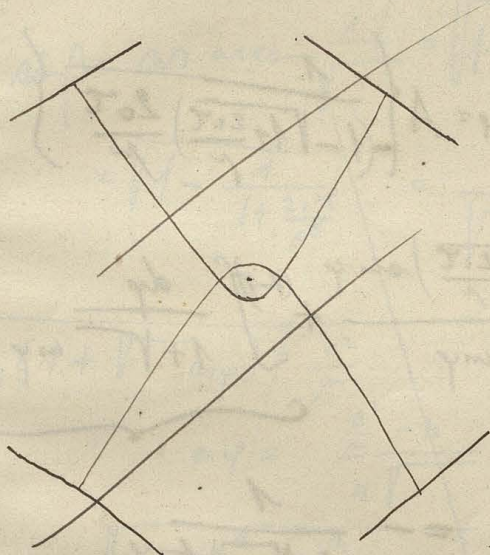
= " " " " " $\left[1 + \left(\text{Verhältnis des Schwerpunktes} \right) \right]$



$$\beta = \varphi_0 - \alpha$$

das, was ^{ds} Red. der Winkel berechnet wird, ist theilweise 149
 die Verkürzung des Weges in Folge der größten Seitenindignität
 im Perikel, theils Wirkung der Dispersion.

10.



// $c: v_0, n_0, \alpha$

~~$c = \lambda \frac{d\phi}{dt}$~~

bei $g = \text{const.}$ S. 101!

$$C = \frac{v_0^2}{2} - \frac{\mu}{n_0}$$

$$z = \frac{\frac{c^2}{\mu}}{1 + \sqrt{1 + \frac{2c^2 C}{\mu^2}} \cos \varphi}$$

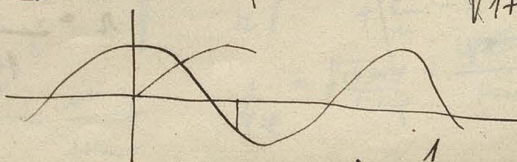
$$\mu = k(m + m_1)$$

$c = \frac{1}{2} n_0 v_0 \sin \alpha$

$$\cos \varphi \sqrt{1 + \frac{2c^2 C}{\mu^2}} = -1$$

$$\varphi = \arccos \left(\frac{-1}{\sqrt{1 + \frac{2c^2 C}{\mu^2}}} \right)$$

$k = \infty$



$$\varphi_0 = \frac{\pi}{2} + \arcsin \frac{1}{\sqrt{1 + \frac{2c^2 C}{\mu^2}}}$$

$$r = \frac{\mu}{1 + \epsilon \cos \varphi}$$

$$\frac{\frac{c^2}{\mu^2}}{\left(1 + \sqrt{1 + \frac{2c^2}{\mu}} \cos \varphi\right)^2} d\varphi = c dt$$

$$\int \frac{A}{(1 + B \cos \varphi)^2} d\varphi = A \left\{ \frac{1}{\left(1 - \sqrt{1 + \frac{2c^2}{\mu}}\right) \frac{2c^2}{\mu}} \right\}$$

$$\left\{ \frac{\left(1 + \frac{2c^2}{\mu} - \sqrt{1 + \frac{2c^2}{\mu}}\right) \sin \varphi}{1 + \sqrt{1 + \frac{2c^2}{\mu}} \cos \varphi} + \int \frac{d\varphi}{1 + \sqrt{1 + \frac{2c^2}{\mu}} \cos \varphi} \right\}$$

$$= - \frac{1}{(1 - \sqrt{1 + \frac{2c^2}{\mu}}) \tan \frac{\varphi}{2}}$$

$$= A \left\{ \frac{1}{\frac{2c^2}{\mu} (\sqrt{1 + \frac{2c^2}{\mu}} - 1)} \right\} \left\{ \left(1 + \frac{2c^2}{\mu} - \sqrt{1 + \frac{2c^2}{\mu}}\right) \sin \varphi \right\}$$

Perihel distance $r(\text{min}) = \varphi = 0$

$$r = \frac{\frac{c^2}{\mu}}{1 + \sqrt{1 + \frac{2c^2}{\mu}}}$$

~~Handwritten scribbles~~
~~Handwritten scribbles~~
lim
r₀ = ∞

$$2 \left(\frac{r_0 \sin \varphi_0 - r \sin \varphi}{\sin \beta} - \frac{r}{r_0} \right) = \gamma \mu^2 \beta$$

$$\sin \beta = \cos \arcsin \frac{1}{\mu} = \sqrt{1 - \frac{1}{\mu^2}}$$
$$= \sqrt{1 - \frac{1}{1 + \frac{2c^2 r}{\mu^2}}} = \frac{\frac{2c^2 r}{\mu^2}}{\sqrt{1 + \frac{2c^2 r}{\mu^2}}}$$

$$r \mu (1 + \sqrt{\cos \varphi}) = \frac{c^2}{\mu}$$
$$\cos \varphi = \frac{\frac{c^2}{\mu} - r}{r \sqrt{\dots}}$$

$$\sin \varphi = \sqrt{1 - \frac{\left(\frac{c^2}{\mu} - r\right)^2}{r^2 \left(1 - \frac{2c^2 r}{\mu^2}\right)}} = \frac{\sqrt{\frac{2c^2 r}{\mu} - \frac{c^4}{\mu^2} - \frac{2c^2 r^2}{\mu^2}}}{r \sqrt{\dots}}$$
$$= \frac{\frac{c^2}{\mu} \left(2r - \frac{c^2}{\mu} - 2c^2 r\right)}{r \sqrt{\dots}}$$

$$\frac{1}{\mu} = \frac{\sqrt{1 - \cos \varphi}}{1 + \cos \varphi} = \frac{r \sqrt{-\frac{c^2}{\mu} + r}}{r \sqrt{+\frac{c^2}{\mu} - r}}$$
$$= \frac{\sqrt{1 - \cos \varphi}}{1 + \cos \varphi} = \frac{\sin \varphi}{1 + \cos \varphi}$$
$$\frac{1}{\mu} = \frac{\sqrt{1 + \cos \varphi}}{1 - \cos \varphi} = \frac{\sin \varphi}{1 - \cos \varphi}$$

$$\left\{ \frac{\left(1 + \frac{2c^2}{\mu} + \sqrt{\dots}\right) \sin \varphi}{1 + \sqrt{\dots} \cos \varphi} - \frac{\sin \varphi}{1 - \cos \varphi} \right\}$$

$$= \frac{2c^2}{\mu} \sin \varphi + \sqrt{\dots} \sin \varphi - \left(\dots \right) \sin \varphi \cos \varphi - 2 \sqrt{\dots} \sin \varphi \cos \varphi$$

$$\frac{(1 + \sqrt{\dots} \cos \varphi)(1 - \cos \varphi)}{(1 + \sqrt{\dots} \cos \varphi)(1 - \cos \varphi)}$$

$$= \frac{\frac{2c^2}{\mu} + \sqrt{1 + \frac{2c^2}{\mu}} - \left(1 + \frac{2c^2}{\mu} + 2\sqrt{1 + \frac{2c^2}{\mu}}\right) \cos \varphi}{(1 + \sqrt{\dots} \cos \varphi)(1 - \cos \varphi)}$$

$$\sin \varphi_0 = \sin(\beta + \alpha)$$

$$= \sin \beta \cos \alpha + \cos \beta \sin \alpha$$

$$= \cos \alpha \sqrt{\frac{2c^2}{\mu^2}} + \sin \alpha \frac{1}{\sqrt{\dots}}$$

$$\sin \varphi_0 = \frac{\frac{2c^2}{\mu^2} \cos \alpha - \sin \alpha}{\sqrt{1 + \frac{2c^2}{\mu^2}}}$$

$$= \frac{C}{\mu^2} \frac{1}{2} r_0^2 v_0^2 \sin^2 \alpha \cos \alpha - \sin \alpha$$

$$= \frac{\sin \alpha}{\sqrt{\dots}} \left(\frac{C}{4\mu^2} r_0^2 v_0^2 \sin^2 \alpha - 1 \right)$$

$$\cos \varphi_0 = \frac{1 + \frac{2c^2 C}{\mu^2} - \frac{4c^4 C^2}{\mu^4} \cos^2 \alpha + \frac{4c^2 C^2}{\mu^2} \cos \alpha \sin \alpha - \sin^2 \alpha}{1 + \frac{2c^2 C}{\mu^2}} \quad 151$$

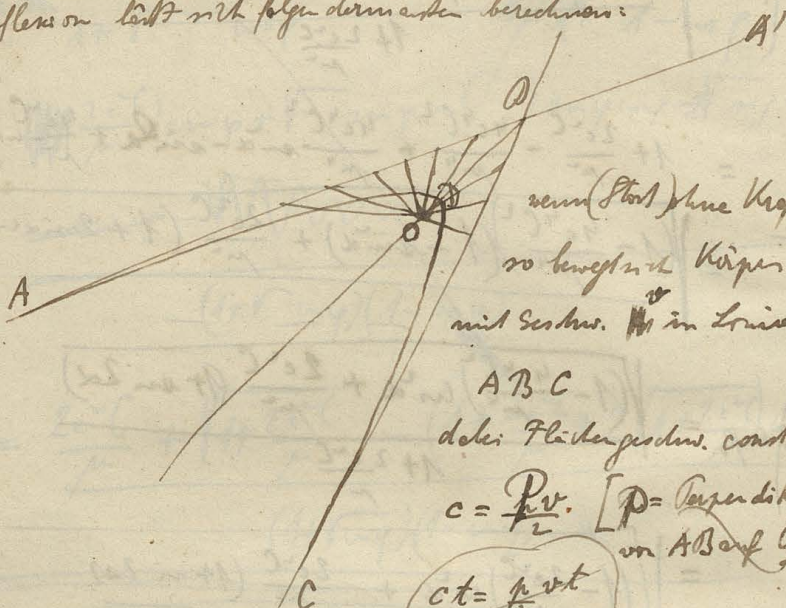
$$= \frac{1 + \frac{2c^2 C}{\mu^2} - \frac{4c^4 C^2}{\mu^4} + \frac{4c^2 C^2}{\mu^2} \sin \alpha - \sin^2 \alpha + \frac{4c^2 C^2}{\mu^2} \cos \alpha \sin \alpha}{1 + \frac{2c^2 C}{\mu^2}}$$

$$= \frac{\left(1 - \frac{4c^4 C^2}{\mu^4}\right) (1 - \sin^2 \alpha) + \frac{2c^2 C}{\mu^2} (1 + 2 \sin \alpha \cos \alpha)}{1 + \frac{2c^2 C}{\mu^2}}$$

$$\sin \varphi_0 = \frac{\sqrt{\left(1 - \frac{4c^4 C^2}{\mu^4}\right) \cos^2 \alpha + \frac{2c^2 C}{\mu^2} (1 + \sin 2\alpha)}}{1 + \frac{2c^2 C}{\mu^2}}$$

$$= \sqrt{\left(1 - \frac{2c^2 C}{\mu^2}\right) \cos^2 \alpha + \frac{2c^2 C}{\mu^2} (1 + \sin 2\alpha)} \quad \frac{1 + \frac{2c^2 C}{\mu^2}}$$

Die Drehbewegung durch die $\frac{1}{2}$ Kraft gegenüber einer Stoß-
Reflexion läßt sich folgen demnach berechnen:



wenn (Stoß) keine Kraft
so bewegt sich Körper
mit Geschw. v in Linie

ABC

daher: Flächengeschw. constant

$$c = \frac{p \cdot v}{2} \quad [p = \text{Perpendikel von } A \text{ auf } O]$$

$$ct = \frac{p}{2} \cdot vt$$

$$F = \frac{p}{2} \cdot s$$

also Zeit = $\frac{\text{Sector Fläche}}{\text{Flächengeschw.}}$

dagegen bei Umdrehungen zwar dieselbe Flächengeschwindigkeit,
aber es entfällt die Fläche zwischen Aequant. und Hyperbel.

$$\text{also Zeitgeraum} = \frac{\text{Fläche } ABCD}{\text{Flächengeschw.}}$$

$$p = \frac{r}{1 + \epsilon \cos \varphi}$$

$$F = \int \frac{p^2 d\varphi}{2} = \frac{r^2}{2} \int \frac{d\varphi}{(1 + \epsilon \cos \varphi)^2}$$

$$\frac{x^2}{a^2} - \frac{y^2}{b^2} = 1$$

152

$$y = \frac{b}{a} \sqrt{x^2 - a^2}$$

$$\int y dx = \frac{b}{a} \int \sqrt{x^2 - a^2} dx$$

$$I = x \sqrt{x^2 - a^2} - \int \frac{x^2}{\sqrt{x^2 - a^2}} dx$$
$$= I + \int \frac{a^2 dx}{\sqrt{x^2 - a^2}}$$

$$2I = x \sqrt{x^2 - a^2} - a^2 \int \frac{dx}{\sqrt{x^2 - a^2}}$$

$$I = \frac{x}{2} \sqrt{x^2 - a^2} - \frac{a^2}{2} \log(x + \sqrt{x^2 - a^2})$$

$$\int y dx = \frac{b}{2} x \sqrt{\left(\frac{x}{a}\right)^2 - 1} - \frac{ab}{2} \log\left(\frac{x}{a} + \sqrt{\left(\frac{x}{a}\right)^2 - 1}\right)$$
$$= \frac{xy}{2} - \frac{ab}{2} \log\left(\frac{x}{a} + \frac{y}{b}\right)$$

$$\text{dies Fläche AOCB} = \frac{ab}{2} \log\left(\frac{x}{a} + \frac{y}{b}\right) \Big|_{x=0}^{x=ab}$$
$$= \frac{ab}{2} + \frac{ab}{2} \log\left(\frac{x}{a} + \frac{y}{b}\right) \Big|_{x=0}^{x=ab} = \infty \dots$$

Zeitpunkt $\rightarrow \infty$; daher muss man als obere Grenze die mittlere ~~Weglänge~~ ^{Entfernung} l einführen

$$\frac{F_{\text{AOCB}}}{2} = \frac{ab}{2} \left\{ 1 + \log\left(\frac{l}{a} + \frac{y_l}{b}\right) \right\} = \frac{ab}{2} \left\{ 1 + \log \frac{l}{a} + \log\left(1 + \frac{a}{l} \frac{y_l}{b}\right) \right\}$$
$$\text{(angenähert)} \quad \frac{ab}{2} \left\{ 1 + \log \frac{2l}{a} \right\}$$

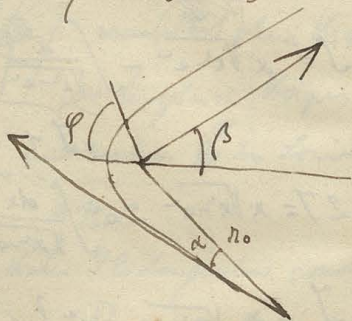
$$l = \text{Entfernung} \cdot \frac{a}{\sqrt{a^2 + b^2}}$$

$$h = \frac{\frac{c^2}{\mu}}{1 + \cos \varphi \sqrt{1 + \frac{2c^2}{\mu^2}}}$$

$$c = \frac{r_0 v_0 \sin \alpha}{2}$$

$$l = \frac{v_0^2}{2} - \frac{\mu}{r_0}$$

$$\mu = k(m + m_1)$$



$$\beta = 180^\circ - \varphi$$

$$\varphi = 180^\circ + \arcsin \frac{1}{\sqrt{1 + \frac{2c^2}{\mu^2}}}$$

$$\text{daraus: } \beta = \arccos \frac{1}{\sqrt{1 + \frac{2c^2}{\mu^2}}}$$

$$\cos \beta = \frac{1}{2}$$

$$\sqrt{a^2 + b^2} = c$$

$$(r)_{\varphi=0} = \sqrt{a^2 + b^2} - a = \frac{\frac{c^2}{\mu}}{1 + \sqrt{1 + \frac{2c^2}{\mu^2}}} = a \left[\sqrt{1 + \left(\frac{b}{a}\right)^2} - 1 \right]$$

$$\frac{b}{a} = \tan \beta = \tan \arccos \frac{1}{\sqrt{1 + \frac{2c^2}{\mu^2}}} = \frac{\sqrt{1 - \frac{1}{1 + \frac{2c^2}{\mu^2}}}}{\frac{1}{\sqrt{1 + \frac{2c^2}{\mu^2}}}} = \sqrt{\frac{2c^2}{\mu^2}}$$

$$= \sqrt{\frac{1 + \frac{2c^2}{\mu^2} - 1}{1 + \frac{2c^2}{\mu^2}}}$$

$$= \sqrt{\frac{\frac{2c^2}{\mu^2}}{1 + \frac{2c^2}{\mu^2}}} = \sqrt{\frac{2c^2}{\mu^2}}$$

$$\frac{c^2}{\mu} = \frac{c^2}{\mu} \sqrt{1 + \frac{2c^2}{\mu}} = \mu \sqrt{1 + \frac{2c^2}{\mu}} - 1$$

$$\frac{c^2}{\mu} = \frac{2c^2}{\mu} \quad \text{153}$$

$$Q = \frac{\mu}{2c}$$

$$\frac{c^2}{\mu} = \mu \left(\sqrt{1 + \frac{2c^2}{\mu}} - 1 \right)$$

Wen sich auf flgedichten =

$$Q = n_0 = \sin \alpha : \sin \beta$$

$$Q = \lim_{n_0 \rightarrow \infty} \frac{n_0 \sin \alpha}{\sin \beta}$$

$$\sin \beta = \sqrt{\frac{2c^2}{1 + \frac{2c^2}{\mu}}}$$

$$\frac{1}{\sin \beta} = \sqrt{\frac{1 + \frac{2c^2}{\mu}}{2c^2}} = \frac{1}{\mu} \sqrt{1 + \frac{1}{\frac{2c^2}{\mu}}} = \frac{1}{\mu} \left[1 + \frac{1}{\frac{2c^2}{\mu}} + \dots \right]$$

$$Q = \lim_{\mu \rightarrow \infty} \frac{n_0 \sin \alpha}{\mu} \left(1 + \frac{1}{\frac{2c^2}{\mu}} \right)$$

$$\lim = 0$$

$$= \frac{2c}{v_0 \mu}$$

$$Q = \frac{\mu}{2c}$$

$$\sqrt{a^2 + b^2} = \sqrt{\frac{\mu^2}{4c^2} + \frac{c^2}{2c}} = \frac{\sqrt{\mu^2 + 2c^2}}{2c}$$

$$b = \sqrt{\frac{c^2}{2c}}$$

$$= \frac{c}{\sqrt{2c}}$$

$$\frac{\sqrt{a^2 + b^2}}{a} = z = \sqrt{1 + \frac{2c^2}{\mu^2}}$$

stammt!

$$\Delta T = \frac{F}{c} = \frac{ab}{c} \left\{ 1 + \log \frac{2l}{a} \right\}$$

$$= \frac{\mu}{\sqrt{(2c)^3}} \left\{ 1 + \log 2 \right\}$$

l = Entfernung des Punktes vom Focus = Entfernung vom
 gemeinsamen Scheitelpunkt = doppelte Entfernung der beiden St.

(angesehen)

$$l = \left\{ \sqrt{a^2 + b^2} + \text{Entfernung vom Focus} \right\} \left\{ \frac{a}{\text{Fest.}} \right.$$

$\underbrace{\hspace{10em}}_{\text{halbe doppelte Entfernung der beiden St. von einander}} \quad \text{"Scheitelpunkt"}$

$$= a + \frac{a \lambda}{2 \sqrt{a^2 + b^2}} = a \left\{ 1 + \frac{\lambda}{2 \sqrt{a^2 + b^2}} \right\}$$

$$= a \left\{ 1 + \frac{\lambda}{2 \frac{\mu}{2c} \sqrt{1 + \frac{2c^2}{\mu^2}}} \right\}$$

$$\Delta T = \frac{\mu}{\sqrt{(2c)^3}} \left\{ 1 + \log 2 \left\{ 1 + \frac{\lambda}{\frac{\mu}{c} \sqrt{1 + \frac{2c^2}{\mu^2}}} \right\} \right\}$$

$$\left[\frac{\mu}{2c} \sqrt{1 + \frac{2c^2}{\mu^2}} - 1 \right] \frac{2c}{\mu} = \dots$$

$$\dots = \dots$$

$$\dots = \dots$$

Druck $\frac{\lambda}{r \cdot V}$ ist = $\frac{\text{Erhöhung der beiden } \sigma}{\text{Erweite. der Dohr}}$ also sehr groß 154

daher ΔT

$$\frac{\Delta T}{r} = \frac{\mu}{(2C)^{\frac{3}{2}}} \left\{ 1 + \log 2 + \log \frac{\lambda}{C \sqrt{1 + \frac{2C^2}{r^2}}} + \log \left(1 + \frac{\frac{\mu}{C} \sqrt{V}}{\lambda} \right) \right\}$$

$$= \frac{\mu}{(2C)^{\frac{3}{2}}} \left\{ 1 + \log \frac{2\lambda}{\frac{\mu}{C} \sqrt{1 + \frac{2C^2}{r^2}}} + \log (1 + \delta) \right\}$$

$\neq 0$

$$\longrightarrow \lambda = 2r_0$$

$$\Delta T = \frac{\mu}{(2C)^{\frac{3}{2}}} \left\{ 1 + \log \frac{4r_0 C}{\frac{\mu}{\epsilon}} \right\}$$

$$\epsilon = \sqrt{1 + \frac{2C^2}{\mu^2}} = \sqrt{1 + \frac{r_0^2 v_0^2 \sin^2 \alpha}{2 \mu^2} \left(\frac{v_0^2}{2} - \frac{\mu}{r_0} \right)}$$

$$\mu = 2 \text{ km}$$

$$\epsilon = \sqrt{1 + \frac{r_0^2 v_0^2 \sin^2 \alpha}{8 \text{ km}^2 \text{ m}^2} \left(\frac{v_0^2}{2} - \frac{2 \text{ km}}{r_0} \right)}$$

wenn man die vertikale Entfernung des Focus von der Asymptote einführt = $b = r_0 \sin \alpha$

$$c = \frac{v_0 b}{2}$$

$$\epsilon = \sqrt{1 + \frac{v_0^2 b^2}{2 \mu^2} \left(\frac{v_0^2}{2} - \frac{\mu}{r_0} \right)}$$

to work as shown

guiding



$$y = v_0 r$$

$$a = \frac{v_0}{r}$$

$$c = v_0 r$$

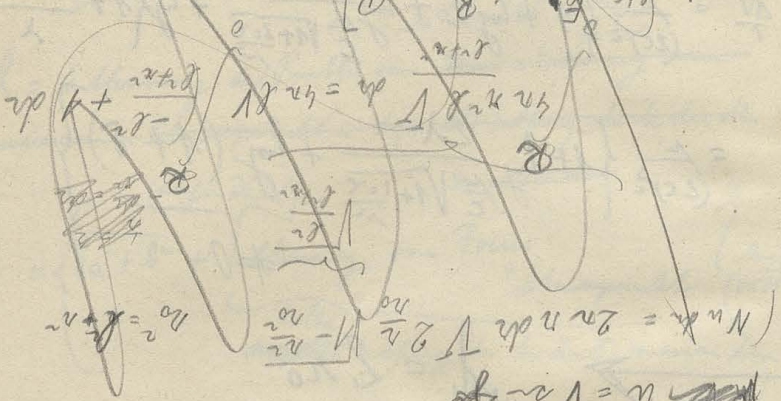
$$l = \frac{v_0 r}{\omega}$$

$$v_0 = \frac{c}{\omega} = \frac{v_0}{\omega}$$

$$u = v_0 \omega$$

$$(N u)_r = 2\pi n d h \sqrt{2\pi} \frac{h \omega}{h \omega} \frac{1}{\sqrt{2\pi}}$$

$$D_0 = k + n$$



$$= + 2\pi \left[1 + \frac{l}{R} \cos \frac{\theta}{2} \right]$$

: $2\pi R V$

$$\frac{w \cos \theta}{\sqrt{2} w} \lambda = \frac{R \lambda}{\sqrt{2} \lambda} \lambda = \lambda$$

~~For $\theta = 90^\circ$, $\lambda = \frac{R \lambda}{\sqrt{2} \lambda} \lambda = \frac{R}{\sqrt{2}} \lambda$~~

~~For $\theta = 0^\circ$, $\lambda = \frac{R \lambda}{\sqrt{2} \lambda} \lambda = \frac{R}{\sqrt{2}} \lambda$~~

$$\frac{w}{\sqrt{2} w} = \frac{w}{\sqrt{2} w} = \frac{1}{\sqrt{2}} = \frac{1}{\sqrt{2}} \neq \frac{1}{\sqrt{2}}$$

$$(\cos \theta / \sqrt{2} - \cos \theta / \sqrt{2}) \lambda =$$

$$(\cos \theta - \cos \theta) \lambda = \lambda - \lambda = 0$$

$$\frac{d}{\lambda} \gamma =$$

$$d \gamma = \lambda$$

$$\frac{d}{\lambda} = \gamma$$

$$c = v \lambda$$

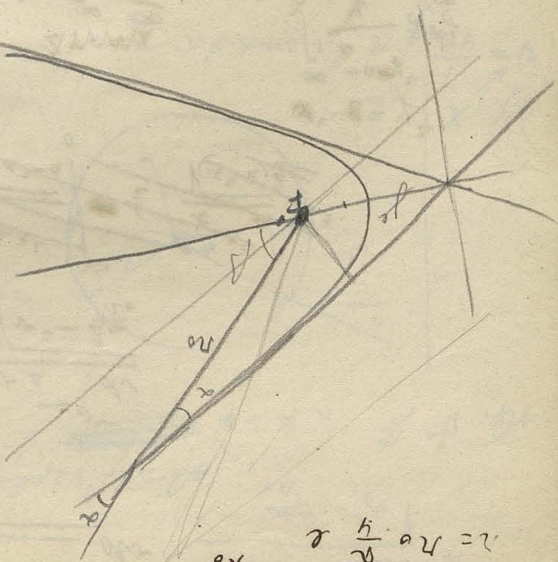
$$c = v \lambda$$

$$c = v \lambda$$

$$\lambda = \frac{c}{v} = \frac{c}{v}$$

$$\lambda = \frac{c}{v}$$

$$\lambda = \frac{c}{v}$$



$$\lambda = \frac{c}{v} = \frac{c}{v}$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}x^2} dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}x^2} dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}x^2} dx$$

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$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}x^2} dx$$

$$N dr = \frac{h_0}{2\pi} \int_{-\infty}^{\infty} \sqrt{1 - \frac{2m}{r_0}} \frac{dr}{r_0}$$

$$((2)) = \int_{-\infty}^{\infty} (h) N dr$$

$$a = \frac{2}{2r}$$

$$= \frac{3 \pi h_0}{2\pi}$$

$$\frac{3}{2} a = (c^2)$$

0.01

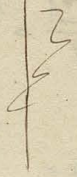
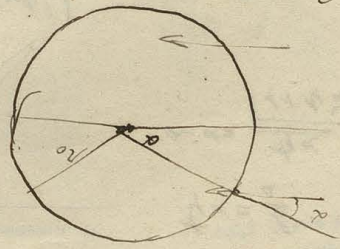
$$N = \frac{1}{2} \frac{h_0}{r} e^{-\alpha r^2}$$

$$= \frac{1}{2} \frac{h_0}{r_0} \frac{\sqrt{v_0^2 - \frac{2m}{r_0}}}{r_0}$$

$$= \frac{1}{2} \int_{-\infty}^{\infty} \frac{1}{r_0} \sin^2 x dx$$

$$((2)) = \frac{1}{2} \int_{-\infty}^{\infty} \pi \sin^2 x dx$$

$$\int_{-\infty}^{\infty} \sin^2 x dx = \int_{-\infty}^{\infty} \frac{1 - \cos 2x}{2} dx$$

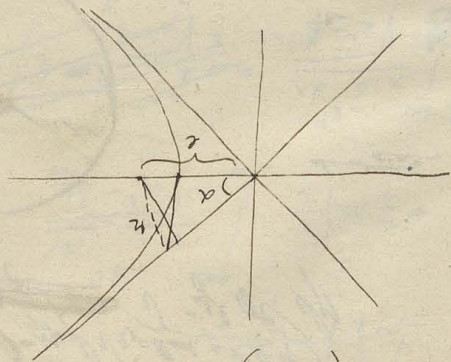


~~10. The ...~~
~~11. The ...~~

$$h = \frac{h_0 v_0}{\sqrt{v_0^2 - \frac{2m}{r_0}}}$$

Die Konstruktion des Theorems.

Watten (Verfahren zur Konstruktion des Theorems)



$$r^2 + h^2 = r^2$$

$$h = r \sin \alpha$$

$$h \alpha = \frac{r}{2}$$

$$r \alpha = \frac{h \alpha}{1 + h \alpha} = \frac{r}{2 + r}$$

$r = \frac{r}{2}$ (oder andersherum)

$$\left. \begin{aligned} h &= r \sin \alpha \\ e &= r_0 \sin \alpha \end{aligned} \right\} \begin{aligned} m + h &= e \\ m + r \sin \alpha &= r \sin \alpha \end{aligned}$$

$$m + r \sin \alpha = r \sin \alpha$$

$$x = x + m$$

$$m^2 y^2 - h^2 x^2 + 2c x \sqrt{m^2 + h^2} - 2m h c x + h^2 c^2 = h^2 c^2 +$$

$$+ c^2 y - 2c m \sqrt{m^2 + h^2}$$

$$= \sqrt{h^2 (m^2 + h^2)} + c^2 y -$$

$$- 2c^2 \sqrt{m^2 + h^2}$$

$$= \frac{h}{2} \sqrt{m^2 + h^2} + c^2 y$$

$$= c^2 y - c^2 \sqrt{m^2 + h^2}$$

$$= -c^2 \frac{h}{2}$$

$$a = \frac{h}{c} \quad b = \frac{r}{c}$$

$$x^2 \frac{h^2}{m^2} - y^2 \frac{c^2}{h} = 1$$

$$1 - \cos \beta \cos \gamma = \sin \beta \sin \gamma$$

$$1 \equiv \cos \beta \cos \gamma + \sin \beta \sin \gamma \neq 0$$

$$1 - \cos \beta \cos \gamma \geq \sin \beta \sin \gamma$$

$$1 - 2 \cos \beta \cos \gamma + (\sin \beta \cos \gamma + \cos \beta \sin \gamma)^2 \geq \sin^2 \beta \sin^2 \gamma$$

$$\begin{aligned} (1 - 2 \cos \beta \cos \gamma + \sin^2 \beta \sin^2 \gamma) &\geq 0 \\ (\cos \beta - \cos \gamma)^2 &\geq 0 \end{aligned}$$

$$1 - \cos \beta \cos \gamma + \sin \beta \sin \gamma$$

$$2(\log 2 - 1)$$

$$\log 2 - 1 + \frac{1}{2} + \log 2 - 1 - 1$$

$$\begin{aligned} \frac{dx}{dt} &= \frac{8 \text{ Ft}}{600.000} + \frac{8 \text{ Ft}}{1100.000} \\ \frac{dx}{dt} &= \frac{8 \text{ Ft}}{600.000} \end{aligned}$$

also in einem Winkel oder
 Einheitskreis mit 10 m
 in einem Jahr = 500 m = 500000

Einheitskreis

4. 7. 5. 1000
2001

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} - \frac{d\sigma}{d\Omega}$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} - \frac{d\sigma}{d\Omega}$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} - \frac{d\sigma}{d\Omega}$$

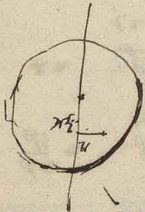
26. n
22
44
86.80

24. 62. 60



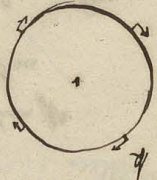
$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} - \frac{d\sigma}{d\Omega}$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} - \frac{d\sigma}{d\Omega}$$



$$= \frac{1}{600,000}$$

$$\frac{1}{k} = \frac{1}{2.400,000} \cdot \frac{1}{24.60}$$



Ergebnis der Rotation

$$\begin{array}{r} 100 \\ 1000 \\ \hline 1000 \\ 100 \\ \hline 1100 \end{array}$$

$$\frac{1}{x^2} = x^{-2}$$

$$\frac{d}{dx} x^{-2} = -2x^{-3}$$

$$= -\frac{2}{x^3}$$

$$= -\frac{2}{x^3}$$

$$= -\frac{2}{x^3}$$

$$= -\frac{2}{x^3}$$

$$= -\frac{2}{x^3}$$

$$= -\frac{2}{x^3}$$

$$\frac{d}{dx} x^{-2} = -\frac{2}{x^3}$$

(Nennwert Wert mit Formeln und Werten)

$\Delta T_{0}^{91x} = 365.021$
 $= 76.000 \quad 153 \text{ sec} = 27 \text{ Minuten!}$
 $\frac{730}{365}$ oder in einem Halbjahr

$$\begin{array}{r} 2988 : 141 = 0.21 \\ \underline{498} \\ 36 : 168 \end{array}$$

$\frac{3}{1-0.21} = \frac{3}{0.79} = 3.80$
 $\frac{23030' / 900.2}{23030' / 2} = 1.00$
 $\frac{37.30}{0.883 \cdot 188.2}$

$\Delta T_{0}^{91x} = 4.91$
 $M(a=5)$
 $\frac{365}{2}$

$\Delta T_{0}^{24x} = 4 \text{ ans}$

≠ ans

$= - \frac{6.60.60}{10.000} \cdot \frac{1}{2} \text{ ans} \cdot \delta = \frac{2.16}{2} \text{ ans} \cdot \delta$

$\Delta T_{0}^{6x} = - 6x \frac{1}{2} \text{ ans} \cdot \delta$

$\varphi = 90^\circ$ Nordpol: also in beiden Lag ≠ 100%
 (Nennwert mit am Berg und Meer)

Nennwert: $\Delta T_0 = - \frac{3}{2} \frac{1}{R} \text{ sec} \neq \frac{1}{2} \text{ sec}$
 $\frac{1}{R} = 1''$

$= - 2.16 \left[\text{ans} \cdot \delta + \frac{1}{2} \text{ ans} \cdot \delta \right]$
 Nennwert = $\frac{1}{2}$

$\Delta T_{0}^{6x} = - \frac{6.60.60}{10.000} \frac{1}{2} \left[\text{ans} \cdot \delta + \frac{1}{2} \text{ ans} \cdot \delta \right]$

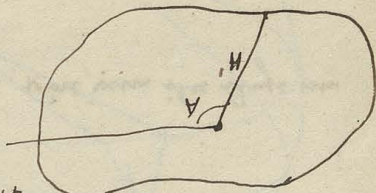
$\varphi = 45^\circ$
 $\sin \varphi = \cos \varphi = \frac{\sqrt{2}}{2}$

6 Stunden große von. Somit für alle Punkte unmittelbar
 Dargestellte Richtung durch, so kann von einem Punkte, um

$$\begin{aligned}
 & H^2 \left\{ m^2 A_{mp} + m^2 p^2 + 2f m A_{mp} \cdot H \right. \\
 & \qquad \qquad \qquad \left. = f^2 (m^2 p^2 - m^2) \right. \\
 & \qquad \qquad \qquad m^2 A_{mp} + 1 - m^2 p^2 \\
 & \qquad \qquad \qquad = 1 - m^2 p^2 \\
 & \qquad \qquad \qquad H^2 + 2H f m A_{mp} = f^2 \frac{1 - m^2 p^2}{1 - m^2 p^2} \\
 & \qquad \qquad \qquad H^2 = \frac{f m A_{mp} \cdot p^2}{m^2 A_{mp} + p^2} + \frac{f^2 m^2 A_{mp}^2}{m^2 A_{mp} + p^2} \\
 & \qquad \qquad \qquad = -f m A_{mp} \cdot p^2 \pm \sqrt{f^2 (m^2 p^2 - m^2) (m^2 A_{mp} + p^2)} + f^2 m^2 A_{mp}^2
 \end{aligned}$$

$$\sin A \sin \delta + \cos A \sin \gamma \cdot H^2 + 2 \sin A \sin \delta \cdot H = \sin \gamma \cdot f - \sin^2 \gamma \cdot H^2$$

$$\sin A \sin \delta + \cos A \sin \gamma = \sin \gamma \cdot f - \sin^2 \gamma \cdot H^2$$



~~$H^2 \sin A \sin \gamma = \sin \gamma \cdot f - \sin^2 \gamma \cdot H^2$~~

~~$H^2 \sin \gamma = f - \sin^2 \gamma \cdot H^2$~~

$$\cos A = \frac{H \sin \gamma}{-f \sin \delta + \sin \gamma \sqrt{f^2 - H^2}}$$

$$\cos A = \frac{H \sin \gamma}{-f \sin \delta + \sin \gamma \sqrt{1 - \frac{H^2}{f^2}}}$$

~~$H^2 \sin \gamma = f - \sin^2 \gamma \cdot H^2$~~

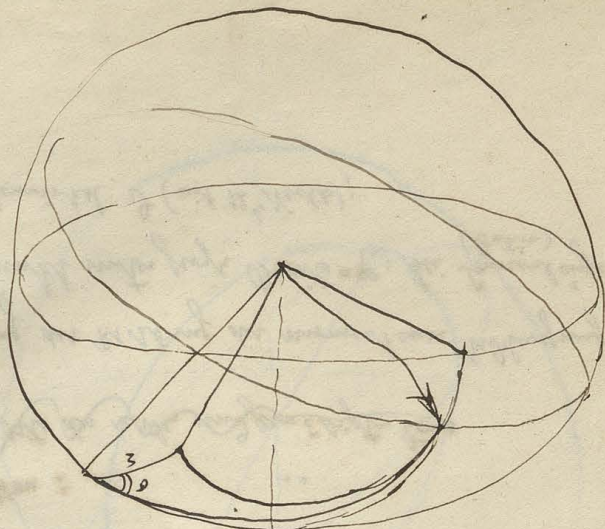
$$\cos \delta = \frac{f}{H}$$

$$H = f \cdot \cos(A)$$

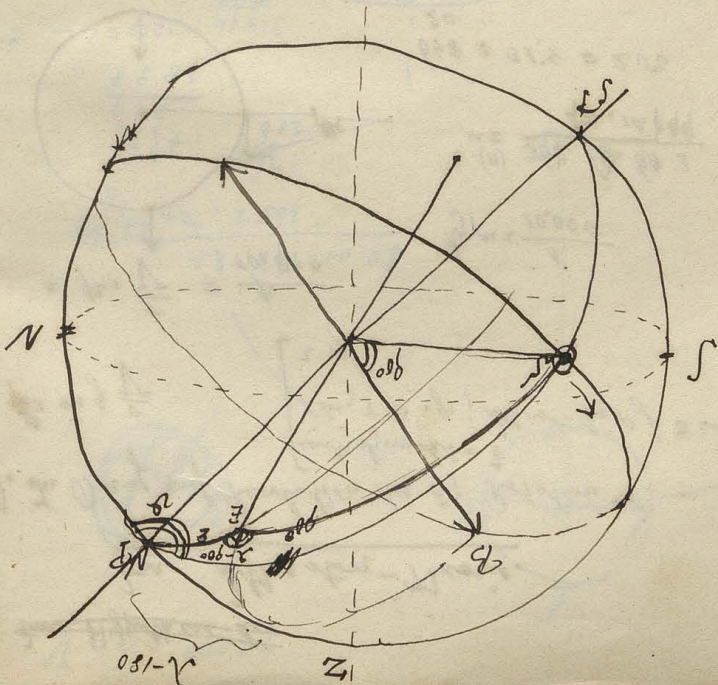
$$\cos A = \frac{\cos \delta \sin \gamma}{-f \sin \delta + \sin \gamma \sqrt{1 - \frac{H^2}{f^2}}}$$

$$H = f \cdot \cos \delta \left[\cos \delta = \sin \gamma \sqrt{1 - \frac{H^2}{f^2}} + \sin \delta \sin A \right]$$

Componente welche in die Vertikale fällt:



Ubergangsrichtung der Erde ist zunächst auf eine Ebene
 welche durch die Sonne und die Ost der Ekliptik gelegt wird.



~~von Hypothesen über~~

von Attraction - Theorie.

I. In Bezug auf Rotation der δ ; Rotation = 0.

$\left[\begin{array}{l} \text{Centrifugalkraft} = f \\ \text{mit } \rho \text{ in } \text{cm}^3, \text{ Abstand } = r \text{ in } \text{cm} \end{array} \right]$

$$f = m g \frac{v}{r}$$

$$= \rho \cdot \frac{V}{r} = \frac{10,000}{r}$$

$$\text{kgm} = \frac{1}{10,000}$$

$$(ii) \frac{368 \cdot 68 \cdot 68 \cdot 3}{2 \cdot 10,000} = n$$

$$698 : 315 = 20.6$$

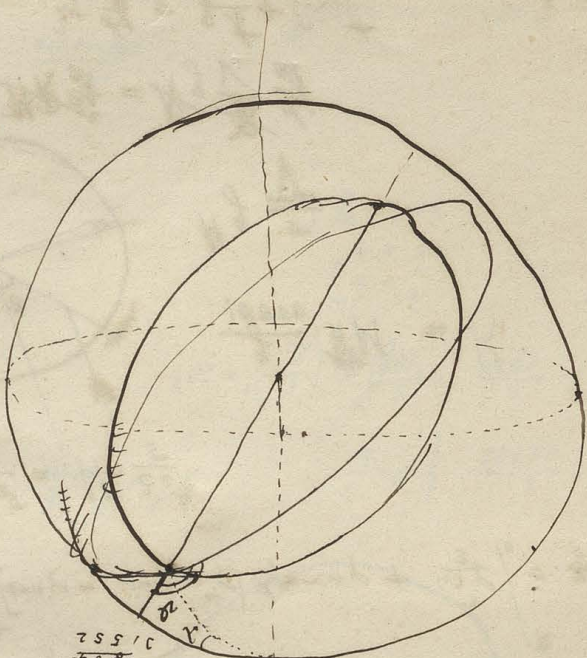
$$n = 20.6''$$

oder in Folge einer Lageänderung des Verticallinien um $40''$!!

Abmessungen:

alle die die Höhe und Breite betreffen

Bestimmung der Richtung der momentanen Ablenkung für
 die Höhen Winkel mit der Größe δ ; die Sonnenhöhe = 24°
 und Stundenwinkel θ (mit 12° North);



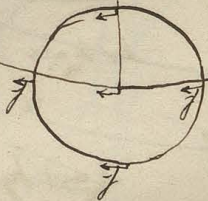
20.02 - 4 - 9888 = 4

$$\begin{array}{r} 7888 \\ 6 \\ \hline 219 \\ 185 \\ \hline 7304 \end{array}$$

36.9 246

$$\frac{365.2 \cdot 24}{6} = 1461.6$$

S



$\Delta = 28$

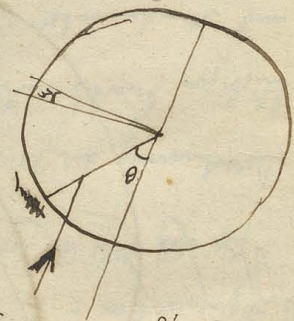
$$C = C_1 e^{\frac{8V}{T}} + C_2$$

$$M \frac{dy}{dt} = \frac{10}{7} \frac{dy}{dt} = \frac{10}{7} \frac{dy}{dt}$$

$$\frac{10}{7} \frac{dy}{dt} = \frac{10}{7} \frac{dy}{dt}$$

$$M g \frac{V}{v}$$

$$K = \frac{W}{g} \cdot \frac{g}{v}$$



$$T = \frac{8\pi a^5}{15} = M \cdot \frac{2a^2}{5}$$

$$\frac{3}{2} = \int_0^{\frac{\pi}{2}} \left[\int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi + \int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi \right] d\theta$$

$$= \frac{2\pi a^5}{5} \int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi$$

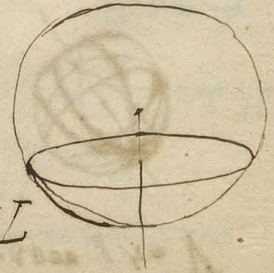
$$T = 2\pi \int_0^{\frac{\pi}{2}} \int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi \, d\theta$$

$$= \frac{2\pi a^5}{5}$$

$$\int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi = \int_0^{\frac{\pi}{2}} \sin^2 \phi \, d\phi \cdot \sin \phi$$

$$= \frac{2\pi a^5}{5} \int_0^{\frac{\pi}{2}} \sin^2 \phi \, d\phi$$

$$I = 2\pi \int_0^{\frac{\pi}{2}} \int_0^{\frac{\pi}{2}} \sin^3 \phi \, d\phi \, d\theta$$



$$\frac{2}{1000000} = 2 \times 10^{-6}$$

$$\frac{39000000}{10^1}$$

$$= 500000 \cdot 10^{-15} = 20000$$

$$\frac{2}{4} = 0.5$$

$$1' = 4$$

$$10' = 4m$$

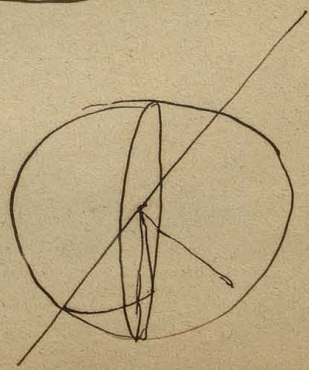
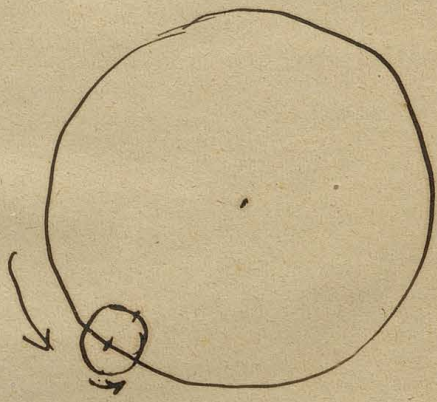
20"

$$0.2 = \left(\frac{5.2}{1}\right)' = \left(\frac{1}{1}\right)' = \frac{1000}{9} = 3$$

$$\frac{15 \frac{1}{2}}{15 \frac{1}{2}} = 1$$

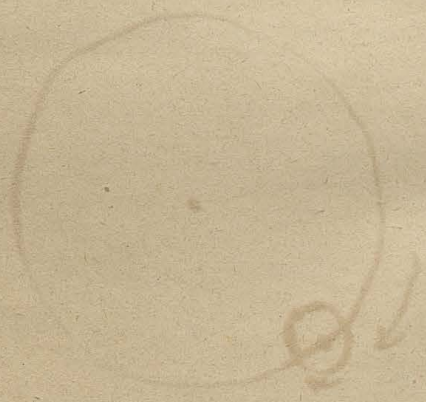
$$\frac{10000}{10000} = 1$$

$$\frac{10000}{1} = 10000$$



$$160.0 + 160 \frac{3}{2}$$

BJ



immer - eindeutige Art der Fortsetzung
der Fortsetzung - eindeutige Art der Fortsetzung
der Fortsetzung - eindeutige Art der Fortsetzung

~~$1 + \sin \varphi = x$~~
 ~~$-\cos \varphi d\varphi = dx$~~

$d\varphi = \frac{-dx}{\sqrt{a^2 - (x-1)^2}}$

$\frac{a}{b} = \frac{\frac{2\sqrt{2}}{2\sqrt{2}} \sqrt{2}}{\frac{2\sqrt{2}}{2\sqrt{2}} \sqrt{2}} = \frac{2\sqrt{2}}{2\sqrt{2}} = \frac{2}{2} = 1$

$\frac{1}{1 + \sin \varphi} = x$

$1 = x + \sin \varphi$

$\sin \varphi = \frac{1-x}{\cos \varphi} \quad \cos \varphi = \frac{\sqrt{a^2 - (1-x)^2}}{\cos \varphi}$

$-\cos \varphi d\varphi = dx$

$\frac{1}{2} \leftarrow \varphi$

$y = \int \frac{x dx}{\sqrt{(a^2-1)x^2 + 2x - 1}}$

$= \frac{-1}{a^2-1} \sqrt{(a^2-1)x^2 + 2x - 1} + \frac{1}{a^2-1} \int \frac{dx}{\sqrt{\dots}}$

$\sqrt{\dots} = \sqrt{\left(\frac{a^2-1}{(1+\sin \varphi)}\right)^2 + \frac{2}{1+\sin \varphi} - 1}$

$= \frac{1}{2\sqrt{(a^2-1)}} \ln \left| \frac{2 + 2(a^2-1)x + 2\sqrt{\dots}}{2 + 2(a^2-1) + 2\sqrt{(a^2-1)}\sqrt{(a^2-1)x^2 + 2x - 1}} \right|$

$= \frac{a^2 - x + 2 + 2\sin \varphi - 1 - 2\sin \varphi - a^2 \sin^2 \varphi}{(1 + \sin \varphi)^2}$

$= \frac{a^2 \sin \varphi}{1 + \sin \varphi}$

$y = \frac{1}{1-a^2} \frac{a^2 \sin \varphi}{1 + \sin \varphi} + \frac{1}{2(a^2-1)^{3/2}} \ln \dots$

$\frac{1 + \frac{a^2-1}{1+\sin \varphi} + \frac{\sqrt{a^2-1} \sin \varphi}{1+\sin \varphi}}{1 + \dots}$

$\frac{1 + \sin \varphi + a^2 - 1 + \sqrt{a^2-1} \sin \varphi}{\cos \varphi + 1 - \sqrt{a^2-1} \sin \varphi}$

$\frac{2\sqrt{a^2-1}}{2\sqrt{a^2-1}} = \frac{2\sqrt{a^2-1}}{2\sqrt{a^2-1}}$

$\frac{2\sqrt{a^2-1}}{2\sqrt{a^2-1}}$

Am 2. September des k. k. Reichs-Ratungsrathes Kaiserl. H. H.

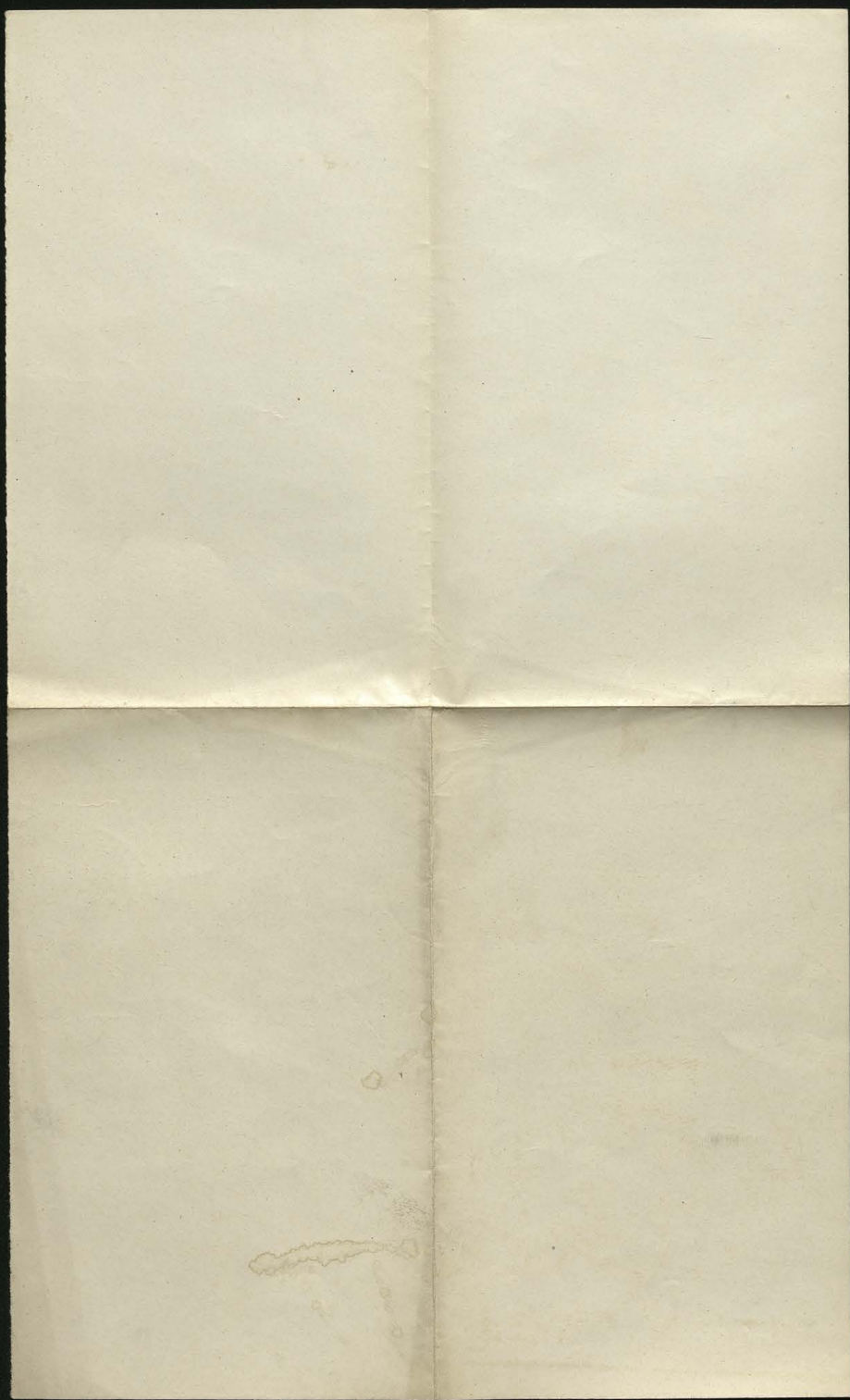
aus dem Jahre 1895

Programm

der dritten Jahrgangweisen Zusammenkunft

- 1) Schönen Geist von Hansch
- 2) Conventus Dichter — Bauer — Leitung von H. H.
- 3) Eine schone Geschichte von Leop. Haber
- 4) Sters varius für H. H. — H. H.
- 5) Die wunliche Verlage — H. H.
- 6) Die wunliche Verlage — H. H.
- 7) Die wunliche Verlage — H. H.
- 8) Die wunliche Verlage — H. H.
- 9) Die wunliche Verlage — H. H.
- 10) Die wunliche Verlage — H. H.
- 11) Die wunliche Verlage — H. H.
- 12) Die wunliche Verlage — H. H.
- 13) Die wunliche Verlage — H. H.
- 14) Die wunliche Verlage — H. H.

Wien, am 11. April 1895 Anfang 2000



John Tenney's 5

